



# Enhancement of performance of wave turbine during stall using passive flow control: First and second law analysis



Ahmed S. Shehata<sup>a, b, \*</sup>, Qing Xiao<sup>a</sup>, Mohamed M. Selim<sup>b</sup>, A.H. Elbatran<sup>b</sup>, Day Alexander<sup>a</sup>

<sup>a</sup> Department of Naval Architecture, Ocean and Marine Engineering, University of Strathclyde, Glasgow G4 0LZ, UK

<sup>b</sup> College of Engineering and Technology, Arab Academy for Science Technology and Maritime Transport, P.O. 1029, AbuQir, Alexandria, Egypt

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## ABSTRACT

Wells turbine is the most common type of self-rectifying air turbine employed by Oscillating Water Column (OWC) wave energy devices due to its technical simplicity, reliability, and design robustness. Because it subjected to early stall, there were many endeavors to improve the energy extraction performance of Wells turbine within the stall regime. Using the multi suction slots as a passive flow control can help obtaining a delayed stall. Two, three and four suction slots were investigated to improve the performance of Wells turbine in the stall regime. In addition the commonly used first law analysis, the present study utilized an entropy generation minimization method to examine the impact of the multi suction slots method on the entropy generation characteristics around the turbine blade. The turbine blade with optimum suction slots number and location was investigated using the oscillating water system based on the real data from the site. To achieve this purpose, two-dimension numerical models for Wells turbine airfoils under sinusoidal wave flow conditions were built and analyze using (ANSYS FLUENT) solver. It is found that the airfoil with three suction slots located at 40%, 55% and 90% from leading edge in chord percentage give the highest torque coefficient by 26.7% before the stall and 51% after the stall.

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## 1. Introduction

Most of fixed-structure OWC systems are located on the shoreline or near the shore. Shoreline devices are characterized by relatively easier maintenance and installation, and they do not require deep water moorings and long underwater electrical cables. The floating OWC devices are slack-moored to the sea bed and so are largely free to oscillate, enhancing the wave energy absorption if the device is properly designed for that purpose [1–6]. The energy conversion from the oscillating air column can be achieved by using a self-rectifying air turbine such as Wells turbine which was invented by A. A. Wells in 1976 [7–12]. The Wells turbine is one of the simplest and probably the most economical turbines for wave energy conversion. It does not require rectifying air valves and can extract power at a low airflow rate, when other turbines would be inefficient. Therefore, it has been extensively researched and developed in many countries. Most self-rectifying air turbines for

wave energy conversion proposed and tested so far are axial-flow machines of two basic types: the Wells turbine and the impulse turbine. The impulse turbine was patented by I. A. Babintsev in 1975 [13]. Its rotor is basically identical to the rotor of a conventional single-stage steam turbine of axial-flow impulse type. Since the turbine is required to be self-rectifying, there are two rows of guide vanes, placed symmetrically on both sides of the rotor, instead of a single row. These two rows of guide vanes are the reflection of each other, with respect to a plane through the rotor disc [14–16]. Therefore, it is more complex and more costly than Wells turbine. The efficiency of Wells turbine is higher than that of the impulse turbine when the flow coefficient is less than the stall point. But after the stall point of Wells turbine, the efficiency of impulse turbine is considerably higher than that of Wells turbine. However, the peak efficiencies are almost the same [17].

There are several factors that influence the design, hence performance, of Wells turbine [3,18,19]. The optimization and improvement of such parameters aim mainly at overcoming the existing disadvantages of the system. The main disadvantage of Wells turbine is the stall condition [12,17,20]. At large flow rates with a large angle of attack, the boundary layers on the blades tend to separate, leading to a drop in torque coefficient and thus the

\* Corresponding author. Department of Naval Architecture, Ocean and Marine Engineering, University of Strathclyde, Glasgow G4 0LZ, UK.

E-mail address: [ahmed.mohamed-ahmed-shehata@strath.ac.uk](mailto:ahmed.mohamed-ahmed-shehata@strath.ac.uk) (A.S. Shehata).

Nomenclature			
A	Total blade area (m <sup>2</sup> )	$S_V$	Viscous entropy generation rate (W/m <sup>2</sup> K)
c	Blade chord (m)	$T_o$	Reservoir temperature (K)
$C_D$	Drag force coefficient	$\bar{u}_i$	Reynolds Averaged velocity component in i direction (m/s)
$C_L$	Lift force coefficient	V	Volume of a computation cell
$C_T$	Torque coefficient	$V_a$	Instantaneous Velocity (m/s)
D	The fluid domain	$V_{am}$	highest speed of axial direction (m/s)
$D_{ss}$	Suction slot diameter (m)	$V_o$	Initial velocity for computation (m/s)
f	cycle frequency (Hz)	$\dot{W}$	The net-work transfer rate
$F_D$	In-line force acting on cylinder per unit length (gf)	$W_{rev}$	Reversible work
G	The filter function	$\eta_F$	The efficiency in first law of thermodynamics
KE	Kinetic Energy (W/K)	$\eta_S$	The second law efficiency
$L_{ss}$	Suction slot location from leading edge in chord percentage %	$\mu$	Viscosity (Kg/ms)
$L_{Rss}$	Reference suction slot location from leading edge in chord percentage %	$\mu_t$	Turbulent viscosity
K	Turbulent kinetic energy	$\rho$	Density (Kg/m <sup>3</sup> )
$\Delta p$	Pressure difference across the turbine (N/m <sup>2</sup> )	$\bar{\theta}$	Flow coefficient
$P_{ss}$	Suction slot pitch distance on x axis	$\omega$	Rotor angular speed (rad/sec)
$\Delta P_{ss}$	Minimum distance between any two suction slot on x axis	$(-\rho \overline{u'_i u'_j})$	Reynolds stress tensor
$R_m$	Mean rotor radius (m)		
$S_{gen}$	local entropy generation rate (W/m <sup>2</sup> K)	<i>List of abbreviations</i>	
$S_G$	Global entropy generation rate (W/K)	CFD	Computational Fluid Dynamics
$S_{ij}$	Mean strain rate	NACA	National Advisory Committee for Aeronautics
$S_t$	Thermal entropy generation rate (W/m <sup>2</sup> K)	OWC	Oscillating Water Column
		2D	Two Dimensional
		3D	Three Dimensional

efficiency. A further increase in the angle of attack, beyond the stall angle, will result in a decrease in lift and also a significant increase in drag. From Refs. [1,12,17] it can be noted that Wells turbine can extract power at low air flow rate, when other turbines would be inefficient [15]. Also, the aerodynamic efficiency increases with the increase of the flow coefficient (angle of attack) up to a certain value, after which it decreases. Thus, most of the past studies aimed to improve the torque coefficient (the turbine output) and improve the turbine behaviour under the stall condition.

The delay of stall onset contributes to improving Wells turbine performance can be achieved by setting guide vanes on the rotor's hub [18,19]. The results of mathematical simulations considering several aerodynamic designs of the Wells turbine are shown in Table 1 [21]. Different guide vanes designs were compared and investigated analytically [22], taking into account the turbine starting characteristics and efficiency in irregular wave conditions. Table 2 shows the best two designs, but for the total performance, G15N11S40 is recommended. R7N08N65 has rectangular blades, a solidity of 0.7, normal blades, 8 blades, and tip gab of 0.65 mm. The G15N11S40 has 1.5 solidity, 11 blades, and the axial spacing between rotor and guide vane equal to 40 mm.

The effects of unsteady flow conditions on the performance of a monoplane Wells turbine without guide vanes during a field experiment on a OWC device are described in Refs. [23,24]. The torque coefficient shows a hysteretic mechanism characterized by a counter-clock-wise loop that appears with high frequency

oscillations. A dynamic stall phenomenon appears with oscillations of very large amplitude, independently from the frequency. A computational model has been used in Ref. [20] to study the performance and aerodynamics of the turbine, quantitatively and qualitatively. In addition, it is used to study the flow coefficient, turbine stalls and the appropriate inlet velocity profile. It is found from the computed results that the wakes behind the turbine blades (NACA0021) merge rigorously in the portion of Radius ratio = 0.45:1.0, which leads the turbine to stall. The tip gab leakage flow is considerably higher in the trailing edge portion. However, as the flow coefficient increases, leakage flow region advances towards the leading edge, causing a large mass flow of air to leak through the gap. A comparison between numerical and experimental investigations is conducted in Ref. [25] for studying uniform tip gab ratio. Regarding turbine efficiency, it is found that the peak efficiency of the turbine decreases, and shifts towards a higher flow coefficient as the tip gab to chord length ratio increases, while the stall margin becomes wider.

The entropy generation, due to viscous dissipation, was investigated by Refs. [26,27] for Wells turbine airfoil sections. The efficiency for four different airfoils in compression cycle is higher than suction cycle at 2° angle of attack. But when the angle of attack increases, the efficiency for suction cycle increases also more than the compression one. This study suggested that a possible existence of critical Reynolds number at which viscous irreversibilities takes minimum values. Exergy analysis [28] is performed using the

**Table 1**  
A summary of the performance data for the different turbines.

Airfoil	NACA0015		HSIM 15-262123-1576	
	Without guide vanes	With guide vanes	Without guide vanes	With guide vanes
Max. Efficiency	58%	71%	55%	67%

**Table 2**  
Maximum efficiency with minimum torque and starting torque coefficients.

Turbine profile	Maximum efficiency	Minimum torque coefficient	Starting torque coefficient
R7N08N65 G15N11S40	44.8%	0.41	0.49
R7N08F65 G15N21S31	47.4%	0.08	0.33

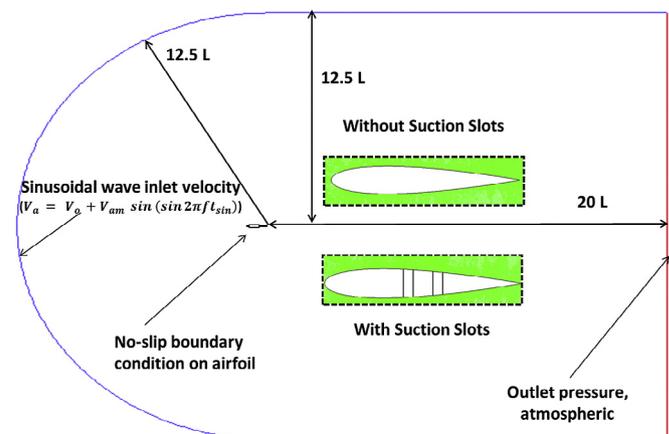
numerical simulation for steady state biplane Wells turbines [29] where the upstream turbine has a design point second law efficiency higher than the downstream turbine second law efficiency by 21.6%. The total entropy generation, due to viscous dissipation, for suggested design for Wells turbine with variable chord was compared with Wells turbine with constant chord in Ref. [30]. The detailed results demonstrate 26.02% average decrease in total entropy generation throughout the full operating range. Recent studies proved that entropy generation within various wave energy extracting systems plays a significant role in determining the overall efficiency of the system [31,32]. Thus, researchers must take into account the entropy analysis for the Wells turbine while investigating the wave energy extracting performance. Since, it has shown very promising result in many applications, such as wind turbine in Refs. [33–38] and gas turbine in Refs. [39–41].

One of the most popular methods that have been used to decrease flow separation around the aerofoil section and delay the stall is the flow control method [42–47]. Different studies have been conducted on flow control techniques. The first scientist who employed boundary layer suction on a cylindrical surface to delay boundary layer separation was Prandtl in 1904 [48]. The earliest known experimental works on boundary layer suction for airfoil and wings were conducted in the late 1930 and the 1940 [49–51]. The idea of passive suction is to use a passive porous surface [52] [53] to mitigate the local pressure gradients and obviate separation to reduce drag. Huang et al. [54] studied the suction and blowing flow control techniques on a NACA0012 airfoil. It can be concluded that perpendicular suction at the leading edge increased lift coefficient more than other suction situations. The tangential blowing at downstream locations was found to lead to the maximum increase in the lift coefficient value. The study in Ref. [55] provide an excellent review of the various periodic excitation methods, mainly steady suction and blowing. This review gives a detailed discussion of the mechanism and also the recent developments in the field. Previous reviews that provide a detailed discussion of the subject include [56–60]. The CFD method has been used to investigate boundary layer control [61–63] such as the effects of blowing and suction jets on the aerodynamic performance of airfoils.

One of the best locations to apply the OWC system with Wells turbine is the northern coast of Egypt [64]. Where, the most energetic coast of the Southern Mediterranean Basin is the Egyptian coast, lying between the Nile Delta and the Libyan borders with a potential of above 3.35 kW/m wave power in summer and 6.8 kW/m in winter [65,66] and the wave energy of about 36003 kWh/m. Wave energy potential of the Eastern Mediterranean Sea Basin was assessed and characterized. Wave fields obtained from 3rd generation spectral wave model for years 1994–2009 by using wind data from European Center for Medium-Range Weather Forecasts (ECMWF) were used in order to calculate the wave powers. Wave model was calibrated using the wave measurements conducted at three different stations. Wave model simulated the wave characteristics such as significant wave heights and mean wave periods with high accuracy. Wave power atlas was generated based on 15-year time-averaged wave data. Also wave power roses and distribution tables in means of periods and heights for different regions were presented [67]. The most energetic sea states have significant wave heights between 1 and 4 m and wave energy

periods between 4 and 8 s. The regions with increased wave energy potential are mainly the western and southern coastlines of Cyprus Island, the sea area of Lebanon and Israel, as well as the coastline of Egypt, especially around Alexandria. The significant differences between the sea in Egypt and other seas are that the sea wave in Egypt is relatively low but also stable. Hence, the potential wave energy can be revealed and exploited. Otherwise, sea states with the wave heights greater than 5 m are not very important for the annual energy [68] as they contribute little to the annual energy. But they need to be considered in design and selection of wave energy converters. Although Eastern Mediterranean and Aegean Sea Basins can be regarded as low potential areas in terms of wave energy, future developments in the wave energy converter technologies can make gathering energy economically viable.

Current researchers are only investigating the aerofoil with the passive flow control methods, such as the suction and blowing slot under steady (non-sinusoidal) flow. Therefore, the objective of the present work is to investigate the passive flow control methods (e.g. the suction and blowing slot) which affect the entropy generation behaviour under sinusoidal flow. While improving the generated torque coefficient on the aerofoil section with the attached slot, it is equally important to accurately model the effect on the entropy generation and the second law efficiency. Therefore, it is essential that to investigate and define the optimum parameters (e.g. slots number, locations, and angle) for single and multi-slots based on the first and second law analysis. Then, the turbine blade with optimum suction slots number and location was investigated using the oscillating water system based on the real data from the northern coast of Egypt. Where, the present work recommends Wells turbine as a suitable choice for the Egyptian coasts due to its simple and efficient operation under low input air flow. Furthermore, in view of the previous research in the section of the connections with other turbines, it can be concluded that the efficiency of Wells turbine is higher than that of the other turbines when the flow coefficient is less than the stall point. But after the stall point of Wells turbine, the efficiency of the others turbine (such as impulse turbine) is considerably higher than that of Wells turbine. In other words, Wells turbine can extract power



**Fig. 1.** Dimensions of whole computational domain and location of airfoil.

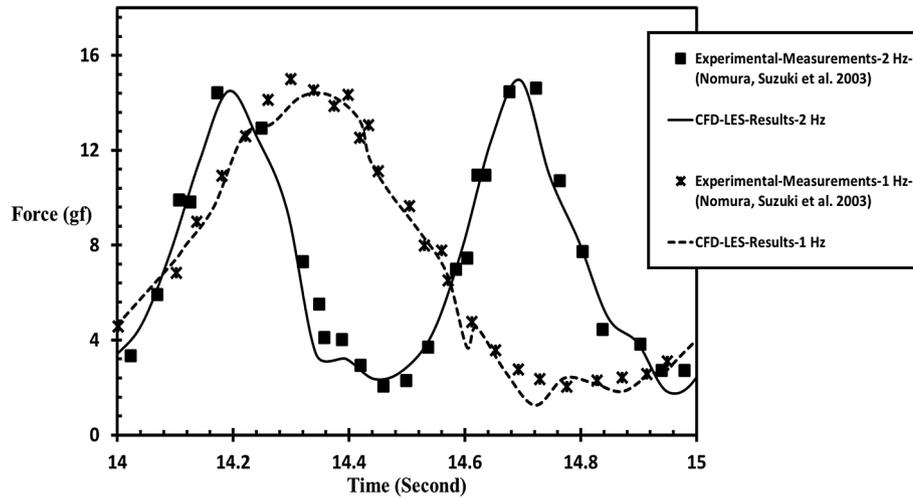


Fig. 2. Measured unsteady in-line force  $F_D$  from reference (Nomura, Suzuki et al., 2003), (angle of attack =  $0^\circ$ ) and  $F_D$  calculated from the present CFD with frequencies 2 Hz and 1 Hz.

