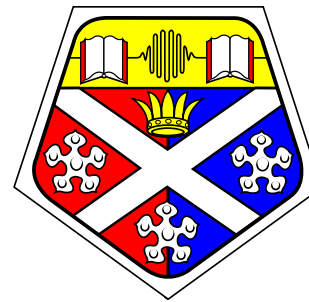


Multigrid Solution of Discrete Convection-Diffusion Equations

Alison Ramage
Dept of Mathematics
University of Strathclyde
Glasgow, Scotland



UNIVERSITY OF
STRATHCLYDE



UNIVERSITY OF
MARYLAND

Howard Elman
Dept of Computer Science
University of Maryland
College Park, MD, USA

Overview

- background
 - convection-diffusion problems
 - multigrid methods
- practical multigrid issues
 - **approximation** and **smoothing** properties
 - convergence analysis
- model problem Fourier analysis
 - matrix transformation
 - comparison with **semiperiodic** problem
 - implications for **Dirichlet** problem

Convection-Diffusion in 2D

$$\begin{aligned} -\epsilon \nabla^2 u(x, y) + \mathbf{w} \cdot \nabla u(x, y) &= f(x, y) \quad \text{in } \Omega \in \mathbb{R}^2 \\ u(x, y) &= g \quad \text{on } \partial\Omega \end{aligned}$$

divergence-free convective velocity ('wind') \mathbf{w}

diffusion parameter $\epsilon \ll 1$

discretisation parameter h

Convection-Diffusion in 2D

$$\begin{aligned} -\epsilon \nabla^2 u(x, y) + \mathbf{w} \cdot \nabla u(x, y) &= f(x, y) \quad \text{in } \Omega \in \mathbb{R}^2 \\ u(x, y) &= g \quad \text{on } \partial\Omega \end{aligned}$$

divergence-free convective velocity ('wind') \mathbf{w}

diffusion parameter $\epsilon \ll 1$

discretisation parameter h

mesh Péclet number $P_h = \frac{\|\mathbf{w}\| h}{2\epsilon}$

Boundary Layers and Oscillations

- Galerkin finite element method

$$\epsilon(\nabla u_h, \nabla v_h) + (\mathbf{w} \cdot \nabla u_h, v_h) = (f, v_h) \quad \forall v_h \in V_h$$

Boundary Layers and Oscillations

- Galerkin finite element method

$$\epsilon(\nabla u_h, \nabla v_h) + (\mathbf{w} \cdot \nabla u_h, v_h) = (f, v_h) \quad \forall v_h \in V_h$$

- solution features:

exponential and **characteristic** boundary layers

Boundary Layers and Oscillations

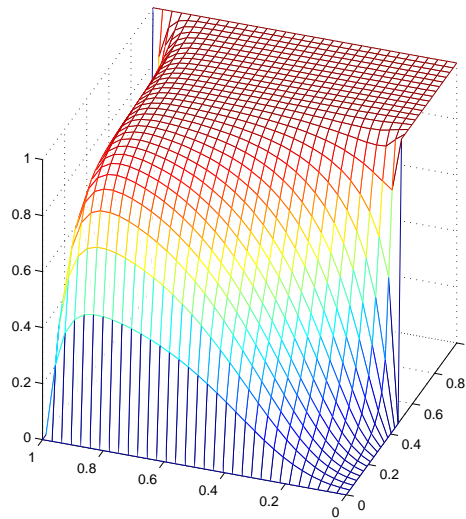
- Galerkin finite element method

$$\epsilon(\nabla u_h, \nabla v_h) + (\mathbf{w} \cdot \nabla u_h, v_h) = (f, v_h) \quad \forall v_h \in V_h$$

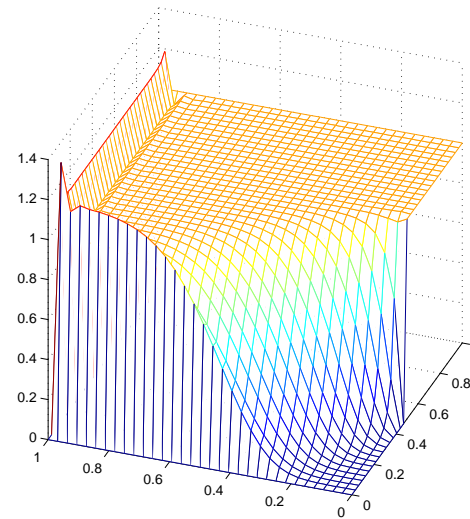
- solution features:

exponential and characteristic boundary layers

- oscillations observed in discrete solutions for $P_h > 1$



$P_e = 0.5$



$P_e = 2$

Streamline Diffusion Method

streamline diffusion FEM, square bilinear elements

$$\begin{aligned} \epsilon(\nabla u_h, \nabla v_h) &+ (\mathbf{w} \cdot \nabla u_h, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (\mathbf{w} \cdot \nabla u_h, \mathbf{w} \cdot \nabla v_h) \\ &= (f, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (f, \mathbf{w} \cdot \nabla v_h) \quad \forall v_h \in V_h \end{aligned}$$

Streamline Diffusion Method

streamline diffusion FEM, square bilinear elements

$$\begin{aligned}\epsilon(\nabla u_h, \nabla v_h) &+ (\mathbf{w} \cdot \nabla u_h, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (\mathbf{w} \cdot \nabla u_h, \mathbf{w} \cdot \nabla v_h) \\ &= (f, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (f, \mathbf{w} \cdot \nabla v_h) \quad \forall v_h \in V_h\end{aligned}$$

- $P_h \leq 1$: $\delta = 0$

Galerkin FEM

- $P_h > 1$: $\delta = \frac{1}{2} - \frac{\epsilon}{h}$

Streamline Diffusion

Multigrid Ideas

- fine grid (h), coarse grid ($2h$)

Multigrid Ideas

- fine grid (h), coarse grid ($2h$)
- decompose a grid function into components in two subspaces

approximate inverse operator
for components in subspace 1

smoothing iteration
rapidly reduces error components in subspace 2

Multigrid Ideas

- fine grid (h), coarse grid ($2h$)
- decompose a grid function into components in two subspaces

approximate inverse operator
for components in subspace 1

smoothing iteration
rapidly reduces error components in subspace 2

- recursive process on nested grids

Multigrid Ideas

- fine grid (h), coarse grid ($2h$)
- decompose a grid function into components in two subspaces

approximate inverse operator
for components in subspace 1

smoothing iteration
rapidly reduces error components in subspace 2

- recursive process on nested grids
- **optimal** in the sense of obtaining convergence rate independent of h

Issues for Convection-Diffusion

- **approximation**: choice of discretisation
 - oscillations on coarser grids?
 - grid transfer operators?

Issues for Convection-Diffusion

- **approximation**: choice of discretisation
 - oscillations on coarser grids?
 - grid transfer operators?
- **smoothing**: choice of relaxation method
 - direction of flow?
 - circular flows?

