

# Efficient iterative solvers for director-based models of LCDs

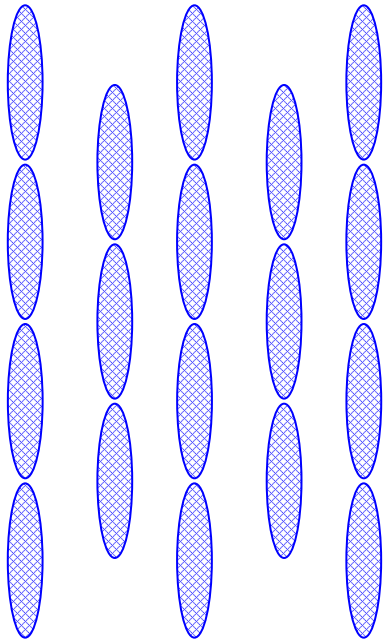
Alison Ramage  
Mathematics and Statistics  
University of Strathclyde  
Glasgow, Scotland



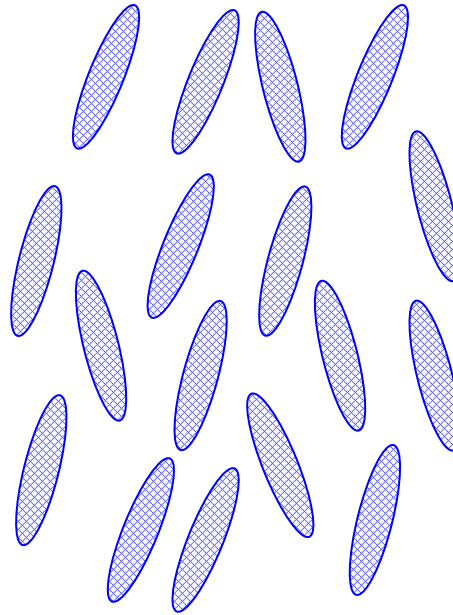
Eugene C. Gartland, Jr.  
Mathematics  
Kent State University  
Ohio, USA

# Liquid Crystals

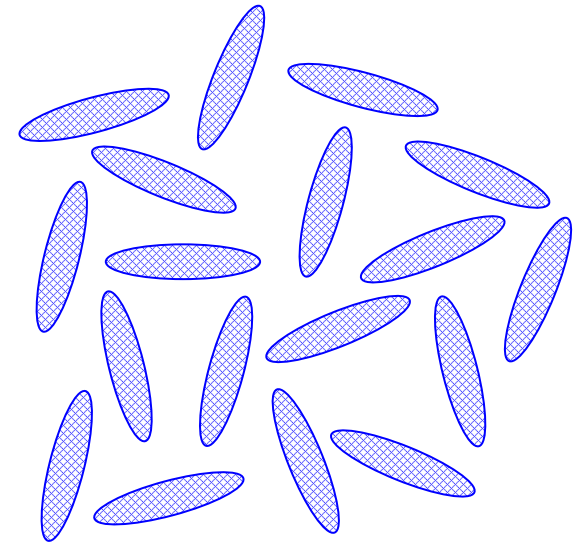
- occur between solid crystal and isotropic liquid states



solid



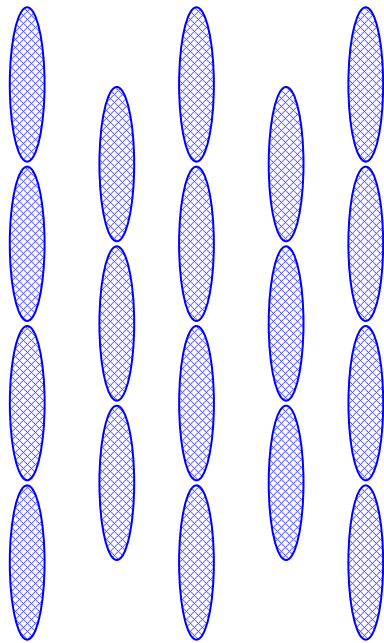
liquid crystal



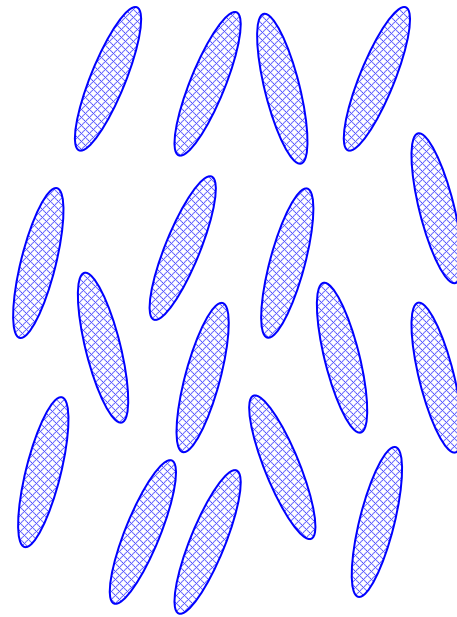
liquid

# Liquid Crystals

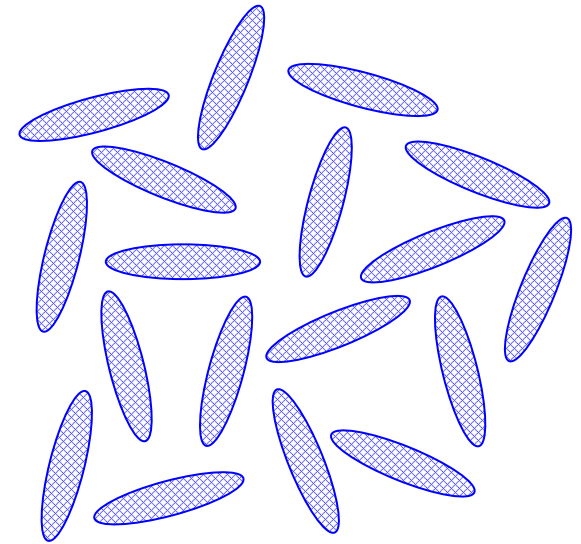
- occur between solid crystal and isotropic liquid states



solid



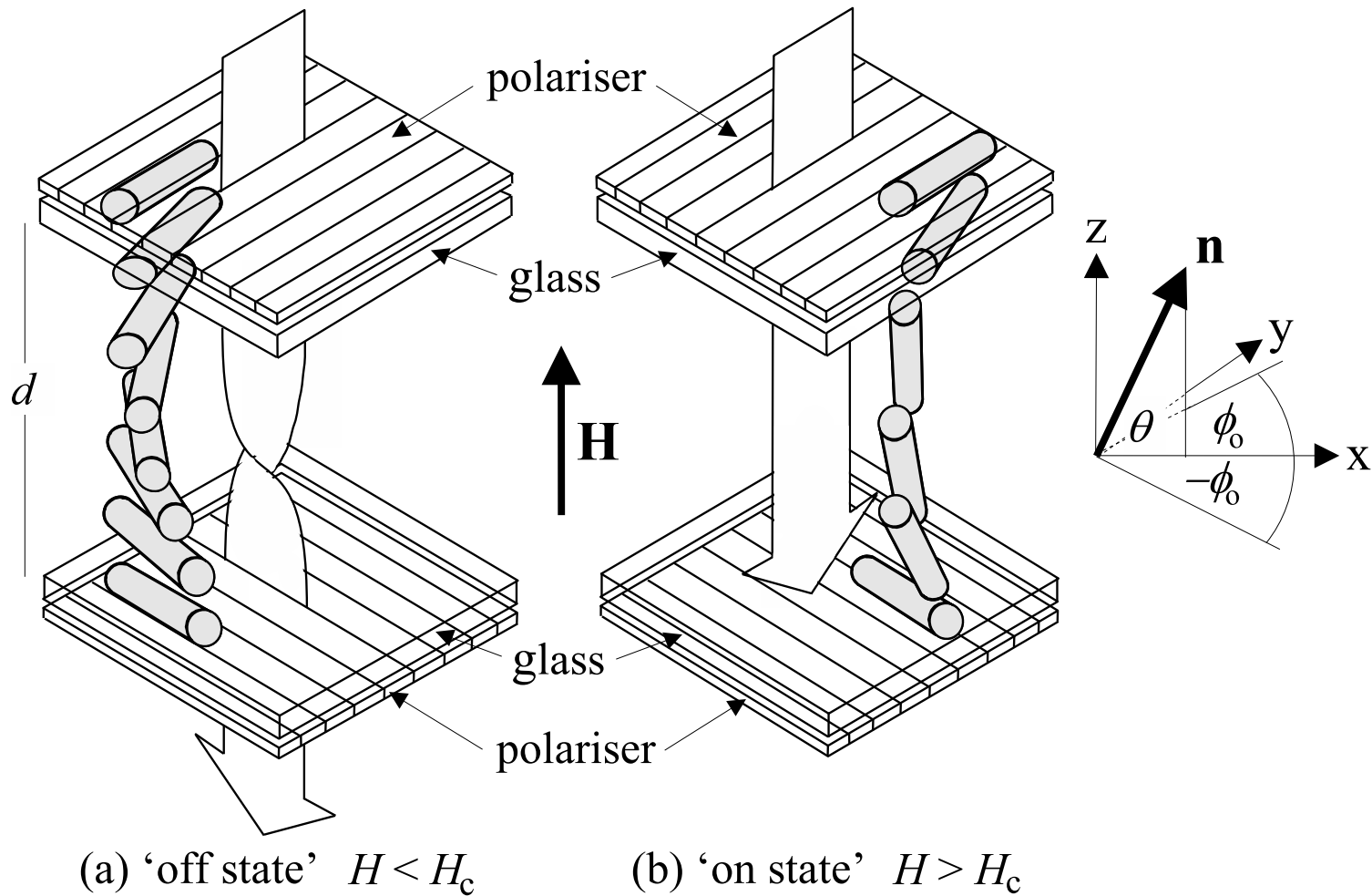
liquid crystal



liquid

- may have different **equilibrium** configurations
- **switch** between stable states by altering applied voltage, magnetic field, boundary conditions, . . .