Table 3  
The error percentage between measured  $F_D$  from reference (Nomura, Suzuki et al., 2003) and calculated  $F_D$  from CFD under unsteady flow with sinusoidal inlet velocity.

	Time (Second)												
Frequency 2 Hz													
$F_D$ (gf)	14.02	14.1	14.12	14.2	14.3	14.34	14.4	14.5	14.6	14.7	14.8	14.9	15
Experimental	3.4	7.7	9.9	14.4	12.9	3.4	4.11	2.3	7.5	14.7	10.7	3.9	2.7
CFD	3.8	7.8	9.8	14.5	12.5	3.5	3.4	2.7	7.8	14.9	10.9	3.8	2.4
Error %	11	1	-1	1	-4	1	-17	17	4	1	2	-2	-11
Frequency 1 Hz													
Experimental	4.5	6.9	12.6	14.1	14.3	13	10.2	7.8	4.7	2.8	2.3	2.5	3
CFD	4.6	7.2	12.6	13.1	14.3	13.2	10.3	8.6	4.5	2.7	2.2	2.6	3.3
Error %	2	4	0	-7	0	1	1	10	-4	1	-4	4	10

at a low air flow rate, when other turbines would be inefficient. To the best of this author’s knowledge, to date no study exists which define the optimum location and number for multi-slots attached to the aerofoil under sinusoidal flow inlet velocity based on the first and second law analysis. Furthermore, to date no specific unsteady CFD study of the multi-slots effect with sinusoidal flow on the entropy generation rate has been performed for Wells turbine.

2. Mathematical formulations and numerical methodology

In order to solve the governing equations of Large Eddy Simulation (LES), the time-dependent Navier-Stokes terms are filtered. This filtering process aims to eliminate eddies that have scale smaller than the filter width or computational mesh spacing. Therefore, the resulting filtered equations govern the dynamics of large eddies. A filtered variable (denoted by an over-bar) could be expressed as [69]:

$$\bar{\phi}(x) = \int_D \phi(x')G(x, x')dx' \tag{1}$$

where  $D$  represents the fluid domain, and  $G$  represents the filter function that determines the scale of the resolved eddies. The unresolved part of a quantity  $\phi$  is defined by:

$$\phi' = \phi - \bar{\phi} \tag{2}$$

And the filtered fluctuations are not zero:

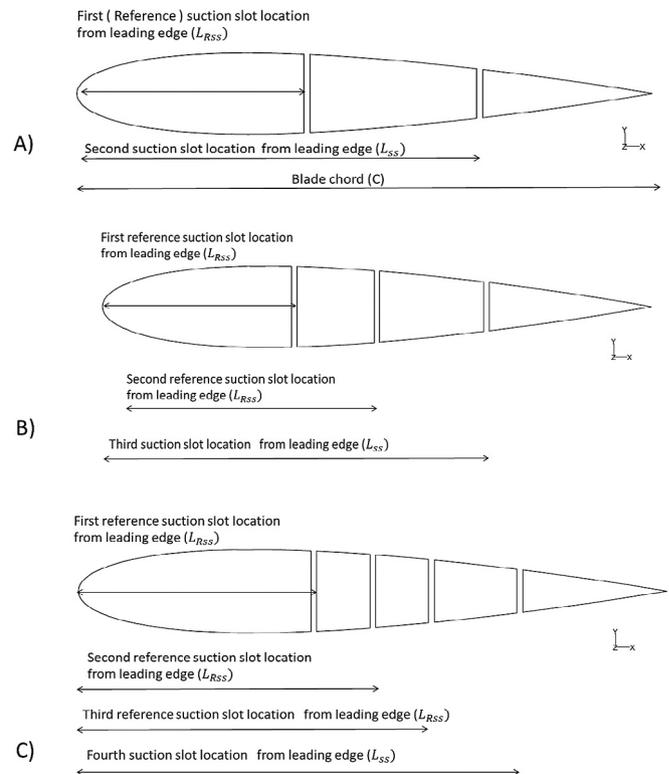


Fig. 3. Aerofoil diagram with two, three and four suction slots A) Two suction slots B) Three suction slots C) four suction slots.

**Table 4**  
Two suction slots with  $P_{ss}$  and  $\Delta P_{ss}$  equal to 0.05.

Reference suction slot	Location of second suction slot															Max. value %	
	20%	25%	30%	35%	40%	45%	50%	55%	60%	65%	70%	75%	80%	85%	90%		
R 20%	-	0.121	0.083	0.12	0.102	0.13	0.126	0.123	0.117	0.11	0.101	0.102	0.094	0.1	0.087	56	
R 25%	0.121	-	0.126	0.124	0.121	0.12	0.123	0.146	0.122	0.114	0.117	0.125	0.108	0.11	0.115	76	
R 30%	0.083	0.126	-	0.125	0.123	0.14	0.118	0.119	0.127	0.118	0.113	0.121	0.117	0.11	0.1	65	
R 35%	0.12	0.124	0.125	-	0.133	0.15	0.125	0.122	0.134	0.125	0.138	0.119	0.116	0.12	0.107	74	
R 40%	0.102	0.121	0.123	0.133	-	0.15	0.144	0.148	0.136	0.133	0.127	0.137	0.134	0.14	0.142	81	
R 45%	0.129	0.123	0.137	0.145	0.151	-	0.131	0.131	0.132	0.137	0.123	0.137	0.128	0.12	0.106	81	
R 50%	0.126	0.123	0.118	0.126	0.144	0.13	-	0.125	0.138	0.142	0.137	0.127	0.138	0.13	0.134	74	
R 55%	0.123	0.146	0.119	0.122	0.148	0.13	0.125	-	0.114	0.122	0.104	0.121	0.131	0.13	0.126	78	
R 60%	0.117	0.122	0.127	0.134	0.136	0.13	0.138	0.114	-	0.108	0.13	0.136	0.128	0.11	0.115	66	
R 65%	0.11	0.114	0.118	0.125	0.133	0.14	0.142	0.122	0.108	-	0.131	0.111	0.131	0.12	0.123	71	
R 70%	0.101	0.117	0.113	0.138	0.127	0.12	0.137	0.104	0.13	0.131	-	0.116	0.117	0.11	0.113	66	
R 75%	0.102	0.125	0.121	0.119	0.137	0.14	0.127	0.121	0.136	0.111	0.116	-	0.118	0.12	0.086	65	
R 80%	0.094	0.108	0.117	0.116	0.134	0.13	0.138	0.131	0.128	0.131	0.117	0.118	-	0.12	0.1	66	
R 85%	0.099	0.111	0.113	0.117	0.135	0.12	0.13	0.129	0.11	0.117	0.113	0.117	0.115	-	0.102	62	
R 90%	0.087	0.115	0.1	0.107	0.142	0.11	0.134	0.126	0.115	0.123	0.113	0.086	0.1	0.10	-	71	
																	Maximum value

$$\overline{\phi'} \neq 0 \tag{3}$$

In FLUENT, the filtering operation is implicitly provided in the finite-volume discretization [1]:

$$\overline{\phi}(x) = \frac{1}{V} \int_V \phi(x') dx', \quad x' \in V \tag{4}$$

where  $V$  is the volume of the control volume. The filter function,  $G(x, x')$ , is expressed as

$$G(x, x') = \begin{cases} 1/V & \text{for } x' \in V \\ 0 & \text{otherwise} \end{cases} \tag{5}$$

While this LES model is treating incompressible flows, not necessarily constant-density flows are assumed. Filtering the incompressible Navier-Stokes terms, one obtains [70].

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho \overline{u}_i}{\partial x_i} = 0 \tag{6}$$

$$\frac{\partial}{\partial t}(\rho \overline{u}_i) + \frac{\partial}{\partial x_j}(\rho \overline{u}_i \overline{u}_j) = \frac{\partial}{\partial x_j} \left( \mu \frac{\partial \overline{u}_i}{\partial x_j} \right) - \frac{\partial \overline{p}}{\partial x_i} - \frac{\partial \tau_{ij}}{\partial x_j} \tag{7}$$

where  $\tau_{ij}$  is the sub-grid-scale stress obtained as

$$\tau_{ij} = \rho \overline{u}_i \overline{u}_j - \rho \overline{u}_i \overline{u}_j \tag{8}$$

Since these sub-grid-scale stresses are unknown, to obtain them an extra modeling step is required. The majority of the eddy viscosity models are of the following form [71]:

$$\tau_{ij} - \frac{1}{3} \tau_{kk} \sigma_{ij} = -2\mu_t \overline{S}_{ij} \tag{9}$$

where  $\overline{S}_{ij}$  is the rate-of-strain tensor which obtained by:

$$\overline{S}_{ij} = \frac{1}{2} \left( \frac{\partial \overline{u}_i}{\partial x_j} + \frac{\partial \overline{u}_j}{\partial x_i} \right) \tag{10}$$

and  $\mu_t$  is the sub-grid-scale turbulent viscosity. The most basic of sub-grid-scale models for “Smagorinsky-Lilly model” was proposed by Smagorinsky [72,73] and further developed by Lilly [74]. In, the eddy viscosity is modeled by Smagorinsky-Lilly model as:

$$\mu_t = \rho L_s^2 |\overline{S}| \tag{11}$$

where  $L_s$  is the mixing length for sub-grid-scale models and  $|\overline{S}| = \sqrt{2\overline{S}_{ij}\overline{S}_{ij}}$ . The  $L_s$  is calculated as follows:

$$L_s = \min(kd, C_s V^{1/3}) \tag{12}$$

where  $C_s$  is the Smagorinsky constant,  $k = 0.42$ ,  $d$  is the distance to the closest wall, and  $V$  is the volume of the computational cell. Lilly derived a value of 0.23 for  $C_s$  from homogeneous isotropic turbulence. However, this value was found to cause excessive damping of large-scale fluctuations in the presence of mean shear or in transitional flows. A dynamic SGS model was not particularly necessary in the LES models due to the turbulence flow at all domain, therefore,  $C_s = 0.1$  has been found to yield the best results for a wide range of flows [1,75,76].

For the first law of thermodynamics, the lift and drag coefficient  $C_L$  and  $C_D$  could be calculated from the post processing software. A single value for the torque coefficient for each angle of attack was calculated using average values for lift and drag coefficients. Thus, the torque coefficient can be calculated as [77]:

$$C_T = (C_L \sin \alpha - C_D \cos \alpha) \tag{13}$$

The flow coefficient,  $\bar{\phi}$ , relating tangential and axial velocities of the rotor is defined as

$$\bar{\phi} = \frac{V_a}{\omega * R_m} \tag{14}$$

where the angle of attack,  $\alpha$ , equals to

$$\alpha = \tan^{-1} \frac{V_a}{\omega R_m} \tag{15}$$

and the torque as:

$$Torque = \frac{1}{2} \rho (V_a^2 + (\omega R_m)^2) AR_m C_T \tag{16}$$

The first law of thermodynamics efficiency ( $\eta_F$ ) is defined as:

$$\eta_F = \frac{Torque * \omega}{\Delta P * Q} \tag{17}$$

According to the second law of thermodynamics, the net-work transfer rate  $\dot{W}$  as [78]:

$$\dot{W} = W_{rev} - T_0 S_{gen} \tag{18}$$

Which, it has been known in engineering as the Gouy–Stodola theorem [79].

It is possible to express the irreversible entropy generation in terms of the derivatives of local flow quantities in the absence of phase changes and chemical reactions. The two dissipative mechanisms in viscous flow are the strain-originated dissipation and the thermal dissipation which correspond to a viscous and a thermal entropy generation respectively [80]. Thus, it can be expressed as,

$$S_{gen} = S_V + S_{th} \tag{19}$$

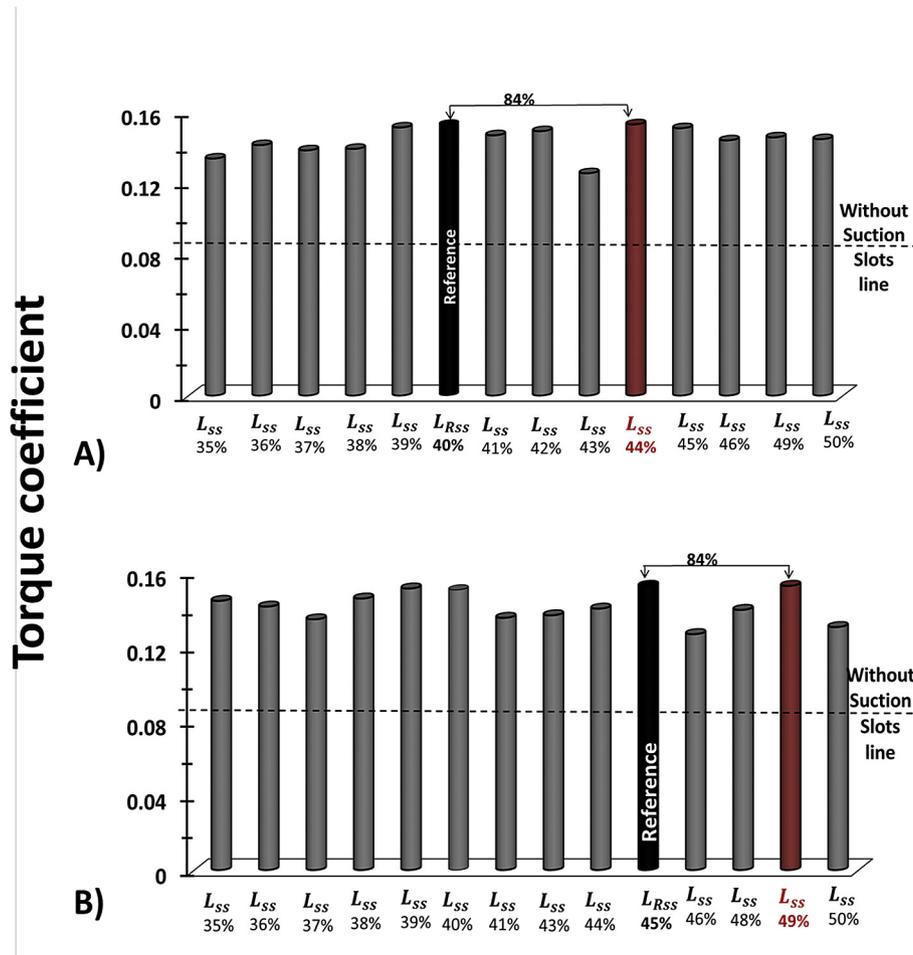


Fig. 4. The effect of tow suction slots on the torque coefficient at the stall angle ( $13.6^\circ$ ) with  $P_{ss}$  and  $\Delta P_{ss}$  equal to 0.01 A)  $L_{RSS} = 40\%$  and  $44\%$ . B)  $L_{RSS} = 45\%$  and  $49\%$ .

**Table 5**

The value of improvement in torque coefficient for the optimum locations for two suction slots at different angles with non-oscillating velocity.