Issues for Convection-Diffusion

- **approximation**: choice of discretisation
 - oscillations on coarser grids?
 - grid transfer operators?
- **smoothing**: choice of relaxation method
 - direction of flow?
 - circular flows?
- multigrid can be implemented effectively for convection-diffusion problems

Convergence Analysis

- standard Poisson-type convergence analysis fails
- ideas for convection-diffusion less well-developed

Convergence Analysis

- standard Poisson-type convergence analysis fails
- ideas for convection-diffusion less well-developed
- various successful approaches
 - **perturbation arguments**
Bank (1981), Bramble, Pasciak and Xu (1988), Mandel (1986), Wang (1993)
 - **matrix-based methods**
Reusken (2002), Olishanskii and Reusken (2002)

Multigrid Method

- two-grid method: N_f (fine grid), N_c (coarse grid)

Multigrid Method

- two-grid method: N_f (fine grid), N_c (coarse grid)
- coefficient matrices: A_f (fine grid), A_c (coarse grid)

direct discretisation on coarse grid

Multigrid Method

- two-grid method: N_f (fine grid), N_c (coarse grid)
- coefficient matrices: A_f (fine grid), A_c (coarse grid)
- prolongation: bilinear interpolation P
- restriction: transpose of prolongation P^T

Multigrid Method

- two-grid method: N_f (fine grid), N_c (coarse grid)
- coefficient matrices: A_f (fine grid), A_c (coarse grid)
- prolongation: bilinear interpolation P
- restriction: transpose of prolongation P^T
- smoothing: line Gauss-Seidel S_A
- ν steps of pre-smoothing, no post-smoothing

Multigrid Convergence

- algebraic error $e_k = \hat{u} - u_k$

Multigrid Convergence

- algebraic error $\mathbf{e}_k = \hat{\mathbf{u}} - \mathbf{u}_k$
- two-grid iteration matrix $M = (I - PA_c^{-1}P^T A_f)S_A^\nu$
- error equation $\mathbf{e}_k = M\mathbf{e}_{k-1} = M^k\mathbf{e}_0$

Multigrid Convergence

- algebraic error $\mathbf{e}_k = \hat{\mathbf{u}} - \mathbf{u}_k$
- two-grid iteration matrix $M = (I - PA_c^{-1}P^T A_f)S_A^\nu$
- error equation $\mathbf{e}_k = M\mathbf{e}_{k-1} = M^k\mathbf{e}_0$
- convergence?

$$\|\mathbf{e}_k\| \leq \|M\|^k \|\mathbf{e}_0\|$$

convergence if $\|M\| < 1$

Two-Grid Convergence Analysis

AIM: find an upper bound for

$$\|M\|_2 = \|(I - PA_c^{-1}P^T A_f)S_A^\nu\|_2$$

Two-Grid Convergence Analysis

AIM: find an upper bound for

$$\|M\|_2 = \|(I - PA_c^{-1}P^T A_f)S_A^\nu\|_2$$

- Approach 1: write

$$M = (A_f^{-1} - PA_c^{-1}P^T)(A_f S_A^\nu) = M_A M_S$$

and bound $\|M_A\|_2$, $\|M_S\|_2$ separately

Two-Grid Convergence Analysis

AIM: find an upper bound for

$$\|M\|_2 = \|(I - PA_c^{-1}P^T A_f)S_A^\nu\|_2$$

- Approach 1: write

$$M = (A_f^{-1} - PA_c^{-1}P^T)(A_f S_A^\nu) = M_A M_S$$

and bound $\|M_A\|_2$, $\|M_S\|_2$ separately

- Approach 2: bound $\|M\|_2$ directly

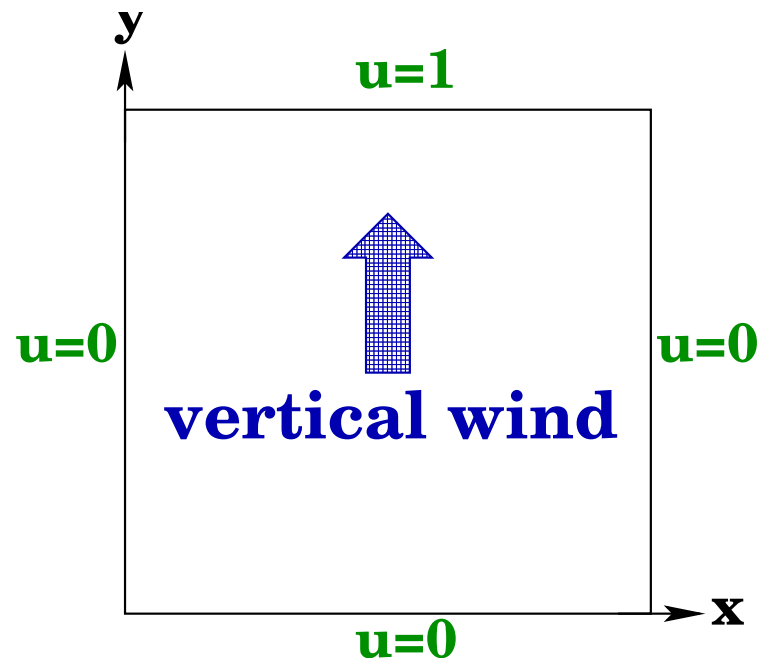
Model Problem

grid-aligned flow with vertical wind and $f = 0$

$$-\epsilon \nabla^2 u(x, y) + (0, 1) \cdot \nabla u(x, y) = 0$$

Dirichlet boundary conditions

square bilinear elements



Computational Molecule

parameters h, ϵ, δ

$$M_2 : \quad -\frac{1}{12}[(2\delta-1)h+4\epsilon] \quad -\frac{1}{3}[(2\delta-1)h+\epsilon] \quad -\frac{1}{12}[(2\delta-1)h+4\epsilon]$$



$$M_1 : \quad \frac{1}{3}(\delta h - \epsilon) \quad \leftarrow \quad \frac{4}{3}(\delta h + 2\epsilon) \quad \rightarrow \quad \frac{1}{3}(\delta h - \epsilon)$$



$$M_3 : \quad -\frac{1}{12}[(2\delta+1)h+4\epsilon] \quad -\frac{1}{3}[(2\delta+1)h+\epsilon] \quad -\frac{1}{12}[(2\delta+1)h+4\epsilon]$$

symmetric stencil

Coefficient Matrix

$$A = \begin{bmatrix} M_1 & M_2 & & & 0 \\ M_3 & M_1 & M_2 & & \\ & \ddots & \ddots & \ddots & \\ & & M_3 & M_1 & M_2 \\ 0 & & & M_3 & M_1 \end{bmatrix}$$