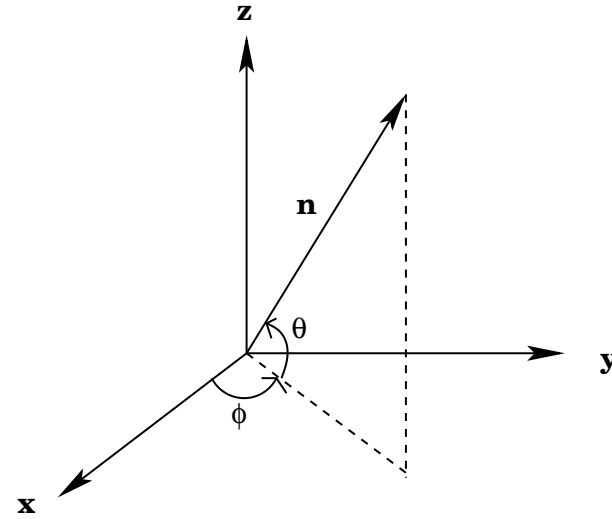
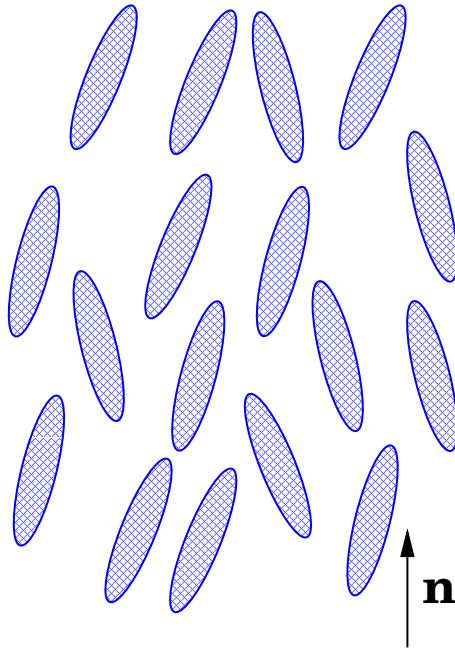
# Liquid Crystal Displays



twisted nematic device

*Static and Dynamic Continuum Theory of Liquid Crystals,*  
*Iain W. Stewart (2004)*

# Modelling: Director-based Models



- **director**: average direction of molecular alignment

unit vector

$$\mathbf{n} = (\cos \theta \cos \phi, \cos \theta \sin \phi, \sin \theta)$$

- **order parameter**: measure of orientational order

$$S = \frac{1}{2} \langle 3 \cos^2 \theta_m - 1 \rangle$$

# Alternative Model: Q-tensor Theory

- tensor order parameter

$$Q = \sqrt{\frac{3}{2}} S (\mathbf{n} \otimes \mathbf{n} - \frac{1}{3} I)$$

- symmetric tensor

$$Q = \begin{bmatrix} q_1 & q_2 & q_3 \\ q_2 & q_4 & q_5 \\ q_3 & q_5 & -q_1 - q_4 \end{bmatrix}$$

$$\text{tr}(Q) = 0, \quad \text{tr}(Q^2) = S^2$$

- five unknowns  $q_1, q_2, q_3, q_4, q_5$

# Finding Equilibrium Configurations

- minimise the **free energy**

$$\mathcal{F} = \int_V F_{bulk}(\theta, \phi, \nabla\theta, \nabla\phi) + \int_{\mathcal{S}} F_{surface}(\theta, \phi) d\mathcal{S}$$

$$F_{bulk} = F_{elastic} + F_{electrostatic}$$

# Finding Equilibrium Configurations

- minimise the **free energy**

$$\mathcal{F} = \int_V F_{bulk}(\theta, \phi, \nabla\theta, \nabla\phi) + \int_{\mathcal{S}} F_{surface}(\theta, \phi) d\mathcal{S}$$

$$F_{bulk} = F_{elastic} + F_{electrostatic}$$

- if fixed boundary conditions are applied, surface energy term can be ignored



# Finding Equilibrium Configurations

- minimise the **free energy**

$$\mathcal{F} = \int_V F_{bulk}(\theta, \phi, \nabla\theta, \nabla\phi) + \int_{\mathcal{S}} F_{surface}(\theta, \phi) d\mathcal{S}$$

$$F_{bulk} = F_{elastic} + F_{electrostatic}$$

- if fixed boundary conditions are applied, surface energy term can be ignored
- solutions with **least** energy are physically relevant

# Elastic Energy

- Frank-Oseen elastic energy

$$\begin{aligned} F_{elastic} &= \frac{1}{2} K_1 (\nabla \cdot \mathbf{n})^2 + \frac{1}{2} K_2 (\mathbf{n} \cdot \nabla \times \mathbf{n})^2 \\ &+ \frac{1}{2} K_3 (\mathbf{n} \times \nabla \times \mathbf{n})^2 \\ &+ \frac{1}{2} (K_2 + K_4) \nabla \cdot [(\mathbf{n} \cdot \nabla) \mathbf{n} - (\nabla \cdot \mathbf{n}) \mathbf{n}] \end{aligned}$$

- Frank elastic constants

$K_1$	splay
$K_2$	twist
$K_3$	bend
$K_2 + K_4$	saddle-splay

# One-Constant Approximation

- set

$$K = K_1 = K_2 = K_3, \quad K_4 = 0$$

# One-Constant Approximation

- set

$$K = K_1 = K_2 = K_3, \quad K_4 = 0$$

- vector identities

$$(\nabla \times \mathbf{n})^2 = (\mathbf{n} \cdot \nabla \times \mathbf{n})^2 + (\mathbf{n} \times \nabla \times \mathbf{n})^2$$

$$\nabla(\mathbf{n} \cdot \mathbf{n}) = 0$$

$$[(\nabla \cdot \mathbf{n})^2 + (\nabla \times \mathbf{n})^2] + \nabla \cdot [(\mathbf{n} \cdot \nabla)\mathbf{n} - (\nabla \cdot \mathbf{n})\mathbf{n}] = \|\nabla \mathbf{n}\|^2$$

# One-Constant Approximation

- set

$$K = K_1 = K_2 = K_3, \quad K_4 = 0$$

- vector identities

$$(\nabla \times \mathbf{n})^2 = (\mathbf{n} \cdot \nabla \times \mathbf{n})^2 + (\mathbf{n} \times \nabla \times \mathbf{n})^2$$

$$\nabla(\mathbf{n} \cdot \mathbf{n}) = 0$$

$$[(\nabla \cdot \mathbf{n})^2 + (\nabla \times \mathbf{n})^2] + \nabla \cdot [(\mathbf{n} \cdot \nabla)\mathbf{n} - (\nabla \cdot \mathbf{n})\mathbf{n}] = \|\nabla\mathbf{n}\|^2$$

- elastic energy  $F_{elastic} = \frac{1}{2}K\|\nabla\mathbf{n}\|^2$

# Electrostatic Energy

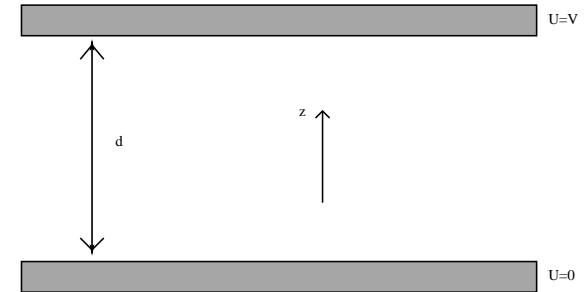
- applied electric field  $\mathbf{E}$  of magnitude  $E$
- electrostatic energy

$$F_{electrostatic} = -\frac{1}{2}\epsilon_0\epsilon_{\perp}E^2 - \frac{1}{2}\epsilon_0\epsilon_a(\mathbf{n} \cdot \mathbf{E})^2$$

- dielectric anisotropy  $\epsilon_a = \epsilon_{\parallel} - \epsilon_{\perp}$
- permittivity of free space  $\epsilon_0$

# Model Problem: Twisted Nematic Device

- two parallel plates distance  $d$  apart



- **strong anchoring** parallel to plate surfaces ( $\mathbf{n}$  fixed)
- rotate one plate through  $\pi/2$  radians
- electric field  $\mathbf{E} = (0, 0, E(z))$ , voltage  $V$

# Equilibrium Equations 1

- equilibrium equations on  $z \in [0, d]$

$$F = \frac{1}{2} \int_0^d \{ K \|\nabla \mathbf{n}\|^2 - \epsilon_0 \epsilon_{\perp} E^2 - \epsilon_0 \epsilon_a (\mathbf{n} \cdot \mathbf{E})^2 \} dz$$

- director  $\mathbf{n} = (u, v, w), \quad |\mathbf{n}| = 1$

- electric potential  $U: \quad E = \frac{dU}{dz}$

- unknowns  $u, v, w, U$



# Equilibrium Equations 2

- nondimensionalise:  $\bar{z} = \frac{z}{d}$ ,  $\bar{U} = \frac{U}{V}$
- nondimensionalised equilibrium equations on  $z \in [0, 1]$

$$F = \frac{1}{2} \int_0^1 [(u_z^2 + v_z^2 + w_z^2) - \alpha^2 \pi^2 (\beta + w^2) U_z^2] dz$$