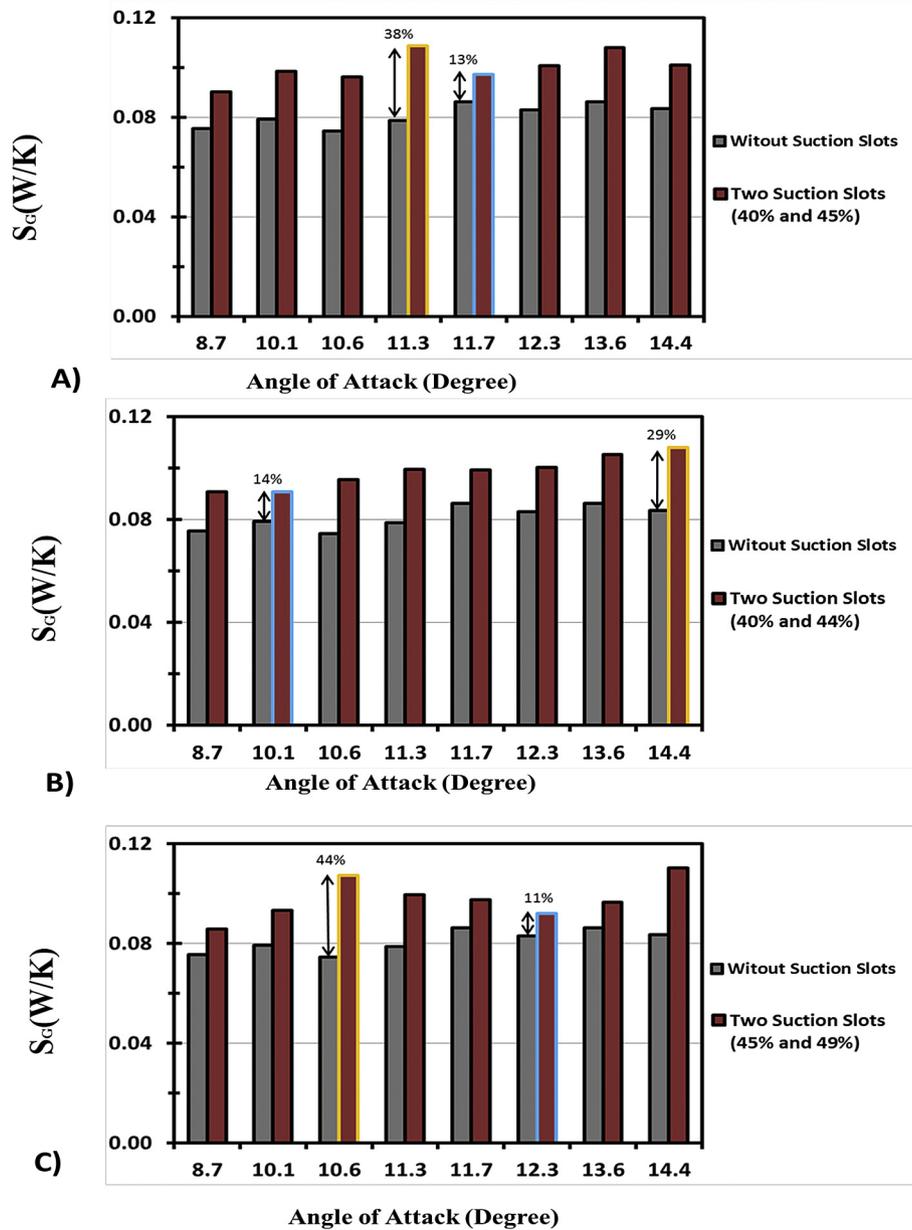
Torque Coefficient	Angle of attack (°)							
	8.7	10.1	10.6	11.3	11.7	12.3	13.6	14.4
Without Suction Slot	0.0509	0.0669	0.0726	0.0793	0.0856	0.0910	0.0830	0.0676
Two Suction Slots at $L_{ss} = 40\%$ and $L_{ss} = 45\%$	0.0654	0.0904	0.0973	0.1122	0.1210	0.1300	0.1505	0.1420
Improvement %	28	35	34	42	41	43	81	110
Two Suction Slots at $L_{ss} = 40\%$ and $L_{ss} = 44\%$	0.0659	0.0848	0.0980	0.1087	0.1175	0.1238	0.1526	0.1428
Improvement %	29	27	35	37	37	36	84	111
Two Suction Slots at $L_{ss} = 45\%$ and $L_{ss} = 49\%$	0.0627	0.0923	0.1011	0.1095	0.1194	0.1297	0.1526	0.1444
Improvement %	23	38	39	38	39	43	84	114

The case in hand represents an incompressible isothermal flow, where the thermal dissipation term vanishes. Therefore, the local viscous irreversibilities can be expressed as [26,27]:

$$S_V = \frac{\mu}{T_0} \varphi \quad (20)$$

where  $\varphi$  is the viscous dissipation term, that is expressed in two dimensional Cartesian coordinates as [80]:

$$\varphi = 2 \left[ \left( \frac{\partial u}{\partial x} \right)^2 + \left( \frac{\partial v}{\partial y} \right)^2 \right] + \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right)^2 \quad (21)$$



**Fig. 5.**  $S_c$  for the optimum locations for two suction slots with non-oscillating velocity A) ( $L_{ss}$  40% and 45%) B) ( $L_{ss}$  40% and 44%). C) ( $L_{ss}$  45% and 49%).

**Table 6**  
Three suction slots with  $P_{ss}$  and  $\Delta P_{ss}$  equal to 0.05.

Reference suction slot	Location of third suction slot															Max. value
	20%	25%	30%	35%	40%	45%	50%	55%	60%	65%	70%	75%	80%	85%	90%	
R 25% and R 55%	0.084	—	0.155	0.149	0.132	0.149	0.100	—	0.136	0.130	0.145	0.130	0.113	0.147	0.144	86%
R 35% and R 45%	0.130	0.131	0.104	—	0.143	—	0.137	0.142	0.140	0.133	0.127	0.142	0.130	0.136	0.138	72%
R 40% and R 45%	0.109	0.144	0.133	0.143	—	—	0.151	0.152	0.155	0.141	0.147	0.151	0.137	0.152	0.126	86%
R 40% and R 50%	0.121	0.104	0.139	0.129	—	0.151	—	0.137	0.146	0.151	0.146	0.142	0.140	0.137	0.144	82%
R 40% and R 55%	0.119	0.132	0.129	0.137	—	0.152	0.137	—	0.147	0.130	0.147	0.159	0.157	0.130	0.161	94%
R 40% and R 90%	0.083	0.135	0.135	0.142	—	0.126	0.144	0.161	0.140	0.137	0.138	0.129	0.140	0.129	—	94%
R 50% and R 65%	0.130	0.146	0.130	0.126	0.151	0.139	—	0.129	0.133	—	0.127	0.133	0.119	0.123	0.130	82%

Maximum value

Equations (20) and (21) were used to create the custom field function file, which is used to calculate the local entropy form the FLUENT software. Then, the global entropy generation rate is hence expressed as:

$$S_G = \int \int_{xy} S_V dy dx \quad (22)$$

Which is also calculated from the FLUENT software by integrating the global value, Equation (23) defines the exergy value as:

$$Exergy = KE + S_G \quad (23)$$

Finally the second law efficiency is defined as [33]:

$$\eta_s = \frac{KE}{Exergy} \quad (24)$$

where  $KE = \frac{1}{2}V^2$ .

From the above equations, it can be concluded that the torque coefficient is related to the first law efficiency, where the increase in torque coefficient leads to an increase in the first law efficiency.

Moreover, the global entropy generation rate is related to the second law efficiency, where the decrease in the global entropy generation rate leads to an increase in the second law efficiency.

### 3. CFD approach

The CFD verification and validation result for such models used in this work is presented. Also, this section is contains a description of the turbulence models used in stall condition, an expression of the discretization methods employed and the boundary conditions for this work.

#### 3.1. Computational model and boundary conditions

2D numerical simulations model for NACA0015 airfoil were conducted and their results were validated against experimental measurements. Both unsteady flow with non-oscillating velocity, as well as, unsteady flow with sinusoidal inlet velocity were investigated. A Cartesian structured mesh was generated using (GAMBIT V-2.4.6) software for discretizing the computational domain. The application matches the Green-Gauss cell based evaluation method for the gradient terms used in the solver (ANSYS FLUENT V-14.5).

**Table 7**  
Three suction slots with  $P_{ss}$  and  $\Delta P_{ss}$  equal to 0.01.

Reference suction slot	Location of third suction slot											Max. value
	95%	94%	93%	92%	91%	90%	89%	88%	87%	86%	85%	
R 40% and R 55%	0.1577	0.1387	0.1369	0.1510	0.1536	0.1610	0.1543	0.1556	0.1542	0.138	0.130	94%
Reference suction slot	60%	59%	58%	57%	56%	55%	54%	53%	52%	51%	50%	Max. value
R 40% and R 90%	0.1397	0.1492	0.1437	0.1421	0.1414	0.1610	0.1425	0.1476	0.1423	0.124	0.144	94%
Reference suction slot	45%	44%	43%	42%	41%	40%	39%	38%	37%	36%	35%	Max. value
R 55% and R 90%	0.1318	0.1401	0.1433	0.1246	0.1386	0.1610	0.1553	0.1256	0.1306	0.125	0.125	94%

Maximum value

For these cases it was found that the second order upwind [81] interpolation scheme yields approximately similar results to those yielded by third order MUSCL scheme in the present situation. Moreover, sometimes the third order MUSCL scheme produces high oscillatory residual during the solution. Therefore, the second order upwind [81] interpolation scheme was used in this study. The Quad-Pave meshing scheme (Structured Grid) was used in this study. It was also found that the solution reaches convergence when the scaled residuals approaches  $1 \times 10^{-5}$ . At this residuals limit, the flow field variables preserve constant values with the application of consecutive iterations.

The axial flow entering Wells turbine was modeled as: 1) a non-oscillating velocity and 2) a sinusoidal wave in this simulation. Therefore, inlet boundary conditions were time dependent. In order to apply the inlet sinusoidal wave boundary condition, inlet velocity with periodic function was generated based on Equation (25). This periodic function was introduced as a user defined function in FLUENT. According to the literature, this equation is used for modeling a sinusoidal wave inlet flow with a non-zero mean velocity over various objects [82,83].

$$V_a = V_o + V_{am} \sin(\sin 2 \pi f t_{sin}) \tag{25}$$

where, the  $V_o$  and  $V_{am}$  are equal to 0.04 and 2.88 m/s respectively,  $t_{sin}$  is the sinusoidal wave time period, and  $f$  is the wave frequency. Based on the conducted literature survey,  $t_{sin}$  for a single period was taken as 6 s ( $f$  equal to 0.167 Hz) in this simulation [75,76,84,85]. In order to satisfy CFL (Courant Friedrichs Lewy) condition equal to 1, the time step was set as 0.000296721 s [86]. For the Egyptian coasts boundary condition, the most energetic sea states have wave energy periods between 4 and 8 s. Therefore, the time period  $t_{sin}$  was taken as 4, 6, and 8 s ( $f$  equal to 0.25, 0.167 and 1.25 Hz) and is set as one period in this simulation, considering the real data from the Egyptian coasts. The sinusoidal wave condition create various Reynolds number up to  $2 \times 10^4$  according to reference [87]. Fig. 1

shows the dimensions of the computational domain and the position of the airfoil.

### 3.2. Grid-independent solution

Several computational grids were tested to ensure a grid-independent solution. The structured grids (Quad-Pave meshing scheme) used in unsteady flow with non-oscillating velocity. The pressure coefficient distribution on the upper and lower surfaces of the NACA0012 aerofoil was computed by the four grids ranging from 112603 up to 446889 cells. Grid 4 (with 446889 cells and  $1 \times 10^{-6}$  for the first cell with growth rate equal to 1.01) required more time than grid 3, yielding similar results for pressure coefficient values. Therefore, grid 3 (with 312951 cells and  $1 \times 10^{-5}$  for the first cell with growth rate equal to 1.012) was chosen to conduct the analysis presented hereafter, more details about grid sensitivity results can be found in Refs. [26,64].

### 3.3. Validation of the CFD model

According to literature [70,88–95], Large Eddy Simulation model excels in predicting flows over airfoil in stall condition. Although LES is a 3D model by definition, there have been numerous successful attempts to use it in 2D applications [6,64]. In this study, two sets of experimental data were used to validate the numerical model from references. First experimental data [87,96,97] was used to simulate and validate the results at stall condition. The first validation case involved the investigation of Wells turbine prototype with the following parameters:

- NACA0015 blade profile
- Hub radius = 101 mm;
- Tip radius = 155 mm;
- Chord length = 74 mm;
- Number of blades = 7;

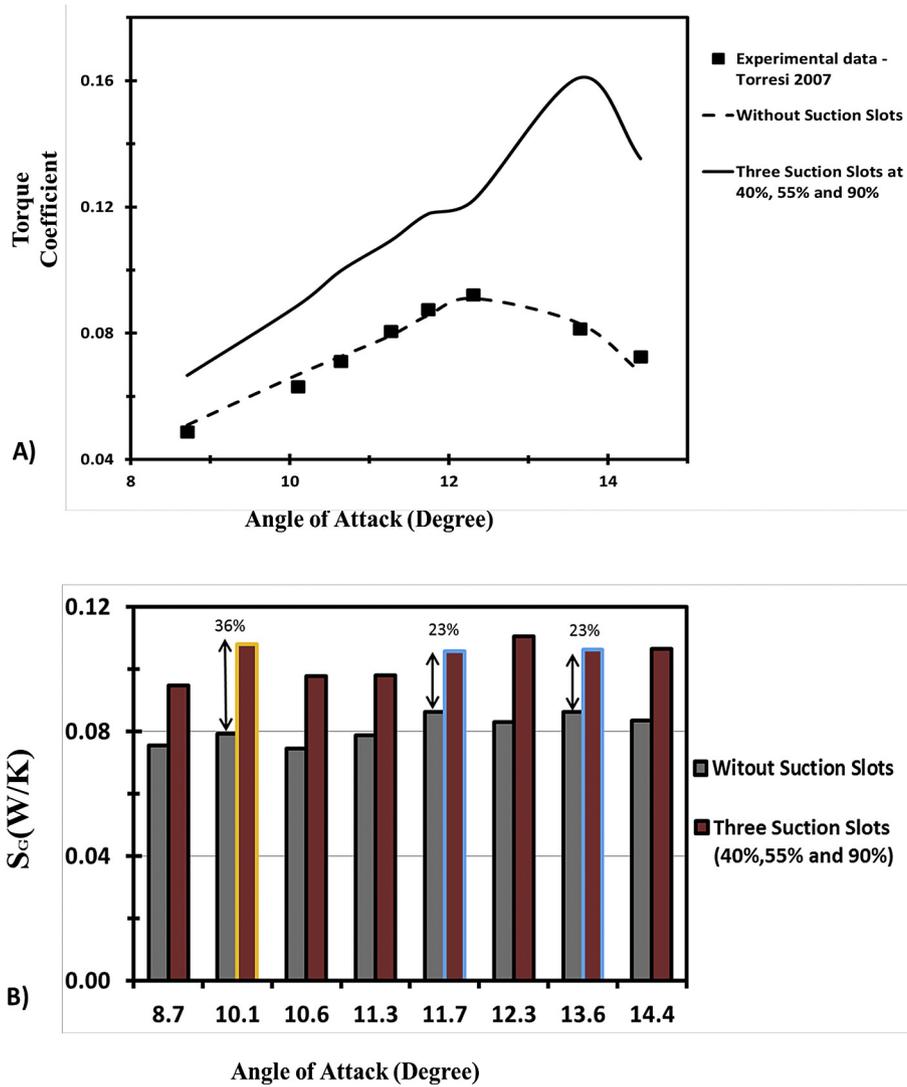


Fig. 6. The optimum location for three suction slots at different angles with non-oscillating velocity. A) Torque coefficient. B) S<sub>c</sub>.

Table 8  
Four suction slots with P<sub>ss</sub> and ΔP<sub>ss</sub> equal to 0.05.