Coefficient Matrix

$$A = \begin{bmatrix} M_1 & M_2 & & & 0 \\ M_3 & M_1 & M_2 & & \\ & \ddots & \ddots & \ddots & \\ & & M_3 & M_1 & M_2 \\ 0 & & & M_3 & M_1 \end{bmatrix}$$

eigenvectors and eigenvalues:

$$\begin{aligned} M_1 \mathbf{v}_j &= \lambda_j \mathbf{v}_j, & \lambda_j &= m_{1c} + 2m_{1r} \cos \frac{j\pi}{N} \\ M_2 \mathbf{v}_j &= \sigma_j \mathbf{v}_j, & \sigma_j &= m_{2c} + 2m_{2r} \cos \frac{j\pi}{N} \\ M_3 \mathbf{v}_j &= \gamma_j \mathbf{v}_j, & \gamma_j &= m_{3c} + 2m_{3r} \cos \frac{j\pi}{N} \end{aligned}$$

$$\mathbf{v}_j = \sqrt{\frac{2}{N}} \left[\sin \frac{j\pi}{N}, \sin \frac{2j\pi}{N}, \dots, \sin \frac{(N-1)j\pi}{N} \right]^T$$

Transformation: Coefficient Matrix (1)

N_f^2 elements, n_f^2 unknowns ($n_f = N_f - 1$)

$$\hat{V}_f = [\mathbf{v}_1 \mathbf{v}_2 \dots \mathbf{v}_{n_f}], \quad V_f = \text{diag}(\hat{V}_f, \dots, \hat{V}_f)$$

$$M_1 \hat{V}_f = \hat{V}_f \Lambda, \quad M_2 \hat{V}_f = \hat{V}_f \Sigma, \quad M_3 \hat{V}_f = \hat{V}_f \Gamma$$

$$V_f^T A_f V_f = \hat{T}_f = \begin{bmatrix} \Lambda & \Sigma & & & 0 \\ \Gamma & \Lambda & \Sigma & & \\ & \ddots & \ddots & \ddots & \\ & & \Gamma & \Lambda & \Sigma \\ 0 & & & \Gamma & \Lambda \end{bmatrix}$$

Transformation: Coefficient Matrix (2)

permute into tridiagonal form:

$$\Pi_f^T \hat{T}_f \Pi_f = T_f = \begin{bmatrix} T_1 & & & & 0 \\ & T_2 & & & \\ & & \ddots & & \\ & & & T_{n_f-1} & \\ 0 & & & & T_{n_f} \end{bmatrix}$$

$$T_j = \text{tridiag}(\gamma_j, \lambda_j, \sigma_j)$$

$$A_f = Q_f T_f Q_f^T \quad Q_f = V_f \Pi_f$$

Transformation: Coefficient Matrix (2)

permute into tridiagonal form:

$$\Pi_f^T \hat{T}_f \Pi_f = T_f = \begin{bmatrix} T_1 & & & & 0 \\ & T_2 & & & \\ & & \ddots & & \\ & & & T_{n_f-1} & \\ 0 & & & & T_{n_f} \end{bmatrix}$$

$$T_j = \text{tridiag}(\gamma_j, \lambda_j, \sigma_j)$$

$$A_f = Q_f T_f Q_f^T \quad Q_f = V_f \Pi_f$$

$$\text{coarse grid: } A_c = Q_c T_c Q_c^T \quad Q_c = V_c \Pi_c$$

Transformation: Smoothing Matrix

block matrix splitting: $A_f = D_A - L_A - U_A$

Gauss-Seidel smoothing matrix:

$$S_A = (D_A - L_A)^{-1}U_A = I - (D_A - L_A)^{-1}A_f$$

Transformation: Smoothing Matrix

block matrix splitting: $A_f = D_A - L_A - U_A$

Gauss-Seidel smoothing matrix:

$$S_A = (D_A - L_A)^{-1}U_A = I - (D_A - L_A)^{-1}A_f$$

transformation:

$$S_A = Q_f S_T Q_f^T$$

where $S_T = I - (D_T - L_T)^{-1}T_f$ is block-diagonal

Transformation: Prolongation Matrix

2D prolongation matrix: $P = L \otimes L$

$$L^T = \begin{bmatrix} \frac{1}{2} & 1 & \frac{1}{2} & & & & & & & & \\ & & \frac{1}{2} & 1 & \frac{1}{2} & & & & & & \\ & & & & & \ddots & & & & & \\ & & & & & & \ddots & & & & \\ & & & & & & & \frac{1}{2} & 1 & \frac{1}{2} & \\ & & & & & & & & & & \end{bmatrix}$$

Transformation: Iteration Matrix (1)

$$\begin{aligned} M &= (I - PA_c^{-1}P^T A_f)S_A^\nu \\ &= (I - PQ_cT_c^{-1}Q_c^T P^T Q_fT_fQ_f^T)S_A^\nu \\ &= Q_f(I - \bar{P}T_c^{-1}\bar{P}^T T_f)Q_f^T(Q_fS_TQ_f^T)^\nu \\ &= Q_f \left(I - \bar{P}T_c^{-1}\bar{P}^T T_f \right) S_T^\nu Q_f^T \end{aligned}$$

$$\Rightarrow M = Q_f \bar{M} Q_f^T$$

where $\bar{M} = (I - \bar{P}T_c^{-1}\bar{P}^T T_f) S_T^\nu$

Transformation: Iteration Matrix (1)

$$\begin{aligned} M &= (I - PA_c^{-1}P^T A_f)S_A^\nu \\ &= (I - PQ_cT_c^{-1}Q_c^T P^T Q_fT_fQ_f^T)S_A^\nu \\ &= Q_f(I - \bar{P}T_c^{-1}\bar{P}^T T_f)Q_f^T(Q_fS_TQ_f^T)^\nu \\ &= Q_f \left(I - \bar{P}T_c^{-1}\bar{P}^T T_f \right) S_T^\nu Q_f^T \end{aligned}$$

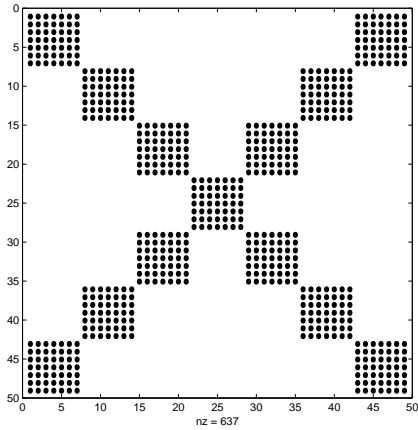
$$\Rightarrow M = Q_f \bar{M} Q_f^T$$

where $\bar{M} = (I - \bar{P}T_c^{-1}\bar{P}^T T_f) S_T^\nu$

Q_f is orthogonal:

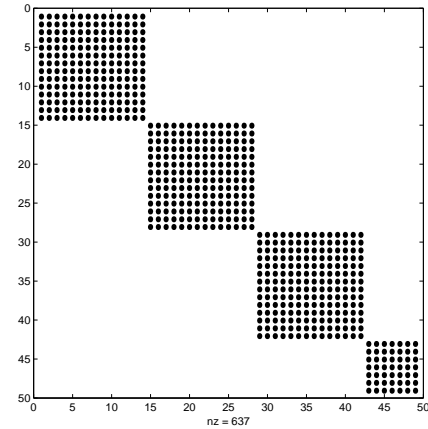
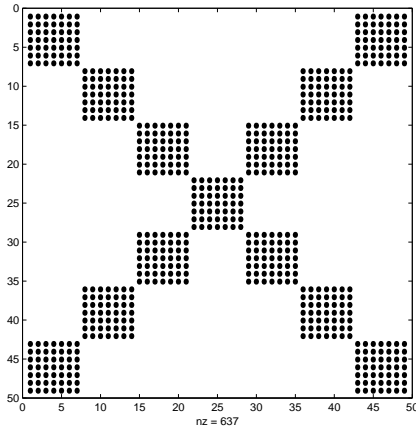
$$\|M\|_2 = \|\bar{M}\|_2$$

Transformed Iteration Matrix (2)



$$\begin{bmatrix} \mathbf{B}_1 & & & & & & \mathbf{C}_1 \\ & \mathbf{B}_2 & & & & & \mathbf{C}_2 \\ & & \mathbf{B}_3 & & & & \mathbf{C}_3 \\ & & & \mathbf{B}_4 & & & \mathbf{C}_4 \\ & & & & \mathbf{B}_5 & & \mathbf{C}_5 \\ & & & & & \mathbf{B}_6 & \mathbf{C}_6 \\ \mathbf{C}_7 & & & & & & \mathbf{B}_7 \end{bmatrix}$$

Transformed Iteration Matrix (2)

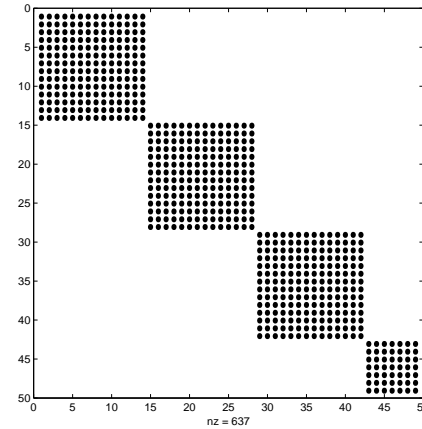
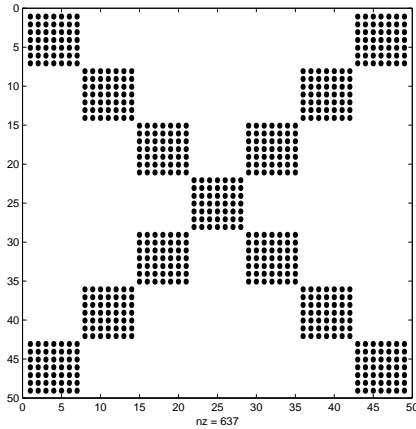


$$\begin{bmatrix} \mathbf{B}_1 & & & & & & \mathbf{C}_1 \\ & \mathbf{B}_2 & & & & & \\ & & \mathbf{B}_3 & & & & \mathbf{C}_2 \\ & & & \mathbf{B}_4 & & & \mathbf{C}_3 \\ & & & & \mathbf{B}_5 & & \\ & & & & & \mathbf{B}_6 & \\ \mathbf{C}_7 & & & & & & \mathbf{B}_7 \end{bmatrix}$$



$$\begin{bmatrix} \mathbf{B}_1 \mathbf{C}_1 & & & & & & \\ \mathbf{C}_7 \mathbf{B}_7 & & & & & & \\ & & \mathbf{B}_2 \mathbf{C}_2 & & & & \\ & & & \mathbf{C}_6 \mathbf{B}_6 & & & \\ & & & & \mathbf{B}_3 \mathbf{C}_3 & & \\ & & & & & \mathbf{C}_5 \mathbf{B}_5 & \\ & & & & & & \mathbf{B}_4 \end{bmatrix}$$

Transformed Iteration Matrix (2)



$$\begin{bmatrix} \mathbf{B}_1 & & & & & & \mathbf{C}_1 \\ & \mathbf{B}_2 & & & & & \mathbf{C}_2 \\ & & \mathbf{B}_3 & & & & \mathbf{C}_3 \\ & & & \mathbf{B}_4 & & & \mathbf{C}_4 \\ & & & & \mathbf{B}_5 & & \mathbf{C}_5 \\ & & & & & \mathbf{B}_6 & \mathbf{C}_6 \\ \mathbf{C}_7 & & & & & & \mathbf{B}_7 \end{bmatrix}$$



$$\begin{bmatrix} \mathbf{B}_1 \mathbf{C}_1 & & & & & & \\ \mathbf{C}_7 \mathbf{B}_7 & & & & & & \\ & \mathbf{B}_2 \mathbf{C}_2 & & & & & \\ & \mathbf{C}_6 \mathbf{B}_6 & & & & & \\ & & \mathbf{B}_3 \mathbf{C}_3 & & & & \\ & & \mathbf{C}_5 \mathbf{B}_5 & & & & \\ & & & & \mathbf{B}_4 & & \end{bmatrix}$$

$$\|\bar{M}\|_2 = \max \left\{ \max_{j=1, \dots, n_c} \left\| \begin{bmatrix} B_j & C_j \\ C_k & B_k \end{bmatrix} \right\|_2, \|B_{N_c}\|_2 \right\}, \quad k = N_f - j$$

The Story So Far...

- $n_f^2 \times n_f^2$ two-grid iteration matrix M
- Fourier transformation converts 2D problem to a set of n_f problems with 1D structure
- $\|M\|_2$ can be found from norms of N_c smaller problems

n_c of size $2n_f \times 2n_f$, 1 of size $n_f \times n_f$

The Story So Far...

- $n_f^2 \times n_f^2$ two-grid iteration matrix M
- Fourier transformation converts 2D problem to a set of n_f problems with 1D structure
- $\|M\|_2$ can be found from norms of N_c smaller problems
 n_c of size $2n_f \times 2n_f$, 1 of size $n_f \times n_f$
- **IDEA**: analyse semiperiodic version of the problem
 n_c of size $2N_f \times 2N_f$, 1 of size $N_f \times N_f$
- gain insight into Dirichlet problem behaviour?

Semiperiodic problem

- B_j, C_j are replaced by periodic versions, e.g.