- dimensionless parameters

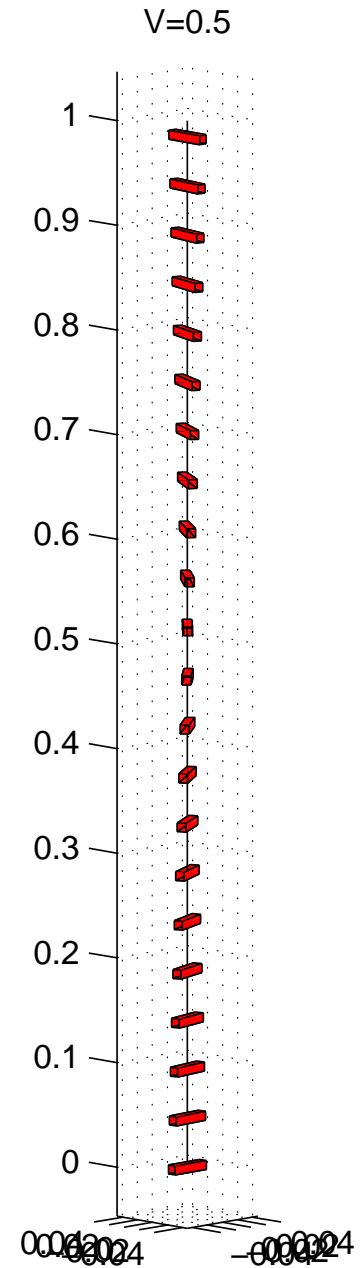
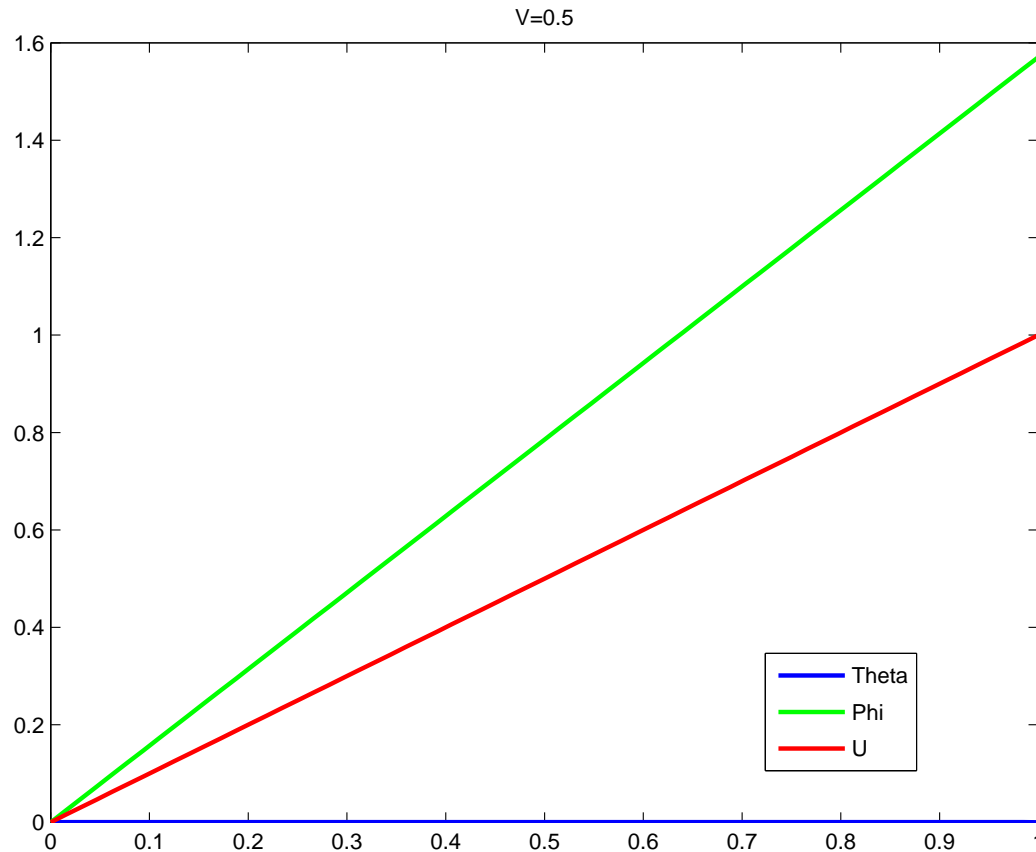
$$\alpha^2 = \frac{\epsilon_0 \epsilon_a V^2}{K \pi^2}, \quad \beta = \frac{\epsilon_{\perp}}{\epsilon_a}$$

- boundary conditions:

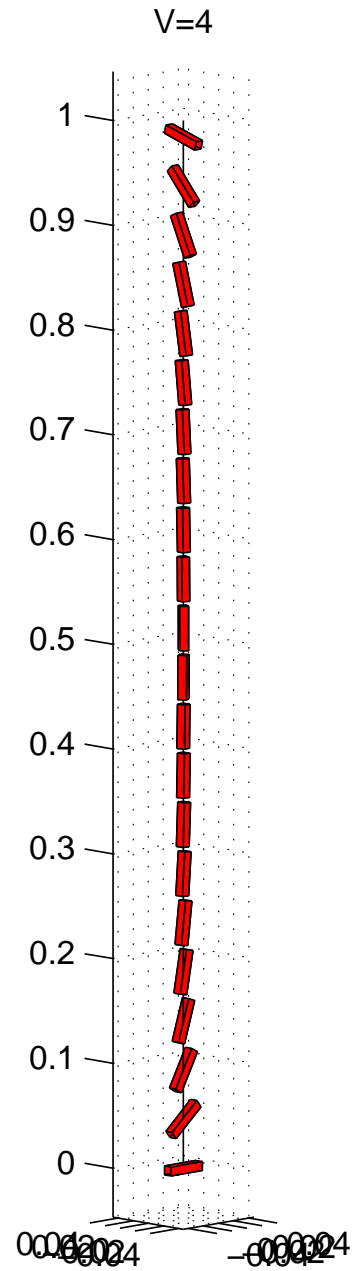
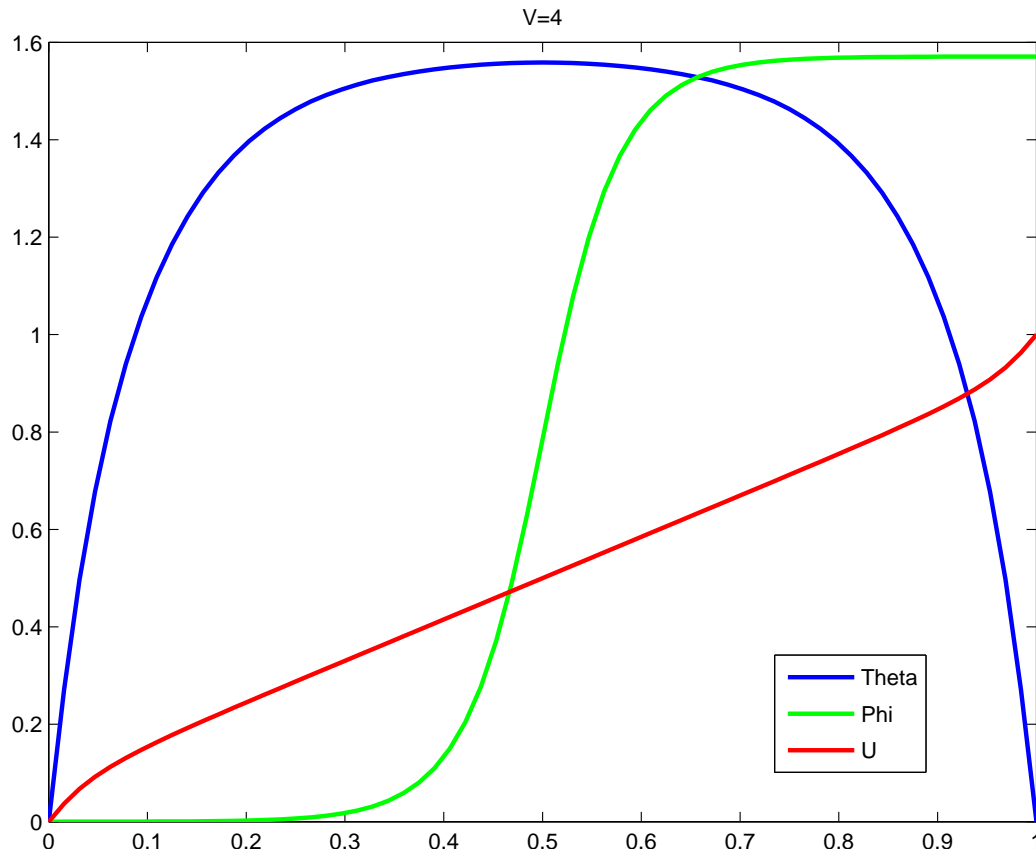
$$\text{at } z = 0: \mathbf{n} = (1, 0, 0), \quad \text{at } z = 1: \mathbf{n} = (0, 1, 0)$$

# Off State

$$\theta(z) \equiv 0, \quad \phi(z) = \frac{\pi}{2}z, \quad U(z) = z$$



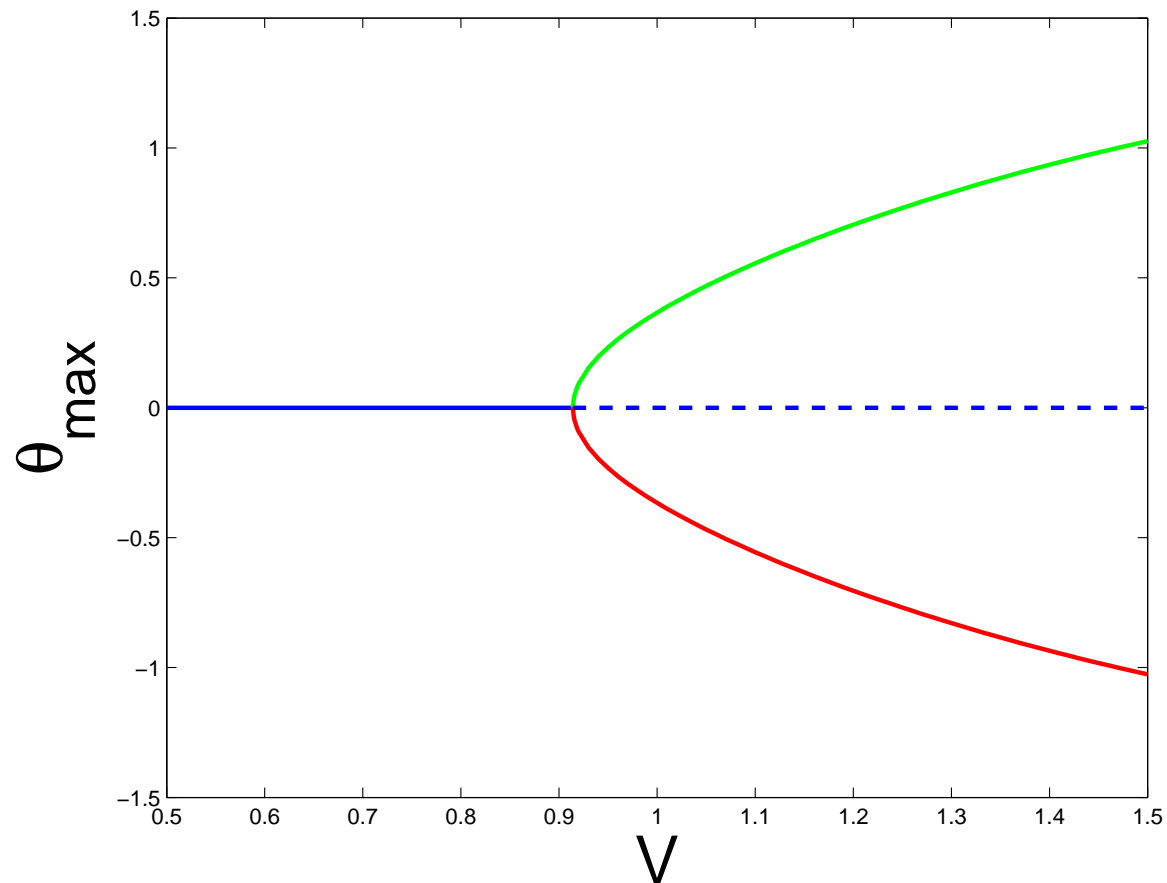
# On State



# Critical Voltage

- switching occurs at

$$V_c = \frac{\pi}{2} \sqrt{\frac{3K}{\epsilon_0 \epsilon_a}}$$



# Discrete Free Energy

- grid of  $N + 1$  points  $z_k$  a distance  $\Delta z$  apart,  $n = N - 1$  unknowns for each variable
- **piecewise linear** approximation, weighted average