Reference suction slot	Location of fourth suction slot															Max. value
	20%	25%	30%	35%	40%	45%	50%	55%	60%	65%	70%	75%	80%	85%	90%	
R 25%, 30% and R 55%	0.143	—	—	0.151	0.13	0.142	0.124	—	0.155	0.122	0.144	0.155	0.123	0.148	0.156	88%
R 40%, 45% and R 50%	0.131	0.122	0.144	0.149	—	—	—	0.13	0.144	0.158	0.153	0.150	0.146	0.156	0.154	90%
R 40%, 45% and R 55%	0.141	0.137	0.132	0.137	—	—	0.143	—	0.160	0.135	0.157	0.151	0.156	0.155	0.125	92%
R 40%, 45% and R 60%	0.132	0.151	0.136	0.142	—	—	0.144	0.10	—	0.148	0.148	0.145	0.150	0.151	0.137	92%
R 40%, 45% and R 75%	0.121	0.148	0.133	0.144	—	—	0.150	0.11	0.145	0.147	0.135	—	0.149	0.145	0.147	92%
R 40%, 45% and R 85%	0.128	0.143	0.136	0.136	—	—	0.156	0.65	0.151	0.142	0.138	0.745	0.136	—	0.139	88%
R 40%, 50% and R 65%	0.147	0.134	0.125	0.140	—	0.158	—	0.13	0.140	—	0.101	0.148	0.142	0.128	0.147	90%
R 40%, 55% and R 75%	0.143	0.099	0.136	0.141	—	0.151	0.157	—	0.150	0.150	0.148	—	0.147	0.139	0.156	89%
R 40%, 55% and R 80%	0.137	0.134	0.131	0.133	—	0.125	0.122	—	0.139	0.151	0.138	0.147	—	0.141	0.136	92%
R 40%, 55% and R 90%	0.141	0.134	0.135	0.149	—	0.134	0.148	—	0.143	0.137	0.155	0.156	0.136	0.157	—	89%

Maximum value

**Table 9**  
Four suction slots with  $P_{ss}$  and  $\Delta P_{ss}$  equal to 0.01.

Reference suction slot	Location of fourth suction slot										Max. value
	65%	64%	63%	62%	61%	60%	59%	58%	57%	56%	
R 40%, 45% and R 55%	0.1348	0.1396	0.1507	0.1477	0.1400	0.1596	0.1501	0.1316	0.1359	0.1229	92%
Reference suction slot	59%	58%	57%	56%	55%	54%	53%	52%	51%	50%	Max. value
R 40%, 45% and R 60%	0.1409	0.1368	0.1306	0.1376	0.1596	0.1410	0.1344	0.1411	0.1427	0.1441	92%
Reference suction slot	50%	49%	48%	47%	46%	45%	44%	43%	42%	41%	Max. value
R 40%, 55% and R 60%	0.1239	0.1326	0.1338	0.1366	0.1291	0.1596	0.1364	0.1253	0.1126	0.1349	92%
Reference suction slot	44%	43%	42%	41%	40%	39%	38%	37%	36%	35%	Max. value
R 45%, 55% and R 60%	0.1302	0.1073	0.1397	0.1421	0.1596	0.1438	0.1318	0.1387	0.1307	0.1369	92%

Maximum value

- Hub-to-tip ratio = 0.65;
- Solidity = 0.64;
- The uncertainty in the measurements = 5%

Second experimental data [83] is adopted to validate the unsteady sinusoidal wave inlet velocity. The second experimental data involved the investigation of unsteady forces ( $F_D$ ) acting on a square cylinder in oscillating flow with nonzero mean velocity. The oscillating air flows are generated by a unique AC servomotor wind tunnel. The generated velocity histories are almost exact sinusoidal waves.

For unsteady flow with non-oscillating and oscillating velocity, it was validated in Ref. [6], where, a comparison between the measured torque coefficient from Refs. [87,96,97] and calculated torque coefficient from CFD results at Reynolds number of  $2 \times 10^4$ . This comparison yields a very good agreement. It can be noted that approximately the same stall condition value of the experimental reference was recorded by the computational model. Furthermore, for an unsteady flow with sinusoidal inlet velocity, Fig. 2 shows a good agreement between measured drag force from Ref. [83] and predicted drag force from CFD at two different frequencies (2 Hz and 1 Hz). It can be shown from Fig. 2 that the computational model has almost the same behaviour of oscillating flow condition as the reference. Table 3 lists the error percentages for each frequency.

#### 4. Analysis and discussion of results

A multi-suction slot with a certain diameter ( $D_{ss}$ ) equal to 0.1% [6] from the blade chord at various locations from the leading edge was created, with a shape of NACA0015 with stall angle equal to  $13.6^\circ$  and Reynolds number equal to  $2 \times 10^5$  from Refs. [87,96,97], see Fig. 3. The locations for the suction slots were changed in order

to obtain an optimum value of  $C_T$ . The test cases investigated were under unsteady flow with non-oscillating velocity at the first to indicate to the best locations and then take the best cases to investigate under sinusoidal wave condition to decide which one has the highest  $C_T$  and which one has the lowest  $S_G$ . The sinusoidal wave inlet flow boundary condition is having the same specifications as that in section 3.1 and equation (25). Finally, a comparative analysis was made based on conditions relevant to northern coast of Egypt with different sinusoidal wave frequencies ( $f$  equal to 0.25, 0.167 and 1.25 Hz).

##### 4.1. Multi suction slots (two, three and four)

The two suction slots were investigated by making the first suction slots as a reference ( $L_{RSS}$ ) and changing the location (x axis direction) of the second suction slots ( $L_{SS}$ ) by pitch distance ( $P_{ss}$ ) equal to 0.05 at each trial. Considering that, the minimum distance between the two suction slots ( $\Delta P_{ss}$ ) was equal to 0.05.

Table 4 provides the details about all two suction slots trial with  $P_{ss}$  and  $\Delta P_{ss}$  equals to 0.05 to improve the torque coefficient at the stall angle  $13.6^\circ$ . It can be noted that the  $L_{RSS}$  equal to 40% and  $L_{SS}$  equal to 45% gives a higher torque coefficient than others, where the torque coefficient increases about 82% higher than the aerofoil without suction slot at the stall angle. Therefore, to get more improvement in the torque coefficient, the value of  $P_{ss}$  and  $\Delta P_{ss}$  was changes to 0.01 around the  $L_{SS}$  40% and  $L_{SS}$  45%. It can be concluded that two suction slots at  $L_{RSS}$  40% and  $L_{SS}$  44% with  $P_{ss}$  equal to 0.01 give a higher torque coefficient than others by 84% at the stall angle in Fig. 4 A). On the other hand, two suction slots at  $L_{RSS}$  45% and  $L_{SS}$  49% give also a higher torque coefficient than others by 84% at the stall angle from Fig. 4 B).

From Table 4 and Fig. 4, it can be noted that the three optimum

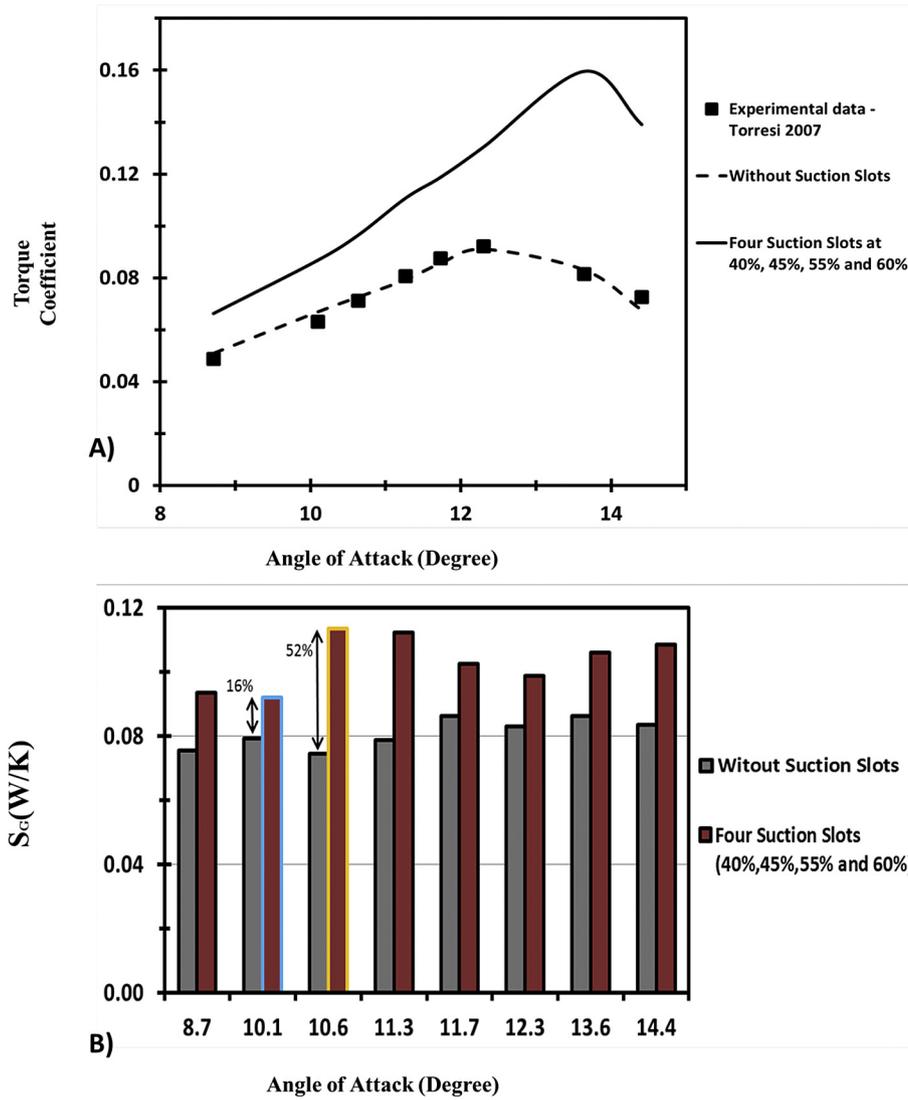


Fig. 7. The optimum location for four suction slots at different angles with non-oscillating velocity A) Torque coefficient B)  $S_c$ .

locations for two suction slots were  $L_{ss}$  40% and 45% with  $P_{ss}$  equal to 0.05, in addition to  $L_{ss}$  40% and 44% and  $L_{ss}$  45% and 49% with  $P_{ss}$  equal to 0.01. Table 5 demonstrates the effect of the three optimum locations for two suction slots on the torque coefficient at different

angles. It can be noted that the two suction slots at  $L_{ss}$  40% and 45% improve the torque coefficient before the stall by 37.2% and after the stall by 95.5%. Also, the two suction slots at  $L_{ss}$  40% and 44% improve the torque coefficient before the stall by 33.5% and after the stall by 97.5%. Finally, the two suction slots at  $L_{ss}$  45% and 49% improve the torque coefficient before the stall by 36.7% and after the stall by 99%.

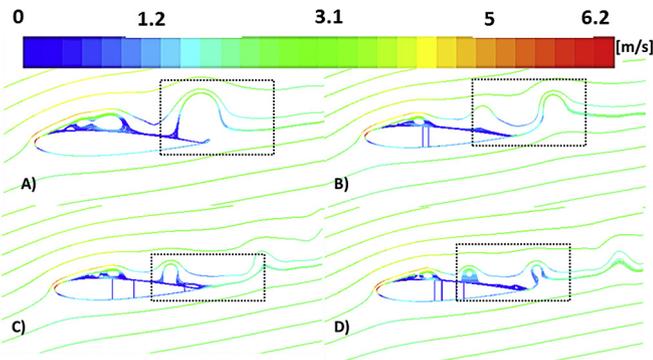


Fig. 8. Path-line coloured by mean velocity magnitude at 13.6°, A) NACA0015 without slots B)  $L_{ss}$  40% and 44% C)  $L_{ss}$  40%, 55% and 90% D)  $L_{ss}$  40%, 45%, 55%, and 60%. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

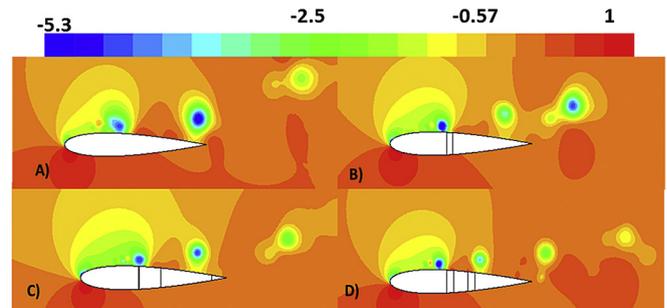


Fig. 9. Contours of pressure coefficient around the aerofoil at 13.6°, A) NACA0015 without slots B)  $L_{ss}$  40% and 44% C)  $L_{ss}$  40%, 55% and 90% D)  $L_{ss}$  40%, 45%, 55%, and 60%.

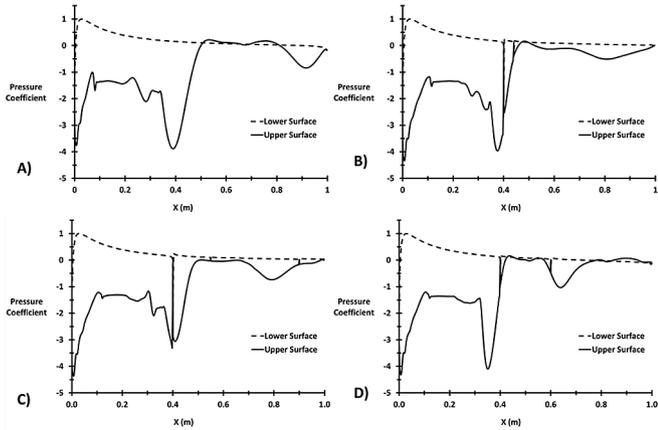


Fig. 10. Pressure distribution on the upper and lower surface at 13.6°. A) NACA0015 without slots B)  $L_{ss}$  40% and 44% C)  $L_{ss}$  40%, 55% and 90% D)  $L_{ss}$  40%, 45%, 55%, and 60%.

The suction slots have a negative effect on the entropy generation, where the global entropy generation rate increases at all angles by 24% before the stall and 23% after the stall due to suction slots at ( $L_{ss}$  40% and 45%). Where, the 11.3° angle of attack (before the stall) has the highest difference in global entropy generation rate by 38%. On the other hand, 11.7° angle of attack (before the stall) has the lowest difference in global entropy generation rate by 13% due to suction slots at Fig. 5 A). Furthermore, the suction slots at ( $L_{ss}$  40% and 44%) cause increase in the global entropy generation rate value as average for all angles by 21% before the stall and 26% after the stall at Fig. 5 B). The 10.1° angle of attack (before the

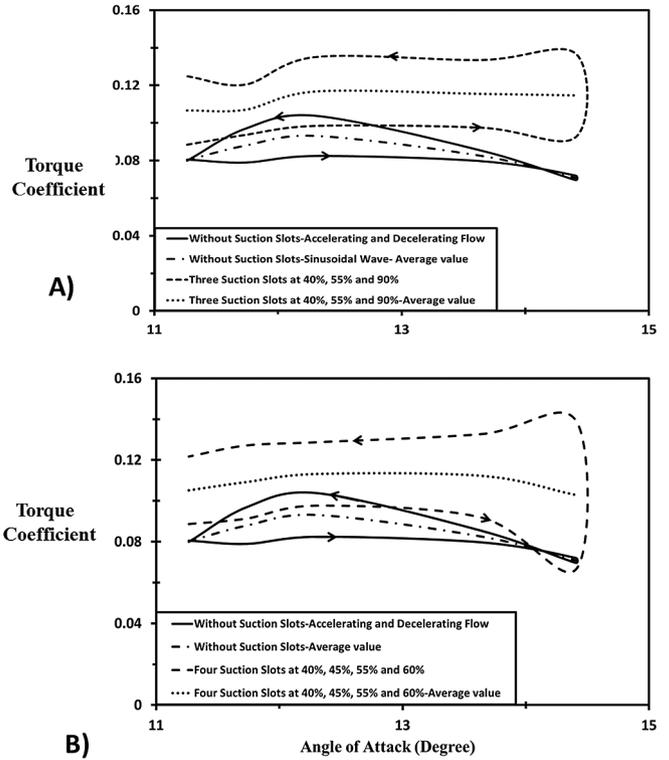


Fig. 12. The hysteretic behaviour and the average torque coefficient for the optimum locations of the suction slots with sinusoidal velocity, A) Three suction slots B) Four suction slots.