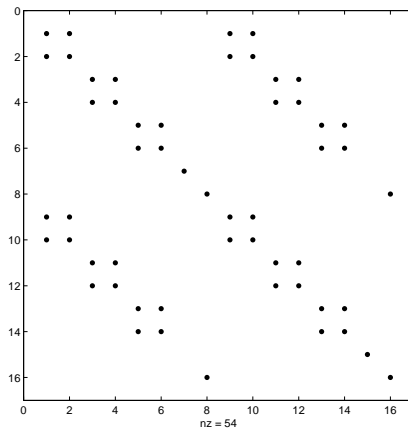
$$B_j^{per} = [I - \bar{P}_j^{per} (T_c^{per})_j^{-1} (\bar{P}_j^{per})^T (T_f^{per})_j] S_j^{per}$$

Semiperiodic problem

- B_j, C_j are replaced by periodic versions, e.g.

$$B_j^{per} = [I - \bar{P}_j^{per} (T_c^{per})_j^{-1} (\bar{P}_j^{per})^T (T_f^{per})_j] S_j^{per}$$

- transform using coarse grid periodic eigenvectors
- B_j^{per}, C_j^{per} become block diagonal with 2×2 blocks

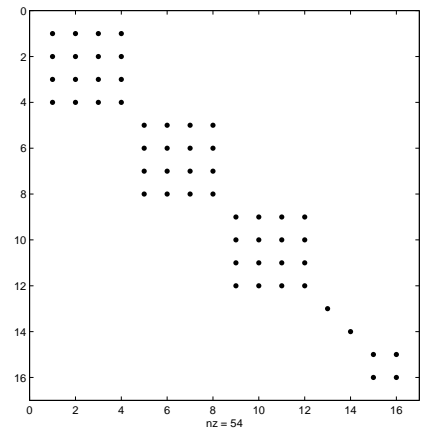
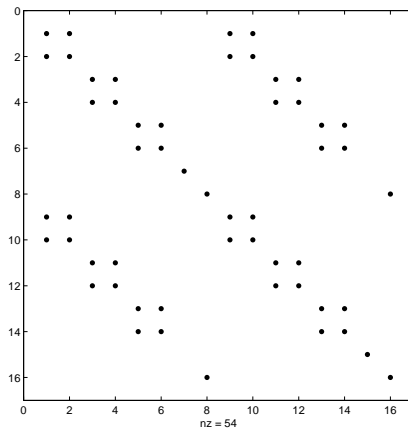


Semiperiodic problem

- B_j, C_j are replaced by periodic versions, e.g.

$$B_j^{per} = [I - \bar{P}_j^{per} (T_c^{per})_j^{-1} (\bar{P}_j^{per})^T (T_f^{per})_j] S_j^{per}$$

- transform using coarse grid periodic eigenvectors
- B_j^{per}, C_j^{per} become block diagonal with 2×2 blocks
- permute into block diagonal form

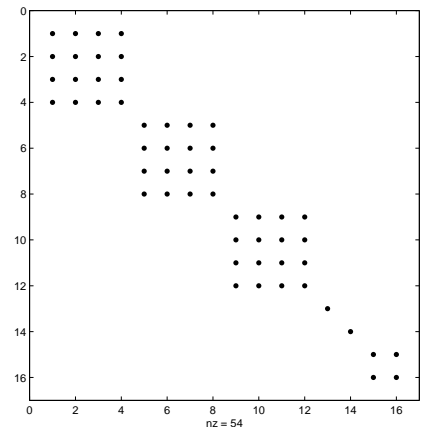
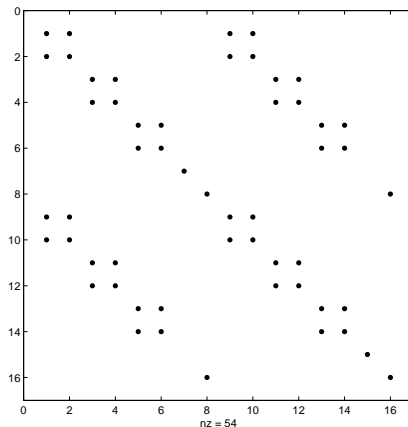


Semiperiodic problem

- B_j, C_j are replaced by periodic versions, e.g.

$$B_j^{per} = [I - \bar{P}_j^{per} (T_c^{per})_j^{-1} (\bar{P}_j^{per})^T (T_f^{per})_j] S_j^{per}$$

- transform using coarse grid periodic eigenvectors
- B_j^{per}, C_j^{per} become block diagonal with 2×2 blocks
- permute into block diagonal form



- 2-norm given by maximum 2-norm of the 4×4 blocks

Analytic result

- with semiperiodic approximation, when $P_h > 1$

$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

independent of ϵ

Analytic result

- with semiperiodic approximation, when $P_h > 1$

$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

- as h is small in practice,

$$\|M^{per}\|_2 \simeq \frac{\sqrt{2}}{5^\nu}$$

Analytic result

- with semiperiodic approximation, when $P_h > 1$

$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

- as h is small in practice,

$$\|M^{per}\|_2 \simeq \frac{\sqrt{2}}{5^\nu}$$

- when $P_h < 1$, analysis is more detailed but good approximations to $\|M^{per}\|_2$ can be derived