$$F \simeq \frac{\Delta z}{2} \sum_{k=0}^{N-1} \left\{ \left[ \frac{u_{k+1} - u_k}{\Delta z} \right]^2 + \left[ \frac{v_{k+1} - v_k}{\Delta z} \right]^2 + \left[ \frac{w_{k+1} - w_k}{\Delta z} \right]^2 - \alpha^2 \pi^2 \left( \beta + \left[ \frac{w_k^2 + w_{k+1}^2}{2} \right] \right) \left[ \frac{U_{k+1} - U_k}{\Delta z} \right]^2 \right\}$$

- equivalent to **mid-point** finite differences, **linear** finite elements

# Constrained Minimisation I

- discrete free energy

$$F \simeq \frac{\Delta z}{2} f(u_1, \dots, u_n, v_1, \dots, v_n, w_1, \dots, w_n, U_1, \dots, U_n)$$

- minimise  $F$  subject to pointwise constraint

$$u_j^2 + v_j^2 + w_j^2 = 1, \quad j = 1, \dots, n$$

- constraints are applied via **Lagrange multipliers**:  
minimise

$$G = \frac{\Delta z}{2} [f - \lambda_1(u_1^2 + v_1^2 + w_1^2 - 1) - \dots - \lambda_n(u_n^2 + v_n^2 + w_n^2 - 1)]$$

# Constrained Minimisation II

- solve  $\nabla G(\mathbf{x}) = \mathbf{0}$  for  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \lambda, \mathbf{U}]$   
 $N + 1$  gridpoints  $\Rightarrow n = N - 1$  unknowns

# Constrained Minimisation II

- solve  $\nabla \mathbf{G}(\mathbf{x}) = \mathbf{0}$  for  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \lambda, \mathbf{U}]$

$N + 1$  gridpoints  $\Rightarrow n = N - 1$  unknowns

- use Newton's method: solve

$$\nabla^2 \mathbf{G}(\mathbf{x}_j) \cdot \delta \mathbf{x}_j = -\nabla \mathbf{G}(\mathbf{x}_j)$$



# Constrained Minimisation II

- solve  $\nabla \mathbf{G}(\mathbf{x}) = \mathbf{0}$  for  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \lambda, \mathbf{U}]$   
 $N + 1$  gridpoints  $\Rightarrow n = N - 1$  unknowns

- use Newton's method: solve

$$\nabla^2 \mathbf{G}(\mathbf{x}_j) \cdot \delta \mathbf{x}_j = -\nabla \mathbf{G}(\mathbf{x}_j)$$

- $5n \times 5n$  coefficient matrix is **Hessian**  $\nabla^2 \mathbf{G}(\mathbf{x})$

$$\nabla^2 \mathbf{G} = \begin{bmatrix} \nabla_{\mathbf{nn}}^2 \mathbf{G} & \nabla_{\mathbf{n}\lambda}^2 \mathbf{G} & \nabla_{\mathbf{nU}}^2 \mathbf{G} \\ \nabla_{\lambda\mathbf{n}}^2 \mathbf{G} & \nabla_{\lambda\lambda}^2 \mathbf{G} & \nabla_{\mathbf{U}\lambda}^2 \mathbf{G} \\ \nabla_{\mathbf{U}\mathbf{n}}^2 \mathbf{G} & \nabla_{\lambda\mathbf{U}}^2 \mathbf{G} & \nabla_{\mathbf{UU}}^2 \mathbf{G} \end{bmatrix}$$

# Hessian Components 1

- matrix notation:  $\nabla_{\mathbf{nn}}^2 \mathbf{G} = \mathbf{A}$

$$\mathbf{A} = \begin{bmatrix} \nabla_{\mathbf{uu}}^2 \mathbf{G} & 0 & 0 \\ 0 & \nabla_{\mathbf{vv}}^2 \mathbf{G} & 0 \\ 0 & 0 & \nabla_{\mathbf{ww}}^2 \mathbf{G} \end{bmatrix} = \begin{bmatrix} A_{uu} & 0 & 0 \\ 0 & A_{vv} & 0 \\ 0 & 0 & A_{ww} \end{bmatrix}$$

- $A_{uu}$ ,  $A_{vv}$  and  $A_{ww}$  are  $n \times n$  **symmetric tridiagonal** blocks

# Hessian Components 1

- matrix notation:  $\nabla_{\mathbf{nn}}^2 \mathbf{G} = \mathbf{A}$

$$\mathbf{A} = \begin{bmatrix} \nabla_{\mathbf{uu}}^2 \mathbf{G} & 0 & 0 \\ 0 & \nabla_{\mathbf{vv}}^2 \mathbf{G} & 0 \\ 0 & 0 & \nabla_{\mathbf{ww}}^2 \mathbf{G} \end{bmatrix} = \begin{bmatrix} A_{uu} & 0 & 0 \\ 0 & A_{vv} & 0 \\ 0 & 0 & A_{ww} \end{bmatrix}$$

- $A_{uu}$ ,  $A_{vv}$  and  $A_{ww}$  are  $n \times n$  **symmetric tridiagonal** blocks

- $A_{uu} = A_{vv} = \frac{1}{\Delta z} \text{tri}(-1, 2 - \Delta z^2 \lambda_j, -1)$

- $A_{ww} = \frac{1}{\Delta z} \text{tri}(-1, 2 - \Delta z^2 \lambda_j - \gamma_j, -1)$

$$\gamma_j = \frac{\alpha^2 \pi^2}{2} [(U_{j+1} - U_j)^2 + (U_j - U_{j-1})^2]$$

# Eigenvalues of $A$

- **off state**: first Newton step, linear  $U$ , constant  $\lambda$

$$\lambda_j = \lambda = \frac{4}{\Delta z^2} \sin^2 \left( \frac{\pi \Delta z}{4} \right)$$

- block matrices are **Toeplitz**

# Eigenvalues of $A$

- **off state**: first Newton step, linear  $U$ , constant  $\lambda$

$$\lambda_j = \lambda = \frac{4}{\Delta z^2} \sin^2 \left( \frac{\pi \Delta z}{4} \right)$$

- block matrices are **Toeplitz**

- $\sigma_{\min}(A_{uu}) = \sigma_{\min}(A_{vv}) \simeq \frac{3\pi^2}{4} \Delta z > 0$

$A_{uu}$  and  $A_{vv}$  are positive definite

- $\sigma_{\min}(A_{ww}) \simeq \left( \frac{3\pi^2}{4} - \alpha^2 \pi^2 \right) \Delta z$

$A_{ww}$  is positive definite iff  $V < V_c$

# Eigenvalues of $A$

- **off state**: first Newton step, linear  $U$ , constant  $\lambda$

$$\lambda_j = \lambda = \frac{4}{\Delta z^2} \sin^2 \left( \frac{\pi \Delta z}{4} \right)$$

- block matrices are **Toeplitz**

- $\sigma_{\min}(A_{uu}) = \sigma_{\min}(A_{vv}) \simeq \frac{3\pi^2}{4} \Delta z > 0$

$A_{uu}$  and  $A_{vv}$  are positive definite

- $\sigma_{\min}(A_{ww}) \simeq \left( \frac{3\pi^2}{4} - \alpha^2 \pi^2 \right) \Delta z$

$A_{ww}$  is positive definite iff  $V < V_c$

- number of negative eigenvalues increases with  $V$

# Hessian Components 2

- matrix notation:  $\nabla_{\mathbf{n}\lambda}^2 \mathbf{G} = B$

- the  $3n \times n$  matrix  $B$  has structure

$$B = -\Delta z \begin{bmatrix} B_u \\ B_v \\ B_w \end{bmatrix}, \quad \begin{aligned} B_u &= \text{diag}(\mathbf{u}) \\ B_v &= \text{diag}(\mathbf{v}) \\ B_w &= \text{diag}(\mathbf{w}) \end{aligned}$$

# Hessian Components 2

- matrix notation:  $\nabla_{\mathbf{n}\lambda}^2 \mathbf{G} = B$

- the  $3n \times n$  matrix  $B$  has structure

$$B = -\Delta z \begin{bmatrix} B_u \\ B_v \\ B_w \end{bmatrix}, \quad \begin{aligned} B_u &= \text{diag}(\mathbf{u}) \\ B_v &= \text{diag}(\mathbf{v}) \\ B_w &= \text{diag}(\mathbf{w}) \end{aligned}$$

- $B^T B = \Delta z^2 I_n$  when constraints are satisfied

- $\text{rank}(B) = \text{rank}(B^T) = \text{rank}(BB^T) = \text{rank}(B^T B) = n$



# Hessian Components 3

- matrix notation:  $\nabla_{\mathbf{U}\mathbf{U}}^2 \mathbf{G} = -\mathbf{C}$
- the  $n \times n$  matrix  $\mathbf{C}$  is **symmetric** and **tridiagonal**