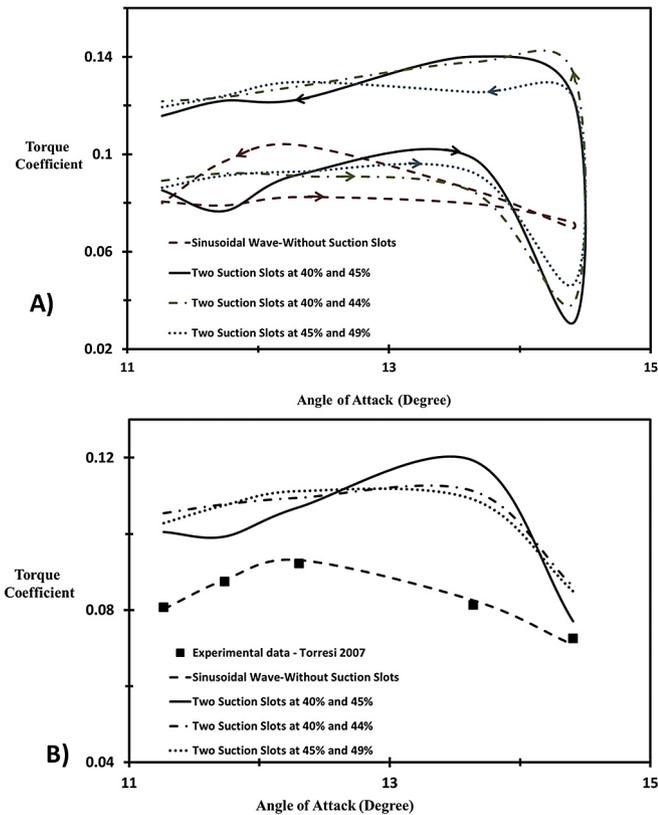


Fig. 11. Comparison between the optimum locations for two suction slots at different angles with sinusoidal inlet velocity A) The hysteretic behaviour B) The average torque coefficient.

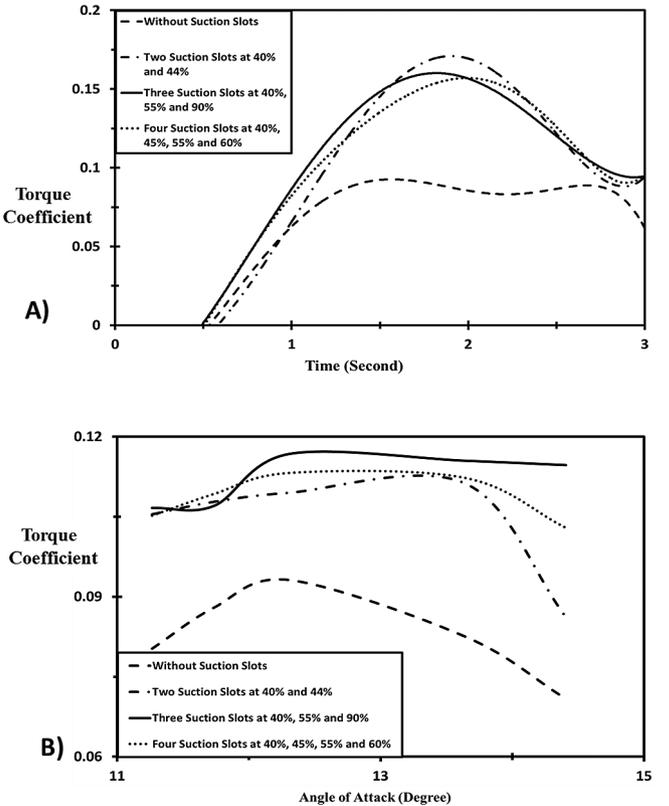
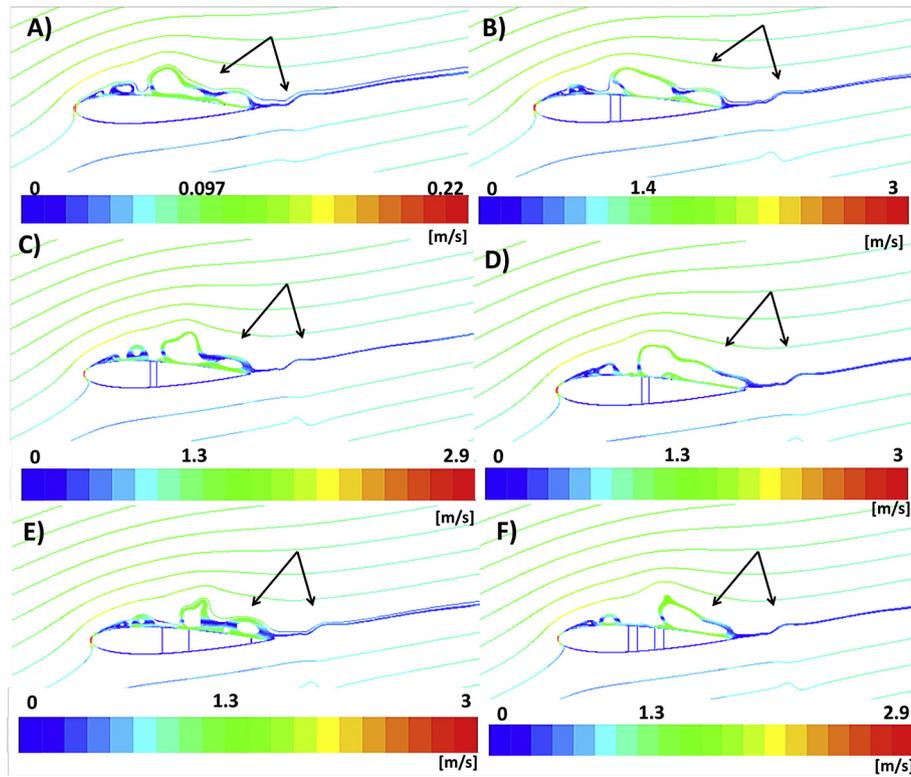


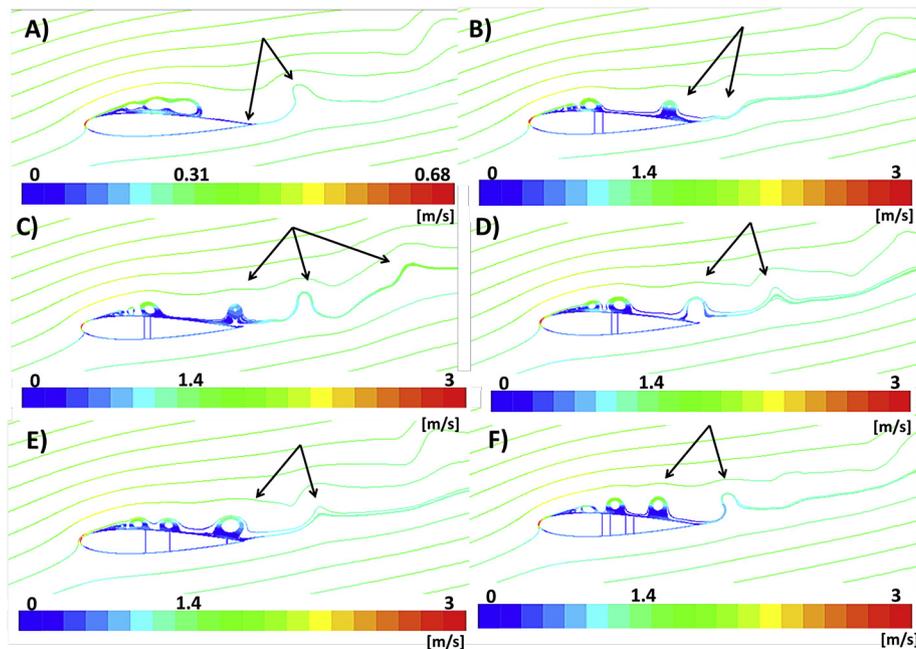
Fig. 13. The comparison between the two suction slots, three suction slots and four suction slots A) The instantaneous torque coefficient B) Average value of torque coefficient.



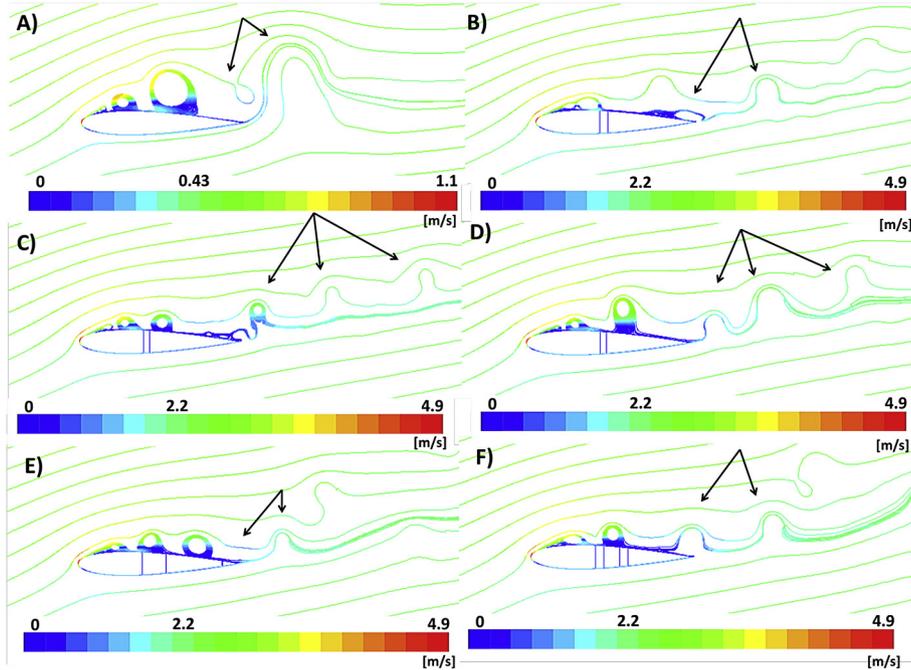
**Fig. 14.** Path-line coloured by mean velocity magnitude at velocity equal to 1.8 m/s (acceleration flow) A) NACA0015 without slots B)  $L_{ss}$  40% and 45% C)  $L_{ss}$  40% and 44% D)  $L_{ss}$  45% and 49% E)  $L_{ss}$  40%, 55% and 90% F)  $L_{ss}$  40%, 45%, 55%, and 60% (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.).

stall) has the lowest difference in global entropy generation rate by 14% due to suction slots and, the  $14.4^\circ$  angle of attack (after the stall) has the highest difference by 29%. Finally, for suction slots at ( $L_{ss}$  45% and 49%) at Fig. 5 C), the global entropy generation rate

increases as average in all angles by 21% before the stall and 22% after the stall. Where, the  $10.6^\circ$  (before the stall) has the highest difference in global entropy generation rate by 44% and, the  $12.3^\circ$  (before the stall) has the lowest by 11% due to suction slots. This



**Fig. 15.** Path-line coloured by mean velocity magnitude at maximum velocity equal to 2.92 m/s A) NACA0015 without slots B)  $L_{ss}$  40% and 45% C)  $L_{ss}$  40% and 44% D)  $L_{ss}$  45% and 49% E)  $L_{ss}$  40%, 55% and 90% F)  $L_{ss}$  40%, 45%, 55%, and 60%. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.).

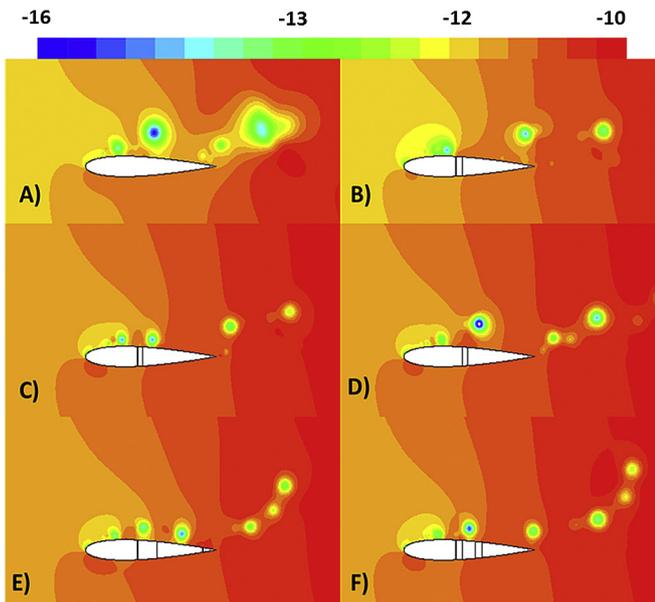


**Fig. 16.** Path-line coloured by mean velocity magnitude at velocity equal to 1.8 m/s (deceleration flow), A) NACA0015 without slots B)  $L_{SS}$  40% and 45% C)  $L_{SS}$  40% and 44% D)  $L_{SS}$  45% and 49% E)  $L_{SS}$  40%, 55% and 90% F)  $L_{SS}$  40%, 45%, 55%, and 60%. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

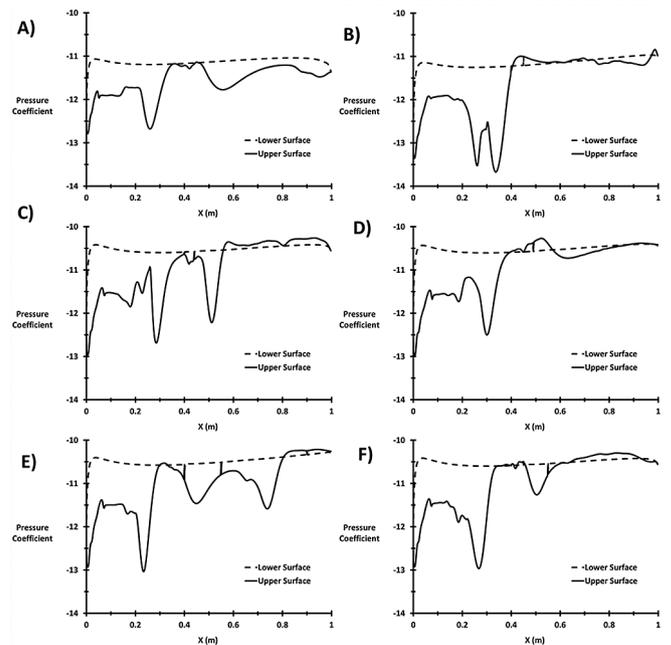
phenomenon suggests that the change in velocity gradient due to the suction slot has a direct impact on the entropy generation.

A third suction slot by  $P_{SS}$  equal to 0.05 was added to all aerofoils with two suction slots that have higher than 70% improvement in the torque coefficient at the stall angle  $13.6^\circ$ . Table 6 provides all three suction slots trial with  $P_{SS}$  and  $\Delta P_{SS}$  equal to 0.05 to improve the torque coefficient at the stall angle ( $13.6^\circ$ ). It can

be noted that the  $L_{RSS}$  equal to 40%–55% and  $L_{SS}$  equal to 90% gives a higher torque coefficient than others, where the torque coefficient increases about 94% higher than the aerofoil without suction slot. Therefore, to get more improvement in the torque coefficient, the value of  $P_{SS}$  and  $\Delta P_{SS}$  was changes to 0.01 around the  $L_{SS}$  40%, 55% and 90%. From Table 7, it can be noted that no improvement on the torque coefficient by change  $P_{SS}$  from 0.05 to 0.01 around the  $L_{SS}$



**Fig. 17.** Contours coloured by pressure coefficient at velocity equal to 1.8 m/s (deceleration flow), A) NACA0015 without slots B)  $L_{SS}$  40% and 45% C)  $L_{SS}$  40% and 44% D)  $L_{SS}$  45% and 49% E)  $L_{SS}$  40%, 55% and 90% F)  $L_{SS}$  40%, 45%, 55%, and 60%. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



**Fig. 18.** Pressure coefficient at upper and lower surface with velocity equal to 1.8 m/s (deceleration flow), A) NACA0015 without slots B)  $L_{SS}$  40% and 45% C)  $L_{SS}$  40% and 44% D)  $L_{SS}$  45% and 49% E)  $L_{SS}$  40%, 55% and 90% F)  $L_{SS}$  40%, 45%, 55%, and 60%.