Analytic result

- with semiperiodic approximation, when $P_h > 1$

$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

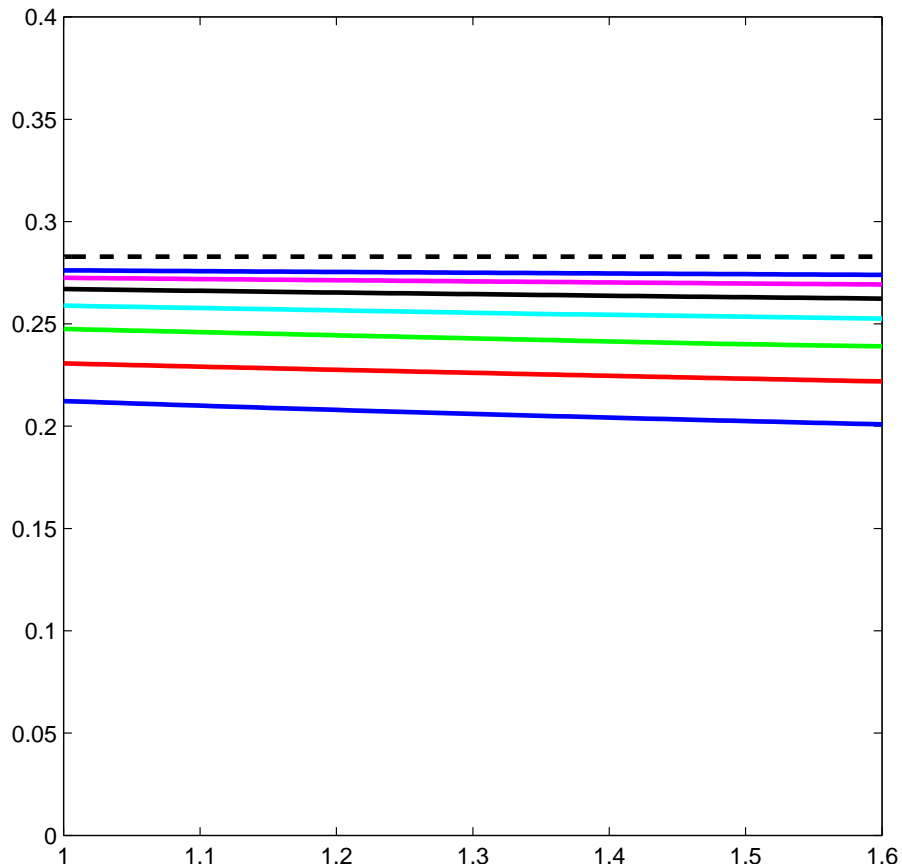
- as h is small in practice,

$$\|M^{per}\|_2 \simeq \frac{\sqrt{2}}{5^\nu}$$

- when $P_h < 1$, analysis is more detailed but good approximations to $\|M^{per}\|_2$ can be derived

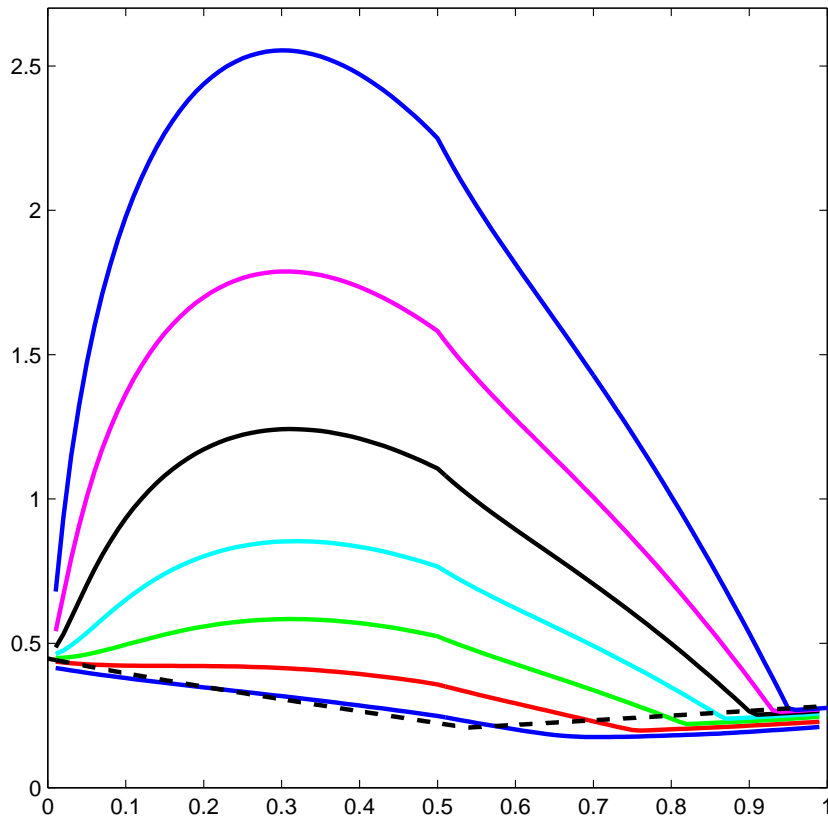
Question: Does this **semiperiodic** analysis correctly predict **Dirichlet** problem behaviour?

Model Problem Results (1)



- $\|M\|_2$ vs P_h
- $P_h \geq 1$ only
- semiperiodic: dashed line
- Dirichlet: solid lines
- h fixed for each line
- $h = \frac{1}{8}$ to $h = \frac{1}{512}$
- $\nu = 1$
- semiperiodic: $\frac{\sqrt{2}}{5} \simeq 0.28$
- Dirichlet $\rightarrow \frac{\sqrt{2}}{5}$

Model Problem Results (2)



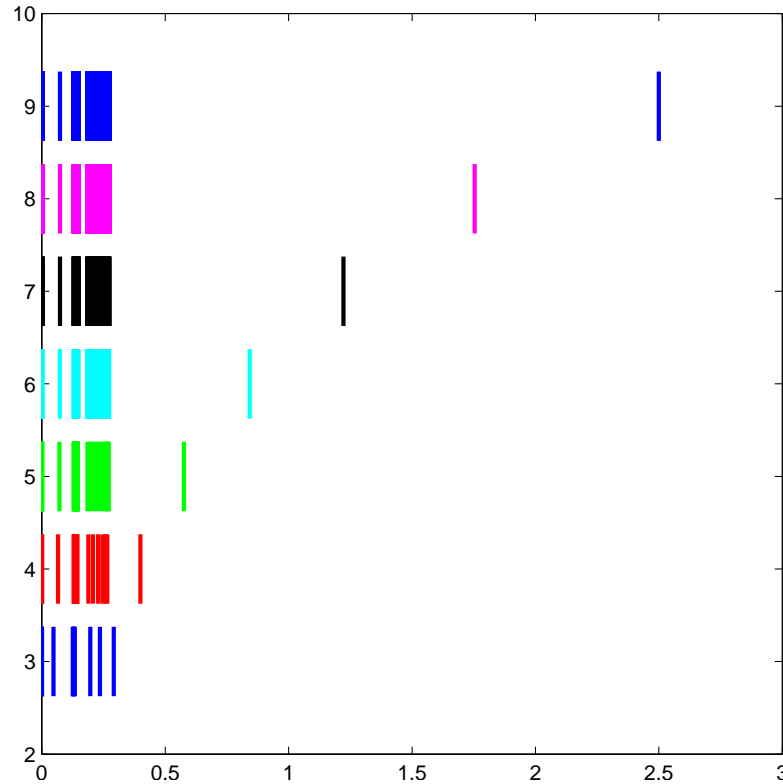
- $\|M\|_2$ vs P_h
- $P_h < 1$ only
- semiperiodic: dashed line
- Dirichlet: solid lines
- h fixed for each line
- $h = \frac{1}{8}$ to $h = \frac{1}{512}$
- $\nu = 1$
- not a good match
- MG may diverge!