# Hessian Components 3

- matrix notation:  $\nabla_{\mathbf{U}\mathbf{U}}^2 \mathbf{G} = -\mathbf{C}$
- the  $n \times n$  matrix  $\mathbf{C}$  is **symmetric** and **tridiagonal**
- $\mathbf{C} = \frac{1}{\Delta z} \text{tri}(-a_{j-\frac{1}{2}}, a_{j-\frac{1}{2}} + a_{j+\frac{1}{2}}, -a_{j+\frac{1}{2}})$

$$a_{j-\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_{j-1}^2 + w_j^2) \right) > 0$$

$$a_{j+\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_j^2 + w_{j+1}^2) \right) > 0$$

# Hessian Components 3

- matrix notation:  $\nabla_{\mathbf{U}\mathbf{U}}^2 \mathbf{G} = -\mathbf{C}$
- the  $n \times n$  matrix  $\mathbf{C}$  is **symmetric** and **tridiagonal**
- $\mathbf{C} = \frac{1}{\Delta z} \text{tri}(-a_{j-\frac{1}{2}}, a_{j-\frac{1}{2}} + a_{j+\frac{1}{2}}, -a_{j+\frac{1}{2}})$

$$a_{j-\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_{j-1}^2 + w_j^2) \right) > 0$$

$$a_{j+\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_j^2 + w_{j+1}^2) \right) > 0$$

- diagonally dominant with positive real diagonal entries

$\mathbf{C}$  is positive definite

# Hessian Components 4

- matrix notation:  $\nabla_{\mathbf{n}U}^2 \mathbf{G} = D$

$$D = \frac{\alpha^2 \pi^2}{\Delta z} \begin{bmatrix} 0 \\ 0 \\ D_w \end{bmatrix}$$

- the  $n \times n$  matrix  $D_w$  is **tridiagonal**

$$D_w = \text{diag}(\mathbf{w}) \text{tri}(U_j - U_{j-1}, U_{j-1} - 2U_j + U_{j+1}, U_j - U_{j+1})$$

# Hessian Components 4

- matrix notation:  $\nabla_{\mathbf{nU}}^2 \mathbf{G} = D$

$$D = \frac{\alpha^2 \pi^2}{\Delta z} \begin{bmatrix} 0 \\ 0 \\ D_w \end{bmatrix}$$

- the  $n \times n$  matrix  $D_w$  is **tridiagonal**

$$D_w = \text{diag}(\mathbf{w}) \text{tri}(U_j - U_{j-1}, U_{j-1} - 2U_j + U_{j+1}, U_j - U_{j+1})$$

- $D_w$  has **complex** eigenvalues in conjugate pairs and one zero eigenvalue ( $N$  even)
- $\text{rank}(D) = n - 1$

# Full Hessian Structure

$$\nabla^2 \mathbf{G} = \begin{bmatrix} \nabla_{\mathbf{nn}}^2 \mathbf{G} & \nabla_{\mathbf{n}\lambda}^2 \mathbf{G} & \nabla_{\mathbf{nU}}^2 \mathbf{G} \\ \nabla_{\lambda\mathbf{n}}^2 \mathbf{G} & \nabla_{\lambda\lambda}^2 \mathbf{G} & \nabla_{\mathbf{U}\lambda}^2 \mathbf{G} \\ \nabla_{\mathbf{U}\mathbf{n}}^2 \mathbf{G} & \nabla_{\lambda\mathbf{U}}^2 \mathbf{G} & \nabla_{\mathbf{UU}}^2 \mathbf{G} \end{bmatrix}$$

$$\nabla^2 \mathbf{G} = \begin{bmatrix} A & B & D \\ B^T & 0 & 0 \\ D^T & 0 & -C \end{bmatrix}$$

saddle-point problem

# Four Saddle-Point Problems

- for unknown vector ordered as  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \mathbf{U}, \lambda]$

$$H = \left[ \begin{array}{c|cc} A & D & B \\ \hline D^T & -C & 0 \\ B^T & 0 & 0 \end{array} \right]$$

$$H = \left[ \begin{array}{cc|c} A & D & B \\ \hline D^T & -C & 0 \\ B^T & 0 & 0 \end{array} \right]$$

- for unknown vector ordered as  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \lambda, \mathbf{U}]$

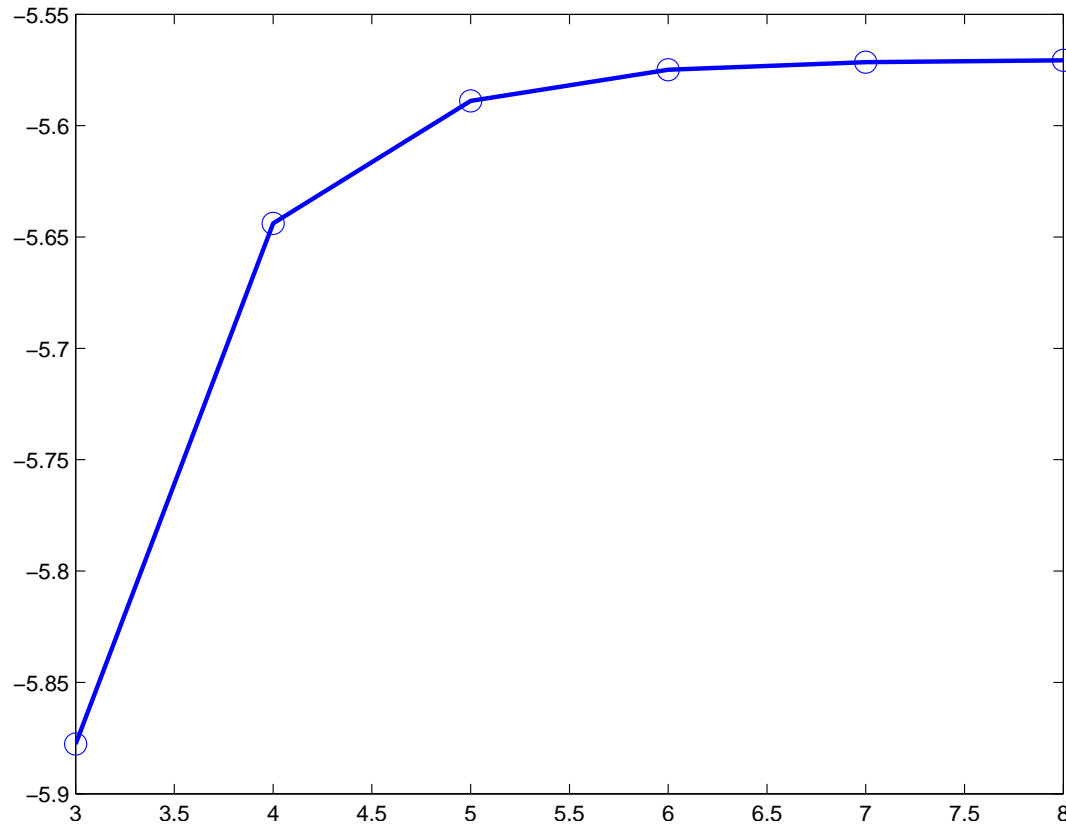
$$H = \left[ \begin{array}{c|cc} A & B & D \\ \hline B^T & 0 & 0 \\ D^T & 0 & -C \end{array} \right]$$

$$H = \left[ \begin{array}{cc|c} A & B & D \\ \hline B^T & 0 & 0 \\ D^T & 0 & -C \end{array} \right]$$

**double** saddle-point structure

# Iterative Solution

- outer iteration: **Newton's method**  $\text{tol}=1e-4$
- inner iteration: **MINRES**  $\text{tol}=1e-4$
- check accuracy by calculating energy of final solution





# MINRES

Paige and Saunders (1975)

Construct iterates  $\mathbf{x}_k = \mathbf{x}_0 + V_k \mathbf{y}_k$  with properties

- $\mathbf{x}_k$  minimises  $\|\mathbf{r}_k\|_2 = \|\mathbf{b} - H\mathbf{x}_k\|_2$
- uses three-term recurrence relation

$$V_k = [\mathbf{v}_1, \mathbf{v}_2, \dots, \mathbf{v}_k]$$

$\mathbf{v}_k$  form an orthonormal basis for Krylov subspace

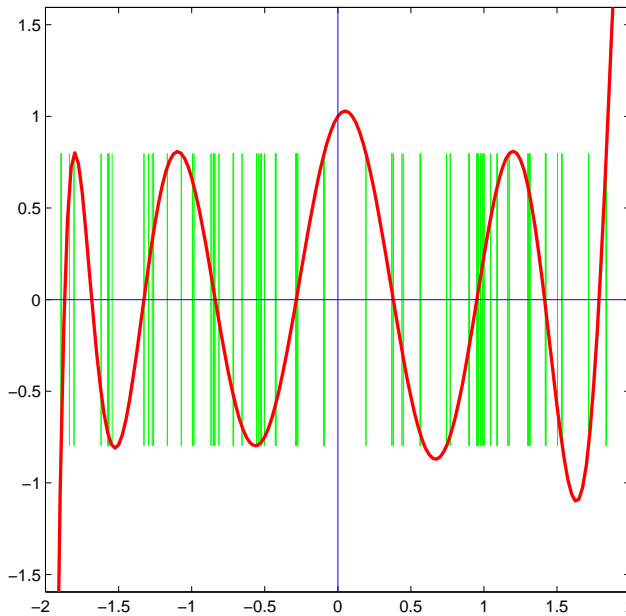
$$\kappa(H, \mathbf{r}_0, k) = \text{span}\{\mathbf{r}_0, H\mathbf{r}_0, \dots, H^{k-1}\mathbf{r}_0\}$$

- use Lanczos method to find  $\mathbf{v}_k$
- solve resulting least squares problem for  $\mathbf{y}_k$  using Givens rotations and QR factorisation

# Convergence of MINRES

- at step  $k$ :