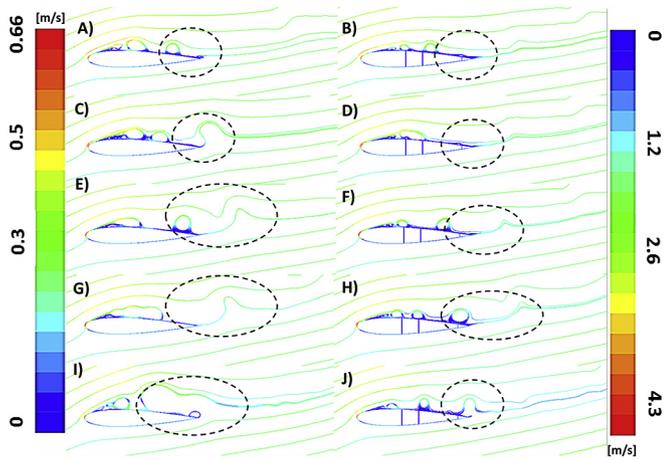


Fig. 19. Path-line coloured by mean velocity magnitude around NACA0015 without and with slots at  $L_{SS}$  40%, 55% and 90% at maximum velocity equal to 2.92 (m/s) A) and B) 11.3, C) and D) 11.7, E) and F) 12.3, G) and H) 13.6, I) and J) 14.4. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

40%, 55% and 90%.

It is clearly noted that the three suction slots at ( $L_{SS}$  40%, 55% and 90%) improve the torque coefficient before the stall by 35.2% and after the stall by 97%, see Fig. 6 A). On the other hand, the global entropy generation rate increases for all angles by 29% before the stall and 25% after the stall as average value at Fig. 6 B). Where, the  $10.1^\circ$  angle of attack (before the stall) has the highest difference in global entropy generation rate by 36% due to suction slots at ( $L_{SS}$  40%, 55% and 90%), and the  $11.7$  (before the stall) and  $13.6$  (after the stall)  $^\circ$  have the lowest difference by 23%.

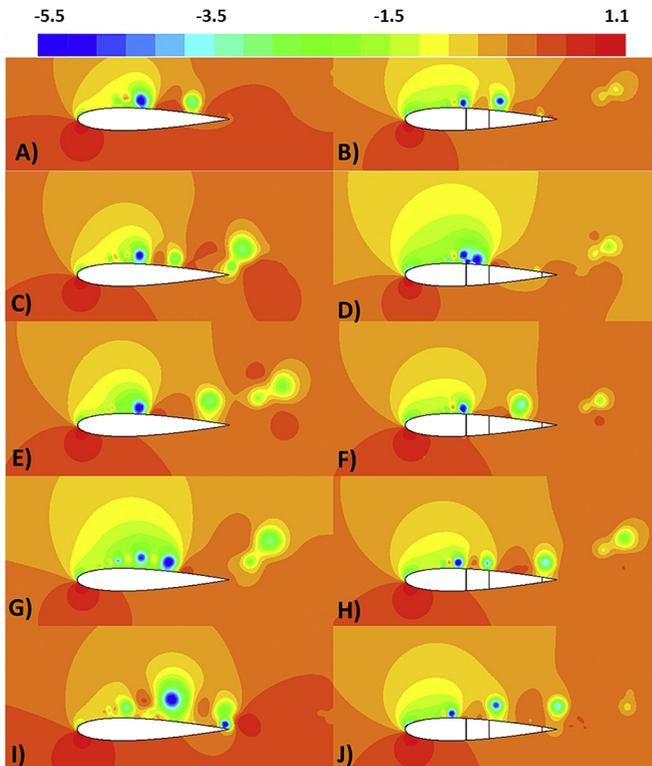


Fig. 20. Contours of the pressure coefficient around NACA0015 A) and B) 11.3, C) and D) 11.7, E) and F) 12.3, G) and H) 13.6, I) and J) 14.4.

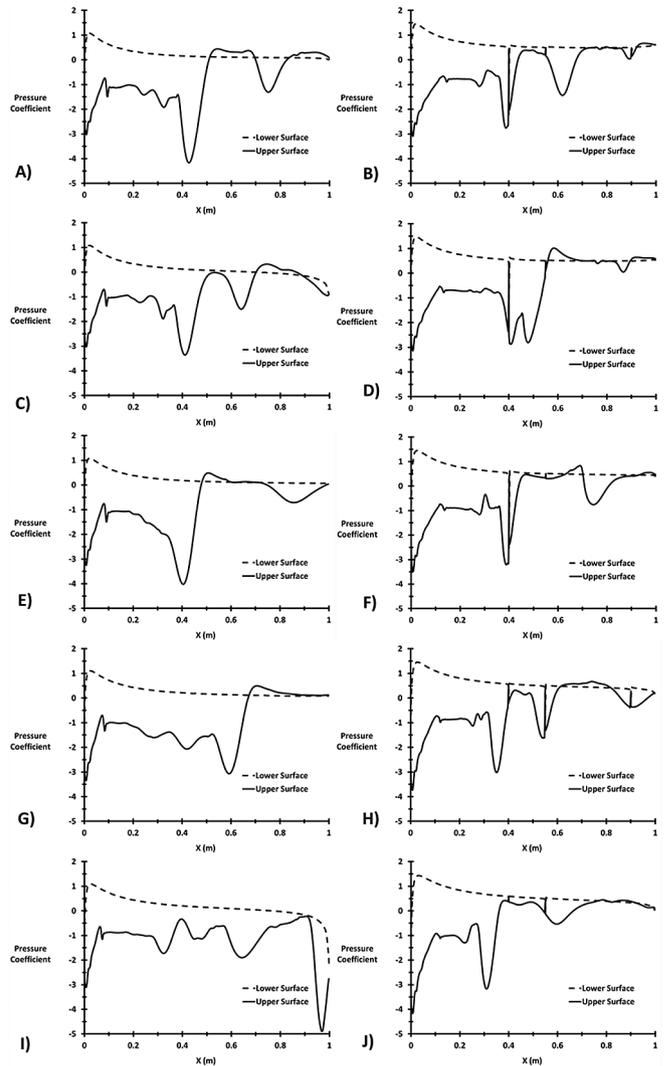


Fig. 21. Pressure coefficient at upper and lower surface of NACA0015 A) and B) 11.3, C) and D) 11.7, E) and F) 12.3, G) and H) 13.6, I) and J) 14.4.

A fourth suction slot with  $P_{SS}$  equal to 0.05 was added to all aerofoils with three suction slots that have higher than 80% improvement in the torque coefficient at the stall angle ( $13.6^\circ$ ). Table 8 shows the effect of four suction slots on the torque coefficient at the stall angle ( $13.6^\circ$ ) with  $P_{SS}$  and  $\Delta P_{SS}$  equal to 0.05 to improvement the torque coefficient. The  $L_{RSS}$  equal to 40%–45% - 55% and  $L_{SS}$  equal to 60% gives a higher improvement in the torque coefficient by 92%. As in two and three suction slots, the value of  $P_{SS}$  and  $\Delta P_{SS}$  changes to 0.01 around the  $L_{SS}$  40%, 45%, 55% and 60% to get more improvement in the torque coefficient. From Table 9, it can be noted that the  $L_{SS}$  40%, 45%, 55% and 60% give highest improvement on the torque coefficient with  $P_{SS}$  equal to 0.01. The four suction slots ( $L_{SS}$  40%, 45%, 55% and 60%) improve the torque coefficient before the stall by 35.8% and after the stall by 99%, see Fig. 7 A). Otherwise, the global entropy generation rate increases by 29% before the stall and 26% after the stall at Fig. 7 B). Where, the  $10.1^\circ$  has the lowest difference in global entropy generation rate by 16% and, the  $10.6^\circ$  has the highest difference by 52%. The path line coloured by the mean velocity magnitude around the NACA0015 without and with slots at the stall angle  $13.6^\circ$  was presented at Fig. 8. The slots effect was very cleared on the separation layers at the trailing edge area and it extends to the area beyond the trailing

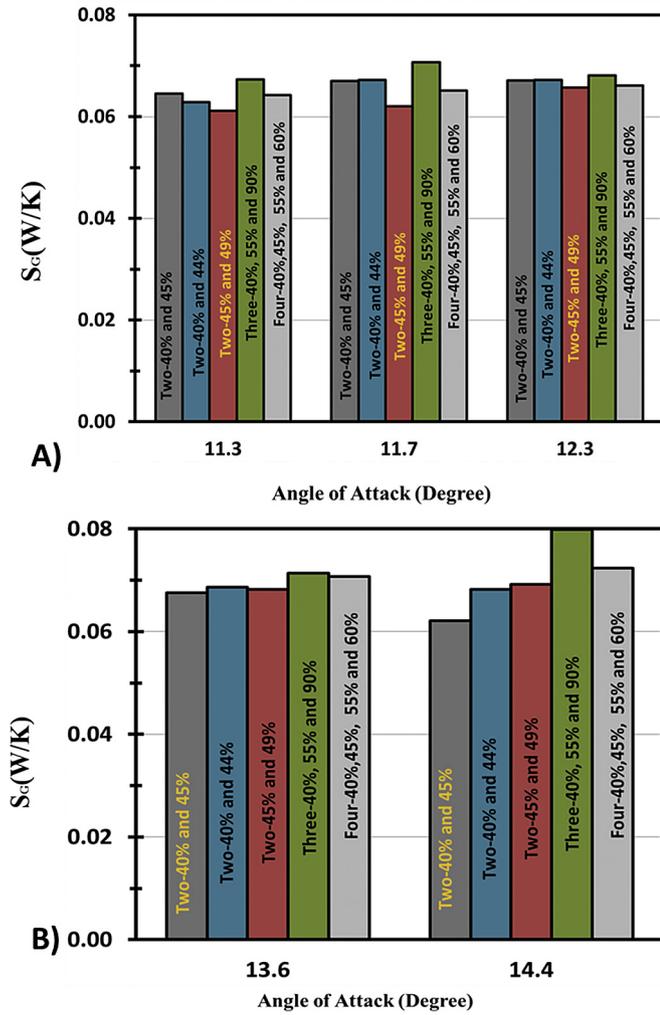


Fig. 22. Comparison between the average value of the  $S_G$  for the compression cycle for NACA0015 with two, three and four suction slots, A) Before the stall B) After the stall.

edge which, leads to delay the stall. Therefore, the NACA0015 with suction slots and Reynolds number equal to  $2 \times 10^5$  not have the stall condition at  $13.6^\circ$ . This improvement can be achieved by two, three or four slots at different location. Where, each case of them has different behaviour.

The low pressure areas, at the trailing edge of the NACA0015 without slots, were caused the separation layer. On the other hand, the suction slots affect directly on these areas and decrease from its value and this leads to decrease the separation layers (Fig. 9). The difference between the upper and lower surface was decreased by the suction slots and this leads to decrease from the disturbance and the separation layers (Fig. 10). The pressure distribution at the upper and lower surface was depending on the number and location of the slots. Therefore, the two, three and four slots were investigated under sinusoidal wave condition in next section.

#### 4.2. Optimum location for multi-suction slots based on first law analysis

From the previous section, it was noted that there are five scenarios for the suction slots location, which gives higher torque coefficient at the stall regime:

- 1 Two Suction Slots ( $L_{ss}$  40% and 45%) with  $P_{ss} = 0.05$

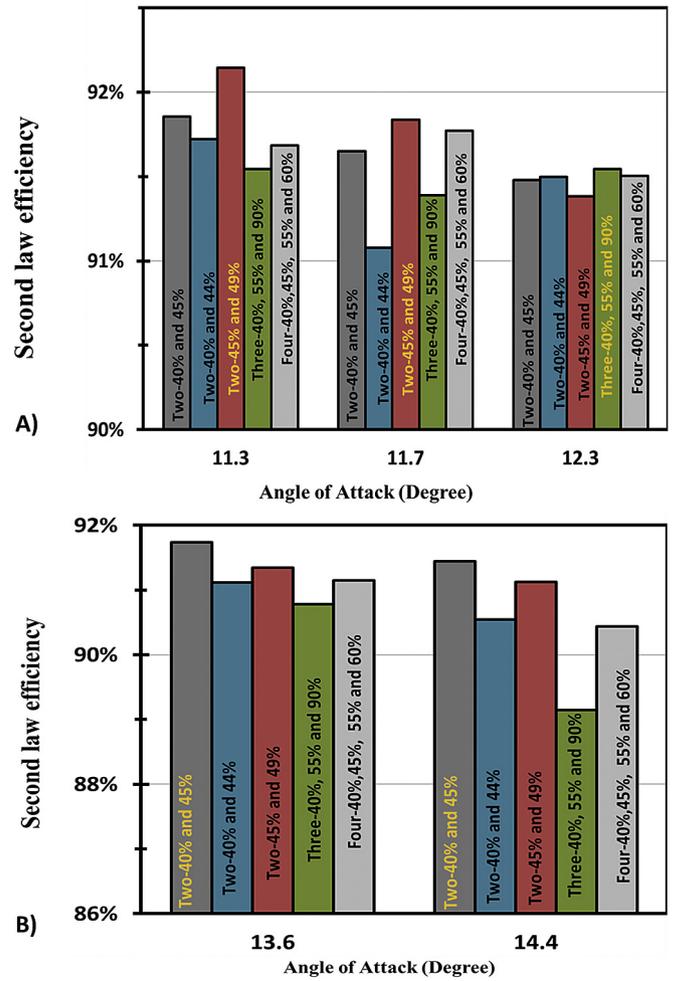


Fig. 23. Comparison between the second law efficiency for the compression cycle for NACA0015 with two, three and four suction slots A) Before the stall B) After the stall.

- 2 Two Suction Slots ( $L_{ss}$  40% and 44%) with  $P_{ss} = 0.01$
- 3 Two Suction Slots ( $L_{ss}$  45% and 49%) with  $P_{ss} = 0.01$
- 4 Three Suction Slots ( $L_{ss}$  40%, 55% and 90%) with  $P_{ss} = 0.05$
- 5 Four Suction Slots ( $L_{ss}$  40%, 45%, 55%, and 60%) with  $P_{ss} = 0.05$

In this section, the optimum locations for multi-suction slots based on the torque coefficient were determined under sinusoidal

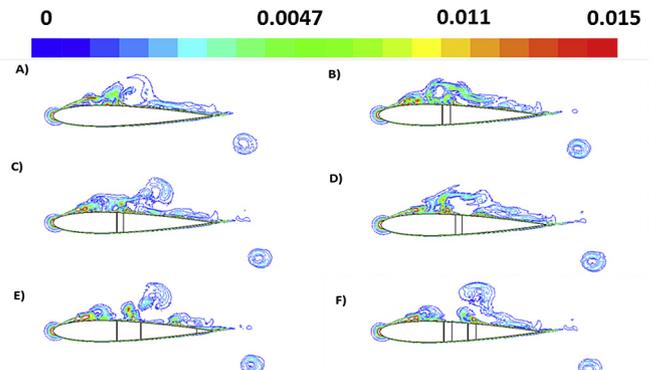
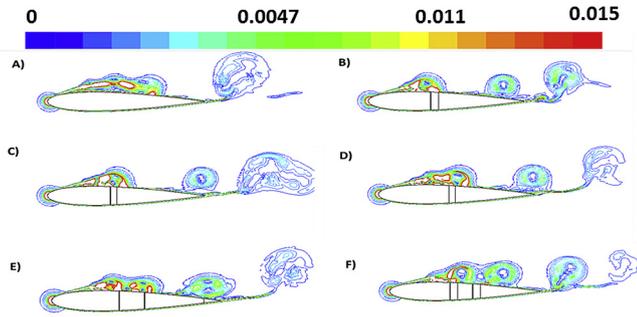


Fig. 24. Contour of  $S_G$  at 1.8 m/s (accelerating) A) NACA0015 without slots B)  $L_{ss}$  40% and 45% C)  $L_{ss}$  40% and 44% D)  $L_{ss}$  45% and 49% E)  $L_{ss}$  40%, 55% and 90% F)  $L_{ss}$  40%, 45%, 55%, and 60%.