Observations

- $\|M\|_2 = \sqrt{|\lambda_1(M^*M)|}$

Observations

- $\|M\|_2 = \sqrt{|\lambda_1(M^*M)|}$
- for $P_h < 1$, matrix blocks have one 'bad' eigenvalue



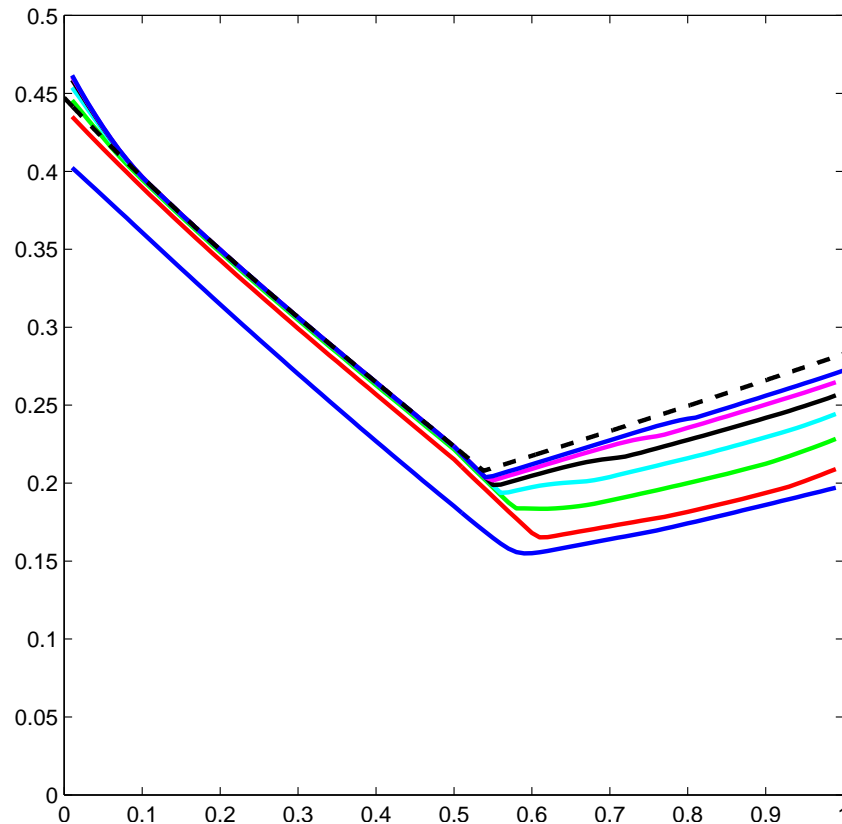
$$\sqrt{|\lambda_1(\mathcal{M}_1^* \mathcal{M}_1)|} \text{ for fixed } P_h = 0.38$$

Alternative Bound?

- artificially 'remove' this eigenvalue: use $\sqrt{|\lambda_2(M^*M)|}$

Alternative Bound?

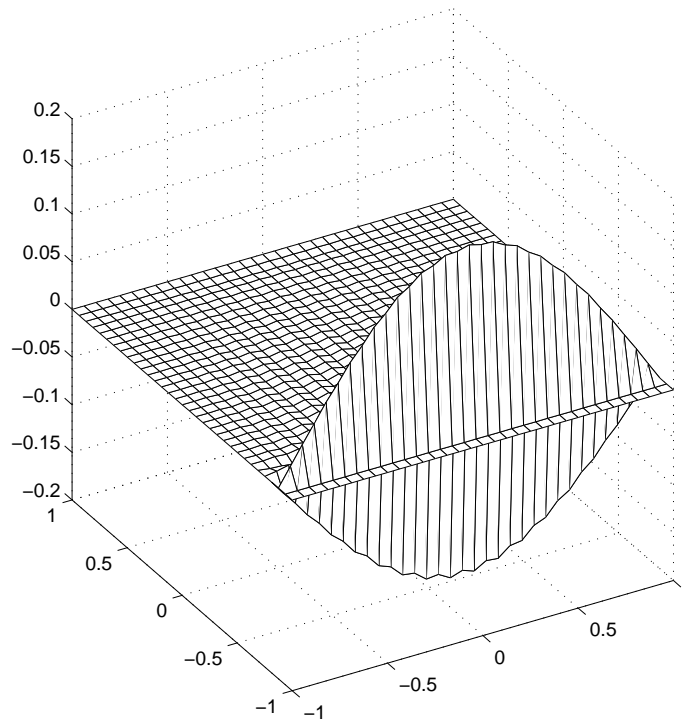
- artificially 'remove' this eigenvalue: use $\sqrt{|\lambda_2(M^*M)|}$



- $P_h < 1$ only
- semiperiodic: $\|M^{per}\|_2$
- Dirichlet: $\sqrt{|\lambda_2(M^*M)|}$

Outlying eigenvalue (1)

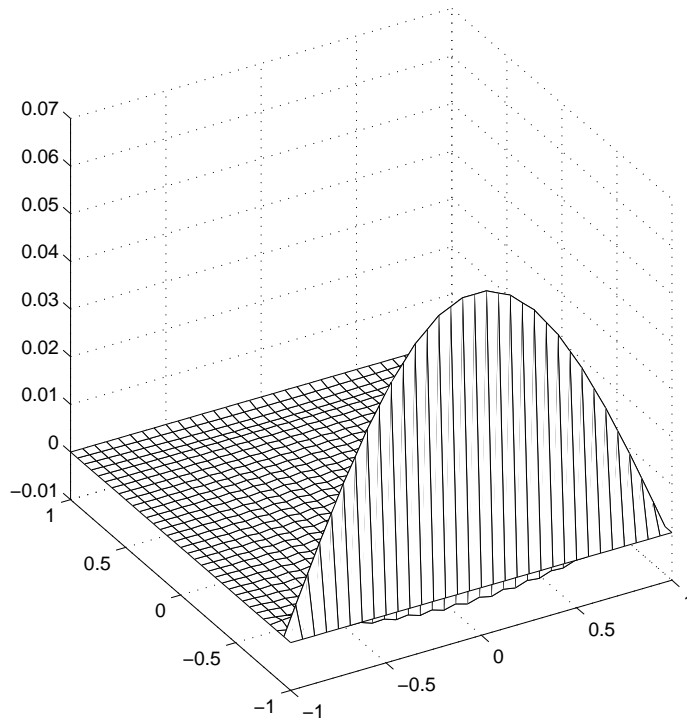
- the **eigenvector** corresponding to the outlying eigenvalue is large only on grid lines very close to the inflow boundary



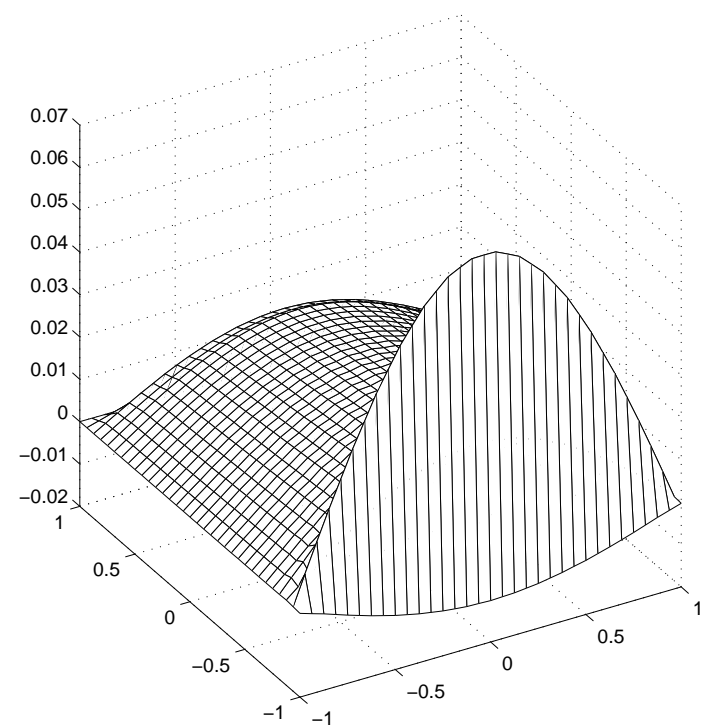
initial error: maximum eigenvector of M^*M