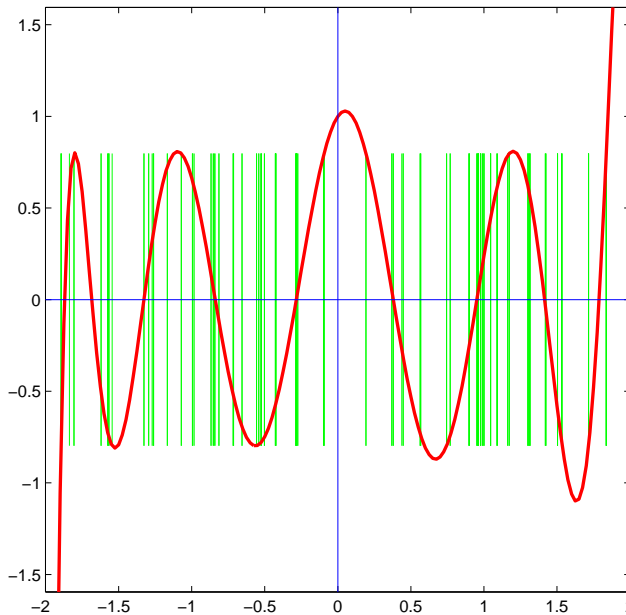
$$\|\mathbf{r}_k\|_2 \leq \min_{p_k \in \Pi_k^1} \max_i |p_k(\lambda_i)| \|\mathbf{r}_0\|_2$$



# Convergence of MINRES

- at step  $k$ :

$$\|\mathbf{r}_k\|_2 \leq \min_{p_k \in \Pi_k^1} \max_i |p_k(\lambda_i)| \|\mathbf{r}_0\|_2$$



- symmetric intervals:  $[-\lambda_{\max}, -\lambda_{\min}] \cup [\lambda_{\min}, \lambda_{\max}]$

$$k \propto \frac{\lambda_{\max}}{\lambda_{\min}}$$

# Minres iterations for full system

$N$	$d$	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
		first step	last step	first step	last step
32	155	226	499	291	691
64	315	728	2,004	1,172	3,571
128	635	2,680	8,528	4,106	17,498
256	1,275	10,253	41,666	15,727	85,784
512	2,555	38,809	194,753	57,499	>200,000
1,024	5,115	150,376	>200,000	>200,000	>200,000

- doubling  $N$  quadruples iteration count

# Nullspace Method I

- Newton system:

$$\begin{bmatrix} A & B & D \\ B^T & 0 & 0 \\ D^T & 0 & -C \end{bmatrix} \begin{bmatrix} \delta \mathbf{n} \\ \delta \lambda \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -\nabla_{\mathbf{n}} G \\ -\nabla_{\lambda} G \\ -\nabla_{\mathbf{U}} G \end{bmatrix}$$

# Nullspace Method I

- Newton system:

$$\begin{bmatrix} A & B & D \\ B^T & 0 & 0 \\ D^T & 0 & -C \end{bmatrix} \begin{bmatrix} \delta \mathbf{n} \\ \delta \lambda \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -\nabla_{\mathbf{n}} G \\ -\nabla_{\lambda} G \\ -\nabla_{\mathbf{U}} G \end{bmatrix}$$

- **Idea:** use information about nullspace of  $B$  to eliminate constraint blocks

# Nullspace Method I

- Newton system:

$$\begin{bmatrix} A & B & D \\ B^T & 0 & 0 \\ D^T & 0 & -C \end{bmatrix} \begin{bmatrix} \delta \mathbf{n} \\ \delta \lambda \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -\nabla_{\mathbf{n}} G \\ -\nabla_{\lambda} G \\ -\nabla_{\mathbf{U}} G \end{bmatrix}$$

- **Idea:** use information about nullspace of  $B$  to eliminate constraint blocks
- use  $Z \in \mathbb{R}^{3n \times 2n}$  whose columns form a basis for the nullspace of  $B^T$

$$B^T Z = Z^T B = 0$$

- $\text{rank}(Z) = 2n$

# Nullspace Method I

- Newton system:

$$\begin{bmatrix} A & B & D \\ B^T & 0 & 0 \\ D^T & 0 & -C \end{bmatrix} \begin{bmatrix} \delta \mathbf{n} \\ \delta \lambda \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -\nabla_{\mathbf{n}} G \\ -\nabla_{\lambda} G \\ -\nabla_{\mathbf{U}} G \end{bmatrix}$$

- **Idea:** use information about nullspace of  $B$  to eliminate constraint blocks
- use  $Z \in \mathbb{R}^{3n \times 2n}$  whose columns form a basis for the nullspace of  $B^T$

$$B^T Z = Z^T B = 0$$

- $\text{rank}(Z) = 2n$
- system size will reduce from  $5n \times 5n$  to  $3n \times 3n$



# Nullspace Method II

$$A\delta\mathbf{n} + B\delta\lambda + D\delta\mathbf{U} = -\nabla_{\mathbf{n}}G \quad (1)$$

$$B^T\delta\mathbf{n} = -\nabla_{\lambda}G \quad (2)$$

$$D^T\delta\mathbf{n} - C\delta\mathbf{U} = -\nabla_{\mathbf{U}}G \quad (3)$$

# Nullspace Method II

$$A\delta\mathbf{n} + B\delta\lambda + D\delta\mathbf{U} = -\nabla_{\mathbf{n}}G \quad (1)$$

$$B^T \delta\mathbf{n} = -\nabla_{\lambda}G \quad (2)$$

$$D^T \delta\mathbf{n} - C\delta\mathbf{U} = -\nabla_{\mathbf{U}}G \quad (3)$$

- write solution of (2) as

$$\delta\mathbf{n} = \widehat{\delta\mathbf{n}} + Z\mathbf{z}$$

- particular solution satisfies  $B^T \widehat{\delta\mathbf{n}} = -\nabla_{\lambda}G$
- $Z\mathbf{z} \in \mathbb{R}^{2n}$  lies in nullspace of  $B^T$

# Nullspace Method II

$$A\delta\mathbf{n} + B\delta\lambda + D\delta\mathbf{U} = -\nabla_{\mathbf{n}}G \quad (1)$$

$$B^T \delta\mathbf{n} = -\nabla_{\lambda}G \quad (2)$$

$$D^T \delta\mathbf{n} - C\delta\mathbf{U} = -\nabla_{\mathbf{U}}G \quad (3)$$

- write solution of (2) as

$$\delta\mathbf{n} = \widehat{\delta\mathbf{n}} + Z\mathbf{z}$$

- particular solution satisfies  $B^T \widehat{\delta\mathbf{n}} = -\nabla_{\lambda}G$
  - $Z\mathbf{z} \in \mathbb{R}^{2n}$  lies in nullspace of  $B^T$
- find  $\widehat{\delta\mathbf{n}}$  via  $\widehat{\delta\mathbf{n}} = -B(B^T B)^{-1} \nabla_{\lambda}G$

# Nullspace Method III

- reduced system:

$$\begin{bmatrix} Z^T A Z & Z^T D \\ D^T Z & -C \end{bmatrix} \begin{bmatrix} \mathbf{z} \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -Z^T (\nabla_{\mathbf{n}} G + A \widehat{\delta \mathbf{n}}) \\ -\nabla_{\mathbf{U}} G - D^T \widehat{\delta \mathbf{n}} \end{bmatrix}$$

# Nullspace Method III

- reduced system:

$$\begin{bmatrix} Z^T A Z & Z^T D \\ D^T Z & -C \end{bmatrix} \begin{bmatrix} \mathbf{z} \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -Z^T (\nabla_{\mathbf{n}} G + A \widehat{\delta \mathbf{n}}) \\ -\nabla_{\mathbf{U}} G - D^T \widehat{\delta \mathbf{n}} \end{bmatrix}$$

- recover full solution from

$$\widehat{\delta \mathbf{n}} = -B(B^T B)^{-1} \nabla_{\lambda} G$$

$$\delta \mathbf{n} = Z \mathbf{z} + \widehat{\delta \mathbf{n}}$$

$$\delta \lambda = (B^T B)^{-1} B^T (-\nabla_{\mathbf{n}} G - A \delta \mathbf{n} - D \delta \mathbf{U})$$

# Nullspace Method III

- reduced system:

$$\begin{bmatrix} Z^T A Z & Z^T D \\ D^T Z & -C \end{bmatrix} \begin{bmatrix} \mathbf{z} \\ \delta \mathbf{U} \end{bmatrix} = \begin{bmatrix} -Z^T (\nabla_{\mathbf{n}} G + A \widehat{\delta \mathbf{n}}) \\ -\nabla_{\mathbf{U}} G - D^T \widehat{\delta \mathbf{n}} \end{bmatrix}$$

- recover full solution from

$$\widehat{\delta \mathbf{n}} = -B(B^T B)^{-1} \nabla_{\lambda} G$$

$$\delta \mathbf{n} = Z \mathbf{z} + \widehat{\delta \mathbf{n}}$$

$$\delta \lambda = (B^T B)^{-1} B^T (-\nabla_{\mathbf{n}} G - A \delta \mathbf{n} - D \delta \mathbf{U})$$

- here  $B^T B$  is **diagonal** so solve is cheap

# Nullspace of $B^T$ I

- permute entries of B:

$$B = -\Delta z \begin{bmatrix} \mathbf{n}_1 & & & \\ & \mathbf{n}_2 & & \\ & & \ddots & \\ & & & \mathbf{n}_n \end{bmatrix}, \quad \mathbf{n}_j = \begin{bmatrix} u_j \\ v_j \\ w_j \end{bmatrix}$$

# Nullspace of $B^T$ I

- permute entries of B:

$$B = -\Delta z \begin{bmatrix} \mathbf{n}_1 & & & \\ & \mathbf{n}_2 & & \\ & & \ddots & \\ & & & \mathbf{n}_n \end{bmatrix}, \quad \mathbf{n}_j = \begin{bmatrix} u_j \\ v_j \\ w_j \end{bmatrix}$$

- eigenvectors of **orthogonal projection**

$$I - \mathbf{n}_j \otimes \mathbf{n}_j = \begin{bmatrix} 1 - u_j^2 & -v_j u_j & -w_j u_j \\ -u_j v_j & 1 - v_j^2 & -w_j v_j \\ -u_j w_j & -v_j w_j & 1 - w_j^2 \end{bmatrix}$$

will be orthogonal to  $\mathbf{n}_j$



# Nullspace of $B^T$ II

- eigenvectors of **orthogonal projection**: e.g.

$$\mathbf{l}_j = \begin{bmatrix} -\frac{v_j}{u_j} \\ 1 \\ 0 \end{bmatrix}, \quad \mathbf{m}_j = \begin{bmatrix} -\frac{w_j}{u_j} \\ 0 \\ 1 \end{bmatrix} \quad (u_j \neq 0)$$