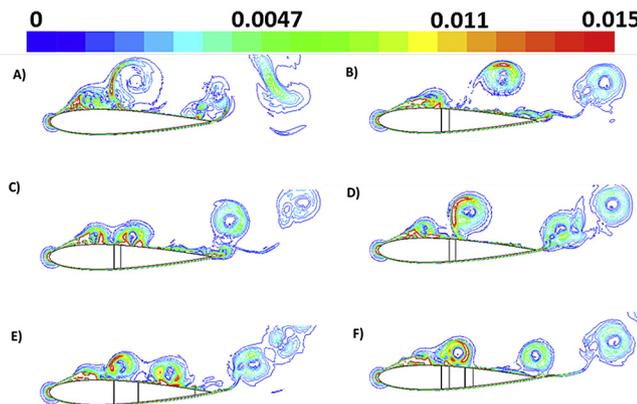


**Fig. 25.** Contour of  $S_G$  at 2.92 m/s (maximum velocity) A) NACA0015 without suction slots B)  $L_{SS}$  40% and 45% C)  $L_{SS}$  40% and 44% D)  $L_{SS}$  45% and 49% E)  $L_{SS}$  40%, 55% and 90% F)  $L_{SS}$  40%, 45%, 55%, and 60%.

wave condition. Fig. 11 compares the torque coefficients for the two suction slots aerofoil at different locations ( $L_{SS}$  40% and 45%), ( $L_{SS}$  40% and 44%) and ( $L_{SS}$  45% and 49%). Fig. 11 A) illustrates the hysteretic behaviour due to the reciprocating flow which shows a delay in the stall regime and an improvement in the torque coefficient. The two suction slots aerofoil with  $L_{SS}$  of 40% and 44% has a higher improvement of torque coefficient than that with  $L_{SS}$  of 40% and 45% by 6.3% before the stall and 1.5% after the stall. Moreover, the former aerofoil also has a higher torque coefficient than that with  $L_{SS}$  of 45% and 49% by 1% before the stall and 2.5% after the stall (Fig. 11 B).

Fig. 12 shows the effect of adding three suction slots at ( $L_{SS}$  40%, 55% and 90%) and four suction slots at ( $L_{SS}$  40%, 45%, 55%, and 60%) under sinusoidal flow condition on the hysteretic behaviour. From this Figure, it can be noted that in both cases a delay in the stall regime occurred. In addition, the torque coefficient was improved by 26.7% before the stall and 51% after the stall due to the addition of three suction slots (Fig. 12 A). However, the addition of four suction slots resulted in torque coefficient improvement by 25.7% before the stall and 40.5% after the stall (Fig. 12 B).

From Fig. 13, it is clearly noted that adding three suction slots at ( $L_{SS}$  40%, 55% and 90%) provided the highest improvement of torque coefficient, from both the instantaneous and average value, compared to all the scenarios that were mentioned in this section. By comparing this aerofoil against the two suction slots aerofoil with optimum locations ( $L_{SS}$  40% and 44%), an improvement of

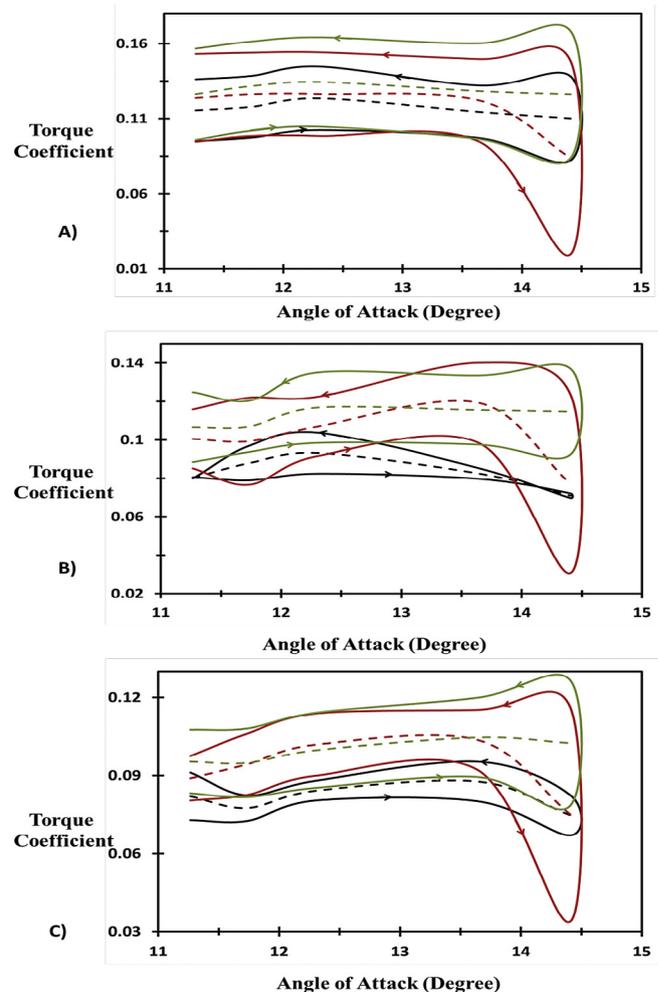


**Fig. 26.** Contour of  $S_G$  at 1.8 m/s (decelerating) A) NACA0015 without slots B)  $L_{SS}$  40% and 45% C)  $L_{SS}$  40% and 44% D)  $L_{SS}$  45% and 49% E)  $L_{SS}$  40%, 55% and 90% F)  $L_{SS}$  40%, 45%, 55%, and 60%.

torque coefficient of 2.7% before the stall and 22.5% after the stall was observed. Moreover, by comparing the same aerofoil against the four suction slots aerofoil ( $L_{SS}$  40%, 45%, 55%, and 60%), an improvement of torque coefficient of 1% before the stall and 10.5% after the stall was observed.

The path-line coloured by mean velocity magnitude highlights the improvement effect of adding a suction slot on the separation layers at the trailing edge region in Fig. 14 (acceleration flow) was small compared with Figs. 15 and 16. Where, the separation layers at the area around the trailing edge increased especially at the deceleration flow in the second half of the compression cycle (Fig. 16). Furthermore, it can be noted that the low pressure areas around the trailing edge decrease due to the slots addition from Figs. 17 and 18. The pressure difference between the lower and upper surfaces was decreased as a result of adding the slots. Therefore, the disturbances in the path line at the trailing edge area and the area extended beyond it was decreased. This leads to delay the stall and improve the torque coefficient.

Fig. 19 shows the effect of adding three suction slots on the boundary layer separation before and after the stall condition via the mean velocity magnitude path-lines. It can be noted that the improvement effect of adding suction slot on separation layers increased in stall regime for both 13.6 and 14.4°. The pressure



**Fig. 27.** The hysteretic behaviour and the average torque coefficient “— without slots –average value”, “- ( $L_{SS}$  40% and 45%) –average value”, “- ( $L_{SS}$  40%, 55% and 90%) –average value”, A)  $t_{sin} = 4$  s B)  $t_{sin} = 6$  s C)  $t_{sin} = 8$  s.

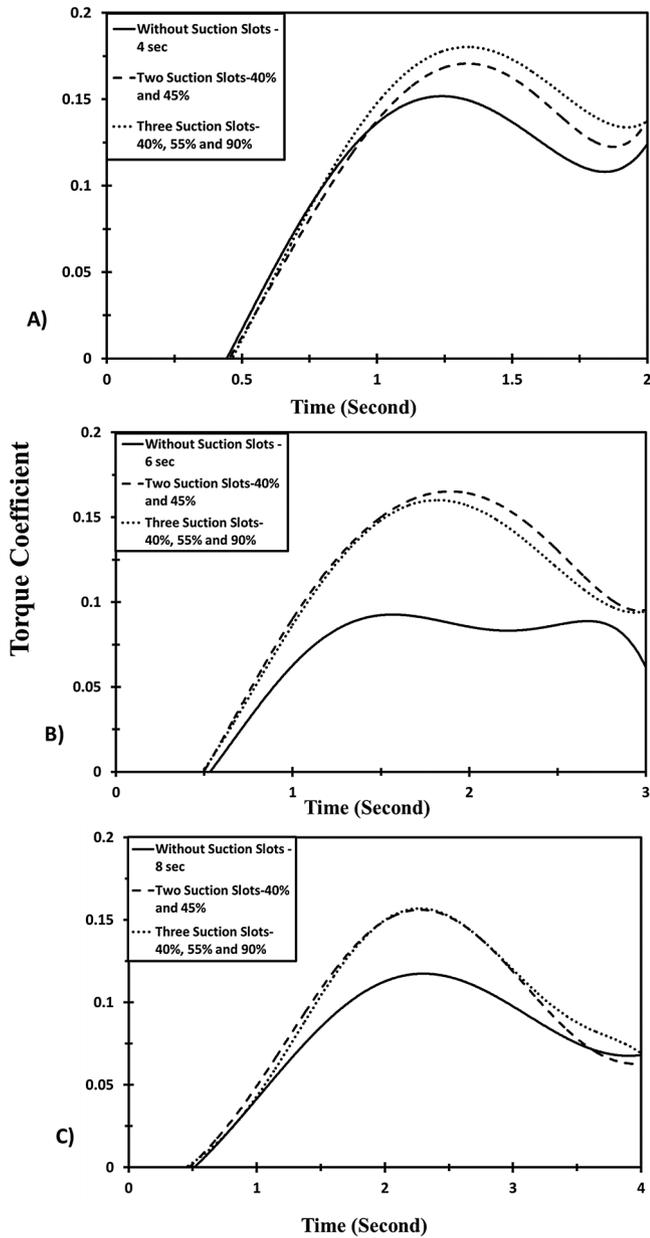


Fig. 28. The instantaneous torque coefficient for the NACA0015 without and with suction slots at optimum locations in 13.6° A)  $t_{sin} = 4$  s B)  $t_{sin} = 6$  s C)  $t_{sin} = 8$  s.

distributions around the aerofoil and at the upper and lower surfaces for different angles of attack were shown in Figs. 20 and 21. Where, the lift column is for the NACA0015 without slots and the right column is for the NACA0015 with three slots at  $L_{SS}$  40%, 55% and 90% with maximum velocity equal to 2.92 m/s. The addition of three slots affects directly the low pressure zones that appear around the trailing edge area and the upper surface of the aerofoil. Where, this low pressure zones were the main reason for the separation layers to be formed. For all angles, the aerofoil with three slots showed an improvement in the pressure distribution and decreased the separation layers especially for the stall angle of 14.4° in Fig. 19I) and J).

4.3. EGM method

The numerical simulations were used to obtain local entropy

viscosity predictions from the different five scenarios for the locations of suction slots. Figs. 22 and 23 highlight the comparison between the ( $L_{SS}$  40% and 45%), ( $L_{SS}$  40% and 44%), ( $L_{SS}$  45% and 49%), ( $L_{SS}$  40%, 55% and 90%) and ( $L_{SS}$  40%, 45%, 55%, and 60%). The comparison was provided as an average value for the compression cycle with different angles of attack. From Fig. 22 A) it can be noted that the minimum value for the global entropy generation rate occurs with ( $L_{SS}$  45% and 49%) by 20.24% increase in  $S_G$  before the stall. On the other hand, the minimum value for the global entropy generation rate occurs with ( $L_{SS}$  40% and 45%) by 14.54% increase in  $S_G$  after the stall; see Fig. 22 B). Furthermore, the two suction slots ( $L_{SS}$  40% and 45%) give minimum  $S_G$  as an average value for the compression cycle before and after the stall by 20.5% increase in  $S_G$  value.

From Fig. 23 it can be concluded that the ( $L_{SS}$  45% and 49%) gives the maximum value of second law efficiency by 0.38% before the stall, and, the ( $L_{SS}$  40% and 45%) gives the maximum value after the stall by 1.19%. Furthermore, the two suction slots ( $L_{SS}$  40% and 45%) give maximum value for the second law efficiency as an average value for the compression cycle before and after the stall by 0.72%. The increases in  $S_G$  (Fig. 22) leads to decrease in second law efficiency in some cases than that without suction slots, such as the two suction slots at  $L_{SS} = 40%$  and  $L_{SS} = 44%$  before the stall which the second law efficiency decreased by (0.01%), and three suction slots at  $L_{SS} = 40%$ ,  $L_{SS} = 55%$  and  $L_{SS} = 90%$  after the stall which the second law efficiency decreased by (0.61%). There were no significant changes in both the global entropy generation rate values and

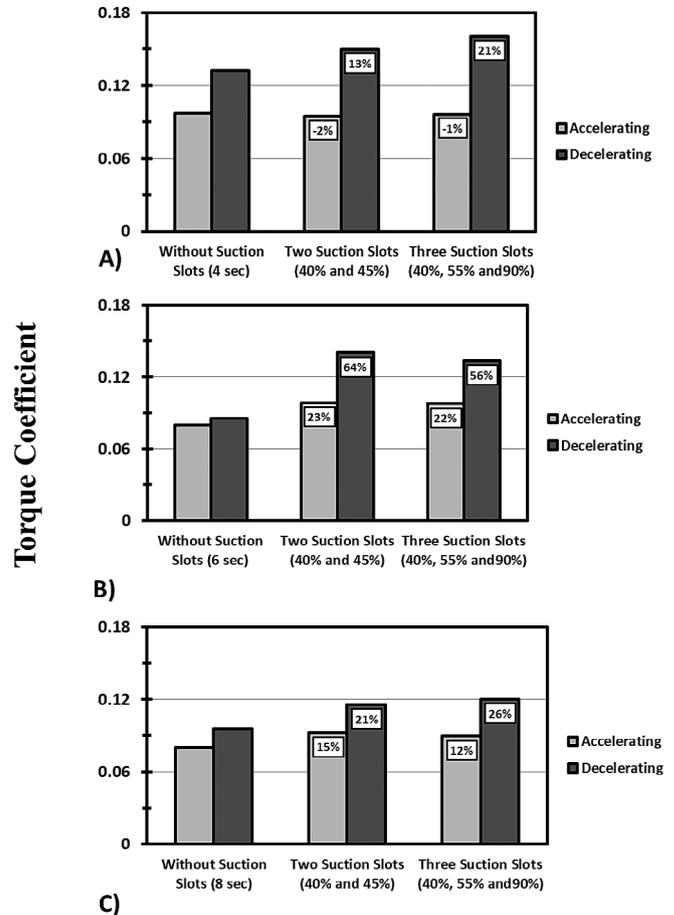


Fig. 29. The average torque coefficient at the accelerating and decelerating flow for the NACA0015 without and with suction slots at optimum locations A)  $t_{sin} = 4$  s B)  $t_{sin} = 6$  s C)  $t_{sin} = 8$  s.

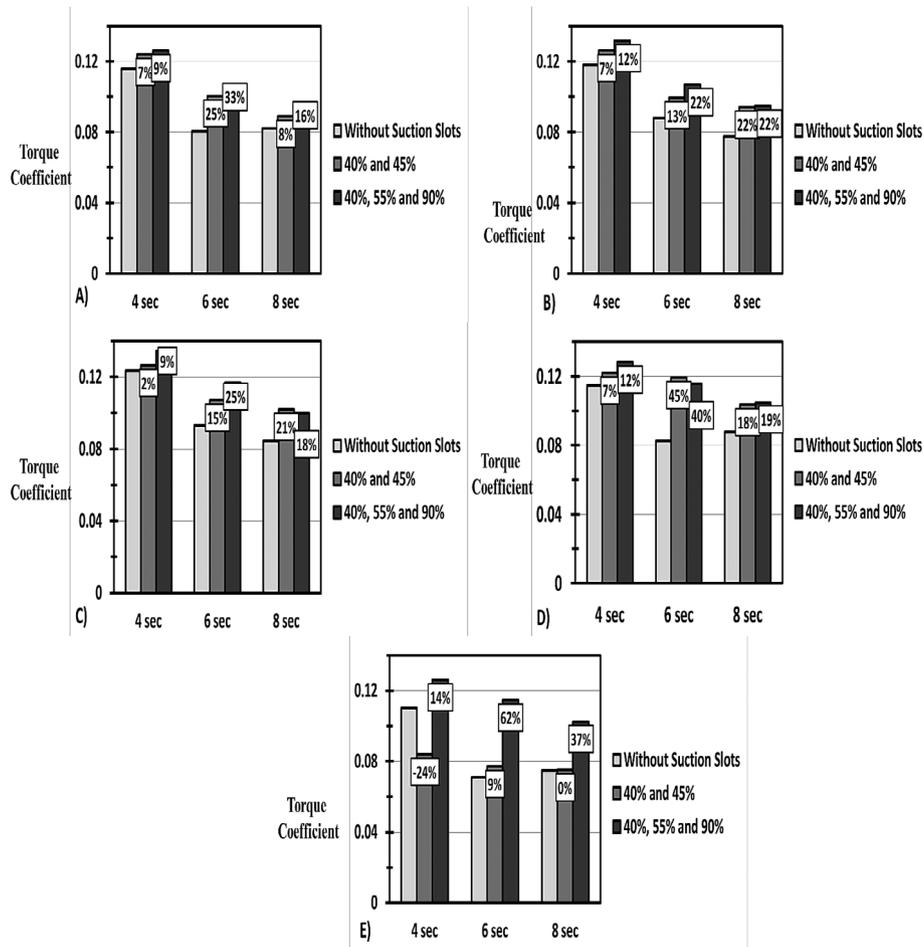


Fig. 30. The average torque coefficient at the compression cycle for the NACA0015 without and with suction slots at optimum locations with sinusoidal velocity A) 11.3 B) 11.7 C) 12.3 D) 13.6 E) 14.4.

the second law efficiency due to the different slots number and location.