Outlying eigenvalue (2)

- in practice, the effect of this outlying eigenvalue is **transient**



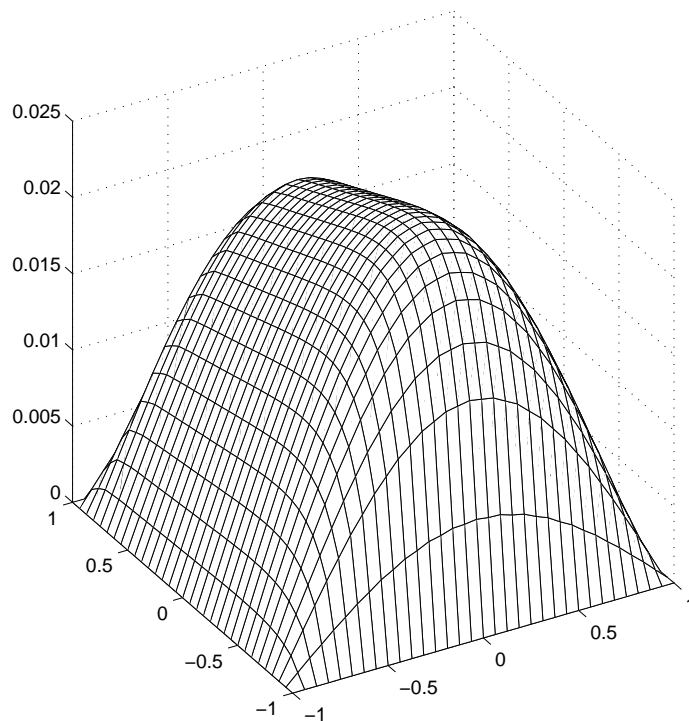
first presmoothing step



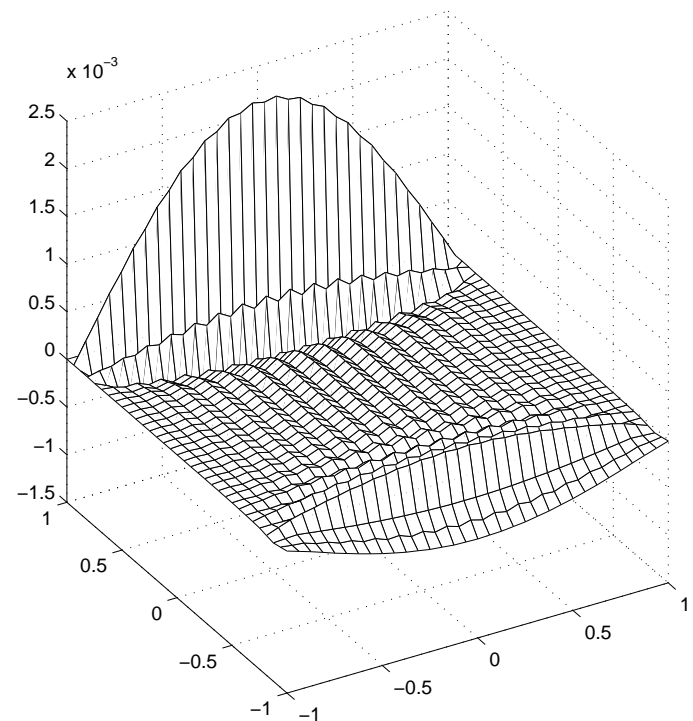
first two-grid iteration

Outlying eigenvalue (3)

- after a few MG iterations, it is **smooth** and so is easily eliminated by coarse grid correction



second presmoothing step



second two-grid iteration

MG Iteration Counts

- these effects do not have an impact on **practical** MG performance

MG Iteration Counts

- these effects do not have an impact on **practical** MG performance

	ϵ										
h	$\frac{1}{2}$	$\frac{1}{4}$	$\frac{1}{8}$	$\frac{1}{16}$	$\frac{1}{32}$	$\frac{1}{64}$	$\frac{1}{128}$	$\frac{1}{256}$	$\frac{1}{512}$	$\frac{1}{1024}$	$\frac{1}{2048}$
$\frac{1}{4}$	5	5	5	5	5	4	4	3	2	2	2
$\frac{1}{8}$	7	7	6	6	5	5	4	4	3	2	2
$\frac{1}{16}$	7	7	7	6	5	5	5	4	4	3	2
$\frac{1}{32}$	7	7	7	7	6	5	5	4	4	3	3
$\frac{1}{64}$	7	7	7	7	6	5	5	4	4	4	3
$\frac{1}{128}$	7	6	6	6	6	6	5	4	4	4	3

$$P_h < 1$$

$$P_h \geq 1$$

- MG-like convergence for any value of P_h

Final Remarks

- Linear algebra gives useful insight into convergence of two-grid iteration.

Final Remarks

- Linear algebra gives useful insight into convergence of two-grid iteration.
- Semiperiodic analysis gives an accurate description of MG behaviour for the full Dirichlet problem.

Final Remarks

- Linear algebra gives useful insight into convergence of two-grid iteration.
- Semiperiodic analysis gives an accurate description of MG behaviour for the full Dirichlet problem.
- Separate approximation and smoothing matrices:
 - semiperiodic analysis for smoothing matrix norm is representative of Dirichlet problem behaviour for all values of P_h ,
 - semiperiodic analysis for approximation matrix norm is representative of Dirichlet problem behaviour for $P_h \geq 1$: for $P_h < 1$, one 'bad' eigenvalue again causes trouble.

Final Remarks

- Linear algebra gives useful insight into convergence of two-grid iteration.
- Semiperiodic analysis gives an accurate description of MG behaviour for the full Dirichlet problem.
- Separate approximation and smoothing matrices:
 - semiperiodic analysis for smoothing matrix norm is representative of Dirichlet problem behaviour for all values of P_h ,
 - semiperiodic analysis for approximation matrix norm is representative of Dirichlet problem behaviour for $P_h \geq 1$: for $P_h < 1$, one 'bad' eigenvalue again causes trouble.
- Replacing the Dirichlet condition by a Neumann condition on the outflow boundary leads to similar computational results.