- orthonormalise:

$$\mathbf{l}_j = \frac{1}{\sqrt{u_j^2 + v_j^2}} \begin{bmatrix} -v_j \\ u_j \\ 0 \end{bmatrix}, \quad \mathbf{m}_j = \frac{1}{\sqrt{u_j^2 + v_j^2}} \begin{bmatrix} -u_j w_j \\ -v_j w_j \\ u_j^2 + v_j^2 \end{bmatrix}$$

- at least one of  $u_j, v_j, w_j$  nonzero as  $|\mathbf{n}_j| = 1$

# Nullspace of $B^T$ III

$$Z = \begin{bmatrix} \mathbf{l}_1 & \mathbf{m}_1 & & & & \\ & & \mathbf{l}_2 & \mathbf{m}_2 & & \\ & & & & \ddots & \\ & & & & & \mathbf{l}_n & \mathbf{m}_n \end{bmatrix}$$

# Nullspace of $B^T$ III

$$Z = \begin{bmatrix} \mathbf{l}_1 & \mathbf{m}_1 & & & & \\ & & \mathbf{l}_2 & \mathbf{m}_2 & & \\ & & & & \ddots & \\ & & & & & \mathbf{l}_n & \mathbf{m}_n \end{bmatrix}$$

- consider  $B^T Z \mathbf{p}$  where  $\mathbf{p} = [p_1, q_1, p_2, q_2, \dots, p_n, q_n]^T$ :

$$B^T Z \mathbf{p} = \begin{bmatrix} \mathbf{n}_1^T & & & & \\ & \mathbf{n}_2^T & & & \\ & & \ddots & & \\ & & & & \mathbf{n}_n^T \end{bmatrix} \begin{bmatrix} p_1 \mathbf{l}_1 + q_1 \mathbf{m}_1 \\ p_2 \mathbf{l}_2 + q_2 \mathbf{m}_2 \\ \vdots \\ p_n \mathbf{l}_n + q_n \mathbf{m}_n \end{bmatrix} = 0$$

- columns of  $Z$  form a **basis** for nullspace of  $B^T$

# Minres iterations for reduced system

$N$	$d$	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
		first step	last step	first step	last step
32	93	59	128	90	172
64	189	187	418	285	557
128	381	660	1,456	1,004	2,002
256	765	2,562	5,455	3,650	7,043
512	1,533	9,983	21,393	13,907	26,504
1,024	3,069	41,267	80,778	55,563	81,821
2,048	6,141	171,385	>200,000	>200,000	>200,000

- doubling  $N$  quadruples iteration count

# Preconditioning

Idea: instead of solving  $\mathcal{H}\mathbf{x} = \mathbf{b}$ , solve

$$\mathcal{P}^{-1}\mathcal{H}\mathbf{x} = \mathcal{P}^{-1}\mathbf{b}$$

for some preconditioner  $\mathcal{P}$

Choose  $\mathcal{P}$  so that

- (i) eigenvalues of  $\mathcal{P}^{-1}\mathcal{H}$  are well clustered
- (ii)  $\mathcal{P}\mathbf{u} = \mathbf{r}$  is easily solved

# Preconditioning

Idea: instead of solving  $\mathcal{H}\mathbf{x} = \mathbf{b}$ , solve

$$\mathcal{P}^{-1}\mathcal{H}\mathbf{x} = \mathcal{P}^{-1}\mathbf{b}$$

for some preconditioner  $\mathcal{P}$

Choose  $\mathcal{P}$  so that

- (i) eigenvalues of  $\mathcal{P}^{-1}\mathcal{H}$  are well clustered
- (ii)  $\mathcal{P}\mathbf{u} = \mathbf{r}$  is easily solved

Extreme cases:

- $\mathcal{P} = \mathcal{H}$ : good for (i), bad for (ii)
- $\mathcal{P} = I$ : good for (ii), bad for (i)

# Diagonal scaling

- set  $\mathcal{P} = \text{diag}(\mathcal{H})$
- $\mathcal{P}$  is cheap to invert

$N$	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
	first step	last step	first step	last step
32	55	129	54	122
64	169	408	167	390
128	573	1,469	565	1,423
256	2,144	5,479	2060	5,301
512	8,254	21,196	8,148	20,804
1,024	33,438	85,154	33,849	80,221
2,048	136,015	>200,000	133,605	>200,000

# Ideal Block Preconditioner

- write  $\bar{A} = Z^T A Z$  and  $\bar{D} = Z^T D$ :

$$\mathcal{H} = \begin{bmatrix} \bar{A} & \bar{D} \\ \bar{D}^T & -C \end{bmatrix}$$



# Ideal Block Preconditioner

- write  $\bar{A} = Z^T A Z$  and  $\bar{D} = Z^T D$ :

$$\mathcal{H} = \begin{bmatrix} \bar{A} & \bar{D} \\ \bar{D}^T & -C \end{bmatrix}$$

- block preconditioner:  $\mathcal{P} = \begin{bmatrix} \bar{A} & 0 \\ 0 & C \end{bmatrix}$

# Ideal Block Preconditioner

- write  $\bar{A} = Z^T A Z$  and  $\bar{D} = Z^T D$ :

$$\mathcal{H} = \begin{bmatrix} \bar{A} & \bar{D} \\ \bar{D}^T & -C \end{bmatrix}$$

- block preconditioner:  $\mathcal{P} = \begin{bmatrix} \bar{A} & 0 \\ 0 & C \end{bmatrix}$
- preconditioned matrix:

$$\tilde{\mathcal{H}} = \mathcal{P}^{-1/2} \mathcal{H} \mathcal{P}^{-1/2} = \begin{bmatrix} I & M^T \\ M & -I \end{bmatrix}$$

$$M = C^{-1/2} \bar{D} \bar{A}^{-1/2}$$

# Preconditioned Spectrum

$$\tilde{\mathcal{H}} = \mathcal{P}^{-1/2} \mathcal{H} \mathcal{P}^{-1/2} = \begin{bmatrix} I & M^T \\ M & -I \end{bmatrix}$$

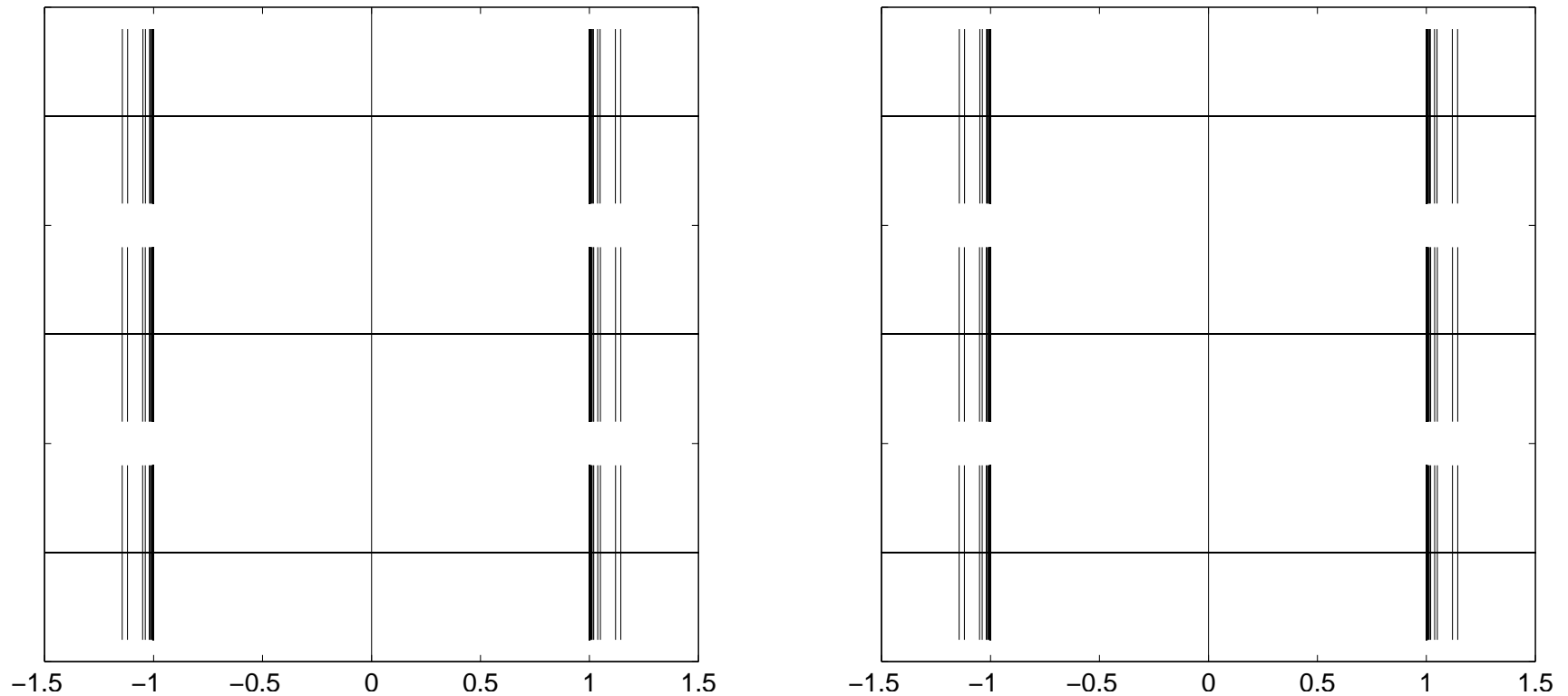
- $M = C^{-1/2} Z^T D (Z^T A Z)^{-1/2}$
- $\text{rank}(M) = n - 1$
- non-zero singular values  $\sigma_k$

# Preconditioned Spectrum

$$\tilde{\mathcal{H}} = \mathcal{P}^{-1/2} \mathcal{H} \mathcal{P}^{-1/2} = \begin{bmatrix} I & M^T \\ M & -I \end{bmatrix}$$

- $M = C^{-1/2} Z^T D (Z^T A Z)^{-1/2}$
- $\text{rank}(M) = n - 1$
- non-zero singular values  $\sigma_k$
- $3n$  eigenvalues of  $\tilde{\mathcal{H}}$  are
  - (i)  $1$  with multiplicity  $n + 1$
  - (ii)  $-1$  with multiplicity  $1$
  - (iii)  $\pm \sqrt{1 + \sigma_k^2}$  for  $k = 1, \dots, n - 1$

# Sample Eigenvalue Plots



eigenvalue plots for  $N = 32, 64, 128$   
first and last Newton iteration

# Estimate of MINRES convergence

- eigenvalues in two symmetric intervals

$$[-\beta, -1] \cup [1, \beta], \quad \beta = \sqrt{1 + \sigma_{\max}^2}$$

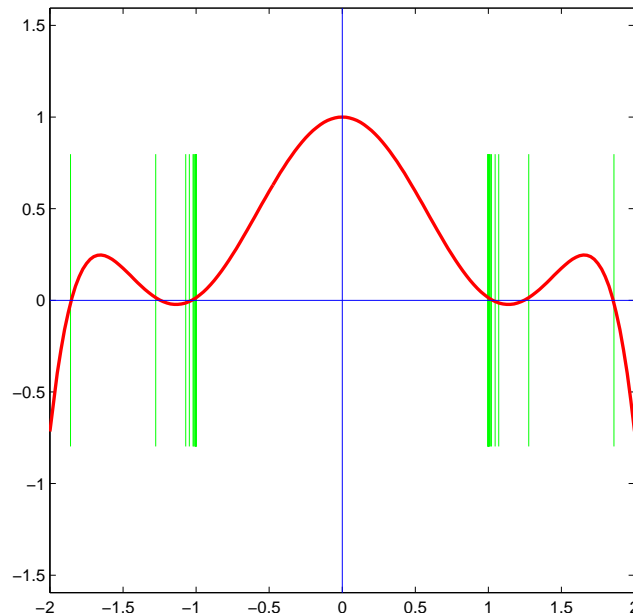
# Estimate of MINRES convergence

- eigenvalues in two symmetric intervals

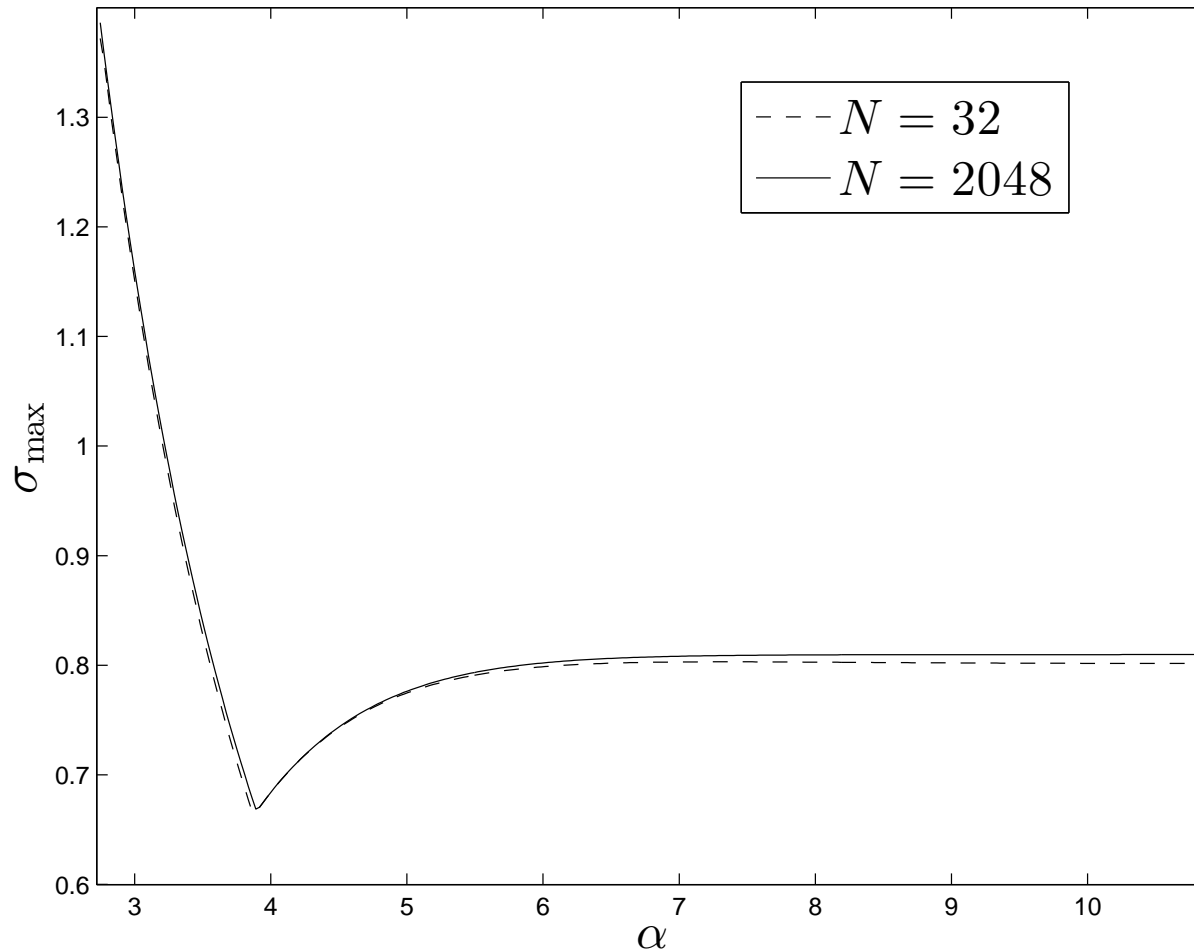
$$[-\beta, -1] \cup [1, \beta], \quad \beta = \sqrt{1 + \sigma_{\max}^2}$$

- to achieve  $\|\mathbf{r}_k\| \leq \epsilon \|\mathbf{r}_0\|$  need

$$k \simeq \frac{1}{2} \sqrt{1 + \sigma_{\max}^2} \ln \left( \frac{2}{\epsilon} \right)$$



# Largest singular value of $M$



- $\sigma_{\max}$  is essentially **independent of  $N$**
- expect a **constant** number of MINRES iterations



# Iteration Counts

$N$	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
	first step	last step	first step	last step
32	4	1	5	7
64	4	1	5	7
128	4	1	5	7
256	4	1	5	7
512	4	1	5	7
1,024	4	1	5	7
2,048	4	1	5	7
4,096	4	1	5	7
8,192	4	1	5	7
16,384	4	1	5	7
32,768	4	1	5	7
65,536	4	1	5	7

# Other options?

- Several other sophisticated saddle point solvers available.
- Many would work here as  $\tilde{A}$  and  $C$  are easy to invert.
- Example: ideal constraint preconditioner

$$\mathcal{P} = \begin{bmatrix} \tilde{A} & \tilde{D} \\ \tilde{D}^T & -C \end{bmatrix}$$

with **P**rojected **P**reconditioned **C**onjugate **G**radient method (Dollar et al. 2006)

- converges in one iteration at each Newton step

# Computing Time

- elapsed time in seconds (tic/toc)

$N$	full direct	reduced direct	ideal block	ideal constraint
1,024	9.95e-02	9.70e-02	3.48e-01	3.08e-01
2,048	1.42e-01	1.36e-01	5.32e-01	8.35e-01
4,096	2.91e-01	2.79e-01	1.05e+00	2.73e+00
8,192	6.02e-01	5.90e-01	2.20e+00	9.74e+00
16,384	1.42e+00	1.29e+00	4.69e+00	3.80e+01
32,768	3.36e+00	2.75e+00	9.70e+00	8.25e+02
65,536	9.27e+00	7.41e+00	2.53e+01	—

# Non-“ideal” versions?

- Block systems can also be solved iteratively.
- Example: use a fixed number of PCG iterations with **AMG** preconditioner (HSL\_MI20).

	1 PCG/AMG iteration				3 PCG/AMG iterations			
	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )		off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
$N$	first	last	first	last	first	last	first	last
32	6	5	7	9	4	1	5	7
128	7	6	7	9	4	1	5	7
512	7	6	8	9	4	1	5	7
2,048	7	6	8	9	4	2	5	7
8,192	7	6	8	9	4	2	5	7

# Non-“ideal” versions?

- Block systems can also be solved iteratively.
- Example: use a fixed number of PCG iterations with **AMG** preconditioner (HSL\_MI20).

	1 PCG/AMG iteration				3 PCG/AMG iterations			
	off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )		off state ( $\alpha = 0.5\alpha_c$ )		on state ( $\alpha = 1.5\alpha_c$ )	
$N$	first	last	first	last	first	last	first	last
32	6	5	7	9	4	1	5	7
128	7	6	7	9	4	1	5	7
512	7	6	8	9	4	1	5	7
2,048	7	6	8	9	4	2	5	7
8,192	7	6	8	9	4	2	5	7

# Summary and other issues

- Nullspace method plus ideal block preconditioner works very well for this simple 1D director model.
- We have also proposed a modified outer iteration (the **Renormalized Newton Method**) with  $\mathbf{n}$  normalised onto the constraint manifold at each iterative step.
- Overall this gives an efficient solution algorithm for repeated solution of liquid crystal director models.

# Summary and other issues

- Nullspace method plus ideal block preconditioner works very well for this simple 1D director model.
- We have also proposed a modified outer iteration (the **Renormalized Newton Method**) with  $\mathbf{n}$  normalised onto the constraint manifold at each iterative step.
- Overall this gives an efficient solution algorithm for repeated solution of liquid crystal director models.
- Nullspace ideas also apply for full 2D and 3D problems.
- Issues remain re how to precondition  $\tilde{A}$  for these more general cases.
- Use of **spectrally equivalent preconditioned iteration** looks promising.

# Summary and other issues

- Nullspace method plus ideal block preconditioner works very well for this simple 1D director model.
- We have also proposed a modified outer iteration (the **Renormalized Newton Method**) with  $\mathbf{n}$  normalised onto the constraint manifold at each iterative step.
- Overall this gives an efficient solution algorithm for repeated solution of liquid crystal director models.
- Nullspace ideas also apply for full 2D and 3D problems.
- Issues remain re how to precondition  $\tilde{A}$  for these more general cases.
- Use of **spectrally equivalent preconditioned iteration** looks promising.

THANKS!