The contours of global entropy generation rate around the NACA0015 at the instantaneous velocity 1.8 m/s for the accelerating (Fig. 24) and decelerating flow (Fig. 26) in addition 2.92 m/s (Fig. 25) were represented. Where, the 2.92 m/s was the maximum velocity which create the peak Reynolds number ( $2 \times 10^5$ ), and 1.8 m/s is approximately at the middle to compare between the accelerating and decelerating flow. It can be shown that the suction slots have a negative effect on the entropy generation, where the global entropy generation rate increases at the three stages, accelerating flow, maximum velocity and decelerating flow at  $13.6^\circ$ . The two-suction slots at ( $L_{SS}$  40% and 45%) and ( $L_{SS}$  45% and 49%)

have the lowest difference in global entropy generation rate by 32% at the accelerating flow in Fig. 24. Otherwise, the three-suction slots at ( $L_{SS}$  40%, 55% and 90%) have the highest difference in global entropy generation rate by 44% at the same Figure. However, the global entropy generation rate has lowest difference due to suction slots at the maximum velocity by 28% with the two-suction slots at ( $L_{SS}$  40% and 45%). Also, the highest value occurs due to the two suction slots at ( $L_{SS}$  40% and 44%) by 35% in Fig. 25.

From Fig. 26 it can be noted that the two suction slots at ( $L_{SS}$  40% and 45%) has the lowest difference in global entropy generation rate by 37% and the highest value occurs due to the three suction slots at ( $L_{SS}$  40%, 55% and 90%) by 53% at the decelerating flow. Finally, the global entropy generation rate around the NACA0015

Table 10

The global entropy generation rate at different angles of attack with wave time period equal to 4 s.

$S_G$ (W/K)	Angle of attack ( $^\circ$ )				
	11.3	11.7	12.3	13.6	14.4
$L_{SS} = 40\%$ and $L_{SS} = 45\%$	0.066	0.066	0.067	0.074	0.063
Average value	Before the stall (6.6%)		After the stall (6.8%)		
$L_{SS} = 40\%$ , $L_{SS} = 55\%$ and $L_{SS} = 90\%$	0.066	0.067	0.071	0.073	0.079
Average value	Before the stall (6.8%)		After the stall (7.6%)		

Table 11

The global entropy generation rate at different angles of attack with wave time period equal to 6 s.

$S_G$ (W/K)	Angle of attack ( $^\circ$ )				
	11.3	11.7	12.3	13.6	14.4
$L_{SS} = 40\%$ and $L_{SS} = 45\%$	0.065	0.067	0.067	0.068	0.062
Average value	Before the stall (6.6%)		After the stall (6.5%)		
$L_{SS} = 40\%$ , $L_{SS} = 55\%$ and $L_{SS} = 90\%$	0.067	0.071	0.068	0.071	0.080
Average value	Before the stall (6.9%)		After the stall (7.6%)		

**Table 12**

The global entropy generation rate at different angles of attack with wave time period equal to 8 s.

$S_G$ (W/K)	Angle of attack ( $^\circ$ )				
	11.3	11.7	12.3	13.6	14.4
$L_{ss} = 40\%$ and $L_{ss} = 45\%$	0.064	0.063	0.064	0.070	0.064
Average value	Before the stall (6.4%)			After the stall (6.7%)	
$L_{ss} = 40\%$ , $L_{ss} = 55\%$ and $L_{ss} = 90\%$	0.066	0.065	0.065	0.073	0.075
Average value	Before the stall (6.5%)			After the stall (7.4%)	

without and with suction slots have the highest value at the maximum velocity and the lowest value at the accelerating flow as a general. From Figs. 14–16 and 19 it can be noted that the attached multi-slots to the aerofoil lead to increase in velocity magnitude around the aerofoil, furthermore, it lead also to increase in the entropy generation in Figs. 24–26. Where, the entropy value depends on the velocity gradient see equation (21).

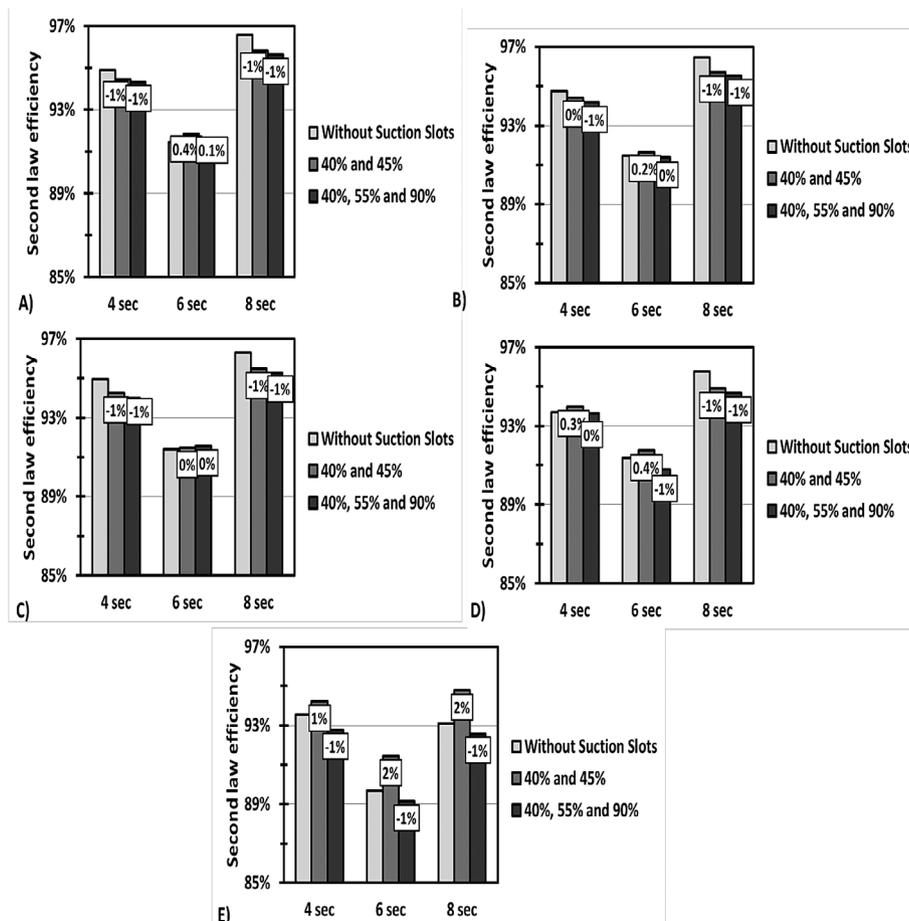
4.4. Comparative analysis based on conditions relevant to northern coast of Egypt

From the previous section, it can be concluded that the three-suction slots ( $L_{ss}$  40%, 55% and 90%) give maximum  $C_T$  and the two-suction slots ( $L_{ss}$  40% and 45%) have the lowest  $S_G$  from the

five scenarios, which give higher torque coefficient. Therefore, these two scenarios were investigated using the oscillating water system based on the real data from the site with different time periods and frequencies ( $f$  equal to 0.25, 0.167 and 1.25 Hz). The hysteretic behaviour due to the reciprocating flow and the total average torque coefficient during the cycle for aerofoil with suction slots at different time periods were shown in Fig. 27. It can be concluded that the aerofoil with three-suction slots ( $L_{ss}$  40%, 55% and 90%) give higher  $C_T$  than that with two-suction slots ( $L_{ss}$  40% and 45%) at 4, 6 and 8 s time period. Also, the increase in time period led to a decrease in the total average torque coefficient in general. At the time period equal to 4 s, the aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) has an average torque coefficient after the stall less than the aerofoil without suction slots by 8.5%. Furthermore, the aerofoil with two suction slots ( $L_{ss}$  40% and 45%) with 8 s time period has improvement in the total average torque coefficient before the stall by 17% and after the stall by 8%.

Figs. 28 and 29 show the instantaneous torque coefficient in addition to average torque coefficient at the accelerating and decelerating cycle for aerofoil with two-suction slots and with three-suction slots. These values were at angle of attack of 13.6° at different time periods (4 s, 6 s and 8 s). It can be seen that the improvement in the torque coefficient has the lowest value at the cycle with time period equal to 4 s. Furthermore, the torque coefficient value and improvement in the torque coefficient at decelerating flow are always higher than that at accelerating flow.

The total average torque coefficients during the compression cycle for different angles of attack were shown in Fig. 30. It can be



**Fig. 31.** The second law efficiency at the compression cycle for the NACA0015 without and with suction slots at optimum locations with sinusoidal velocity A) 11.3 B) 11.7 C) 12.3 D) 13.6 E) 14.4.

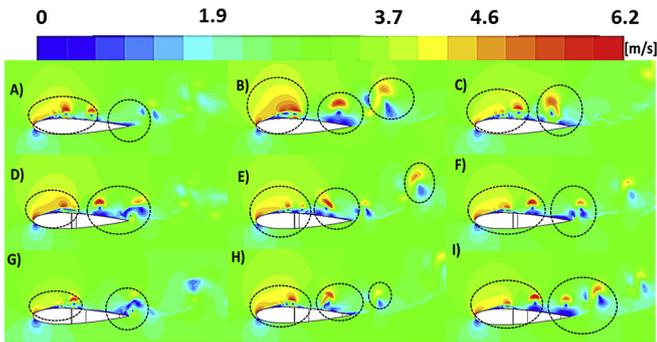


Fig. 32. Contour of velocity magnitude at 12.3° A), B) and C) NACA0015 without slots with  $t_{sin} = 4, 6$  and  $8$  s respectively D), E) and F) two slots with  $t_{sin} = 4, 6$  and  $8$  s respectively G), H) and I) three slots with  $t_{sin} = 4, 6$  and  $8$  s respectively.

observed that for all angles, the suction slot increases the torque coefficient except at the 14.4° Fig. 30 E), where the torque coefficient for the aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) was lower than that without suction slots by 24% at time period 4 s. Also, the torque coefficient at time period 8 s for the aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) was same for that without suction slots. The aerofoil with three-suction slots ( $L_{ss}$  40%, 55% and 90%) mostly has a higher torque coefficient than that of the two-suction slots ( $L_{ss}$  40% and 45%) at different time period.

Tables 10–12 show the comparison between the global entropy generation rate before and after the stall for the aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) and with three-suction slots ( $L_{ss}$  40%, 55% and 90%) at different time periods (4 s, 6 s and 8 s). There were no significant changes in the global entropy generation rate values due to the different time periods. As an average for all time period, the aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) has a lower difference in  $S_G$  before and after the stall than the aerofoil with three-suction slots ( $L_{ss}$  40%, 55% and 90%).

Suction slots have a negative effect on both the entropy behaviour and the second law efficiency. Therefore, most of cases at Fig. 31 have lower second law efficiency for aerofoils with slots than the aerofoils without slots. As it noted in the entropy behaviour, there were also no significant changes in the second law efficiency value due to the different slots number and location. However, the second low efficiency at 14.4° for the aerofoil with two suction slots ( $L_{ss}$  40% and 45%) was the highest value at 4, 6 and 8 s by 1%, 2% and 3% respectively, Fig. 31 E). The wave cycle with 8 s has the highest value of the second law efficiency as a general. On the other hand, the wave cycle with 6 s has the lowest value. The aerofoil with two-suction slots ( $L_{ss}$  40% and 45%) always has higher second law

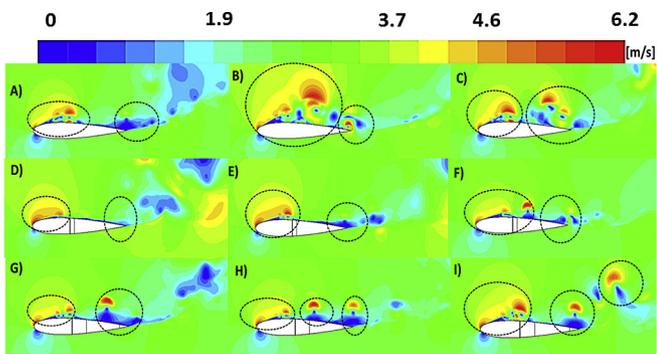


Fig. 33. Contour of velocity magnitude at 14.4° A), B) and C) for NACA0015 without slots with  $t_{sin} = 4, 6$  and  $8$  s respectively D), E) and F) two slots with  $t_{sin} = 4, 6$  and  $8$  s respectively G), H) and I) three slots with  $t_{sin} = 4, 6$  and  $8$  s respectively.

efficiency than that with three-suction slots ( $L_{ss}$  40%, 55% and 90%) at the different time periods.

The flow structures over the NACA0015 aerofoil in oscillating flow was shown in Fig. 32 at angle of attack equal to 12.3° (before the stall) and 14.4° (after the stall) in Fig. 33. The improvement effect of suction slot on flow structures was clear when comparing the NACA0015 without and with suction slots, especially in the separated layer regime at the end of aerofoil, which leads to an improvement in the separation regime.

## 5. Conclusions

More than 450 cases were solved to determine optimum location for multi suction slots based on the first and second law of thermodynamics. They aimed to investigate the effect of aerofoil with those optimum parameters on the entropy generation due to viscous dissipation as well as the torque coefficient and stall condition. After that, the comparative analysis based on real data relevant to northern coast of Egypt was applied using the aerofoil with optimum suction slot parameters.

The modeling results show that the optimum locations for two-suction slots aerofoil ( $L_{ss}$  of 40% and 44%), for three-suction slots aerofoil ( $L_{ss}$  of 40%, 55% and 90%), and for four-suction slots aerofoil ( $L_{ss}$  of 40%, 45%, 55%, and 60%). The three-suction slots aerofoil with  $L_{ss}$  of 40%, 55% and 90% gives the highest torque coefficient with 26.7% before the stall and 51% after the stall when compared to the aerofoil without suction slots. On the other hand, the two-suction slots aerofoil with  $L_{ss}$  of 40% and 45% gives the highest second law efficiency by 0.72% compared to the aerofoil without suction slots. The aerofoils with optimum locations for multi-suction slots under conditions relevant to northern coast of Egypt with different wave frequencies were investigated. For NACA0015, adding three-suction slots at optimum locations ( $L_{ss}$  of 40%, 55% and 90%) mostly gives a torque coefficient higher than that of adding two suction slots at optimum locations ( $L_{ss}$  of 40% and 45%) for different  $t_{sin}$  (4, 6 and 8 s). However, adding two-suction slots at optimum locations ( $L_{ss}$  of 40% and 45%) always gives a second law efficiency higher than that of adding three-suction slots at optimum locations ( $L_{ss}$  of 40%, 55% and 90%) for different  $t_{sin}$  (4, 6 and 8 s).

The main reason behind the improvement in the torque coefficient after the stall is due to the delay of stall condition. The suction slot increases the torque coefficient and delays the stall angle which further leads to an increase of first law efficiency. On the other hand, it increases the entropy generation rate which leads to decreasing the second law efficiency. The main reason also behind this increase in the entropy generation rate is due to the increases in velocity magnitude around the aerofoil lead to increase also in the entropy generation. Where, the entropy value depends on the velocity gradient. At the present study, the optimization parameters have been varied within a certain range with a fixed increment and all parameter values have been analyzed. An alternative approach that could save time and effort would be to use an automated optimization technique, [98]. Furthermore, Wells turbine impeller using the suction slot with optimum parameters needs to be investigated experimentally in the future. Finally, the operating conditions for the northern coast of Egypt are very suitable for the oscillating system with Wells turbine as a wave energy extractor. So, it is essential that to look at the wave energy in Egypt as the way to reduce fossil fuel usage.